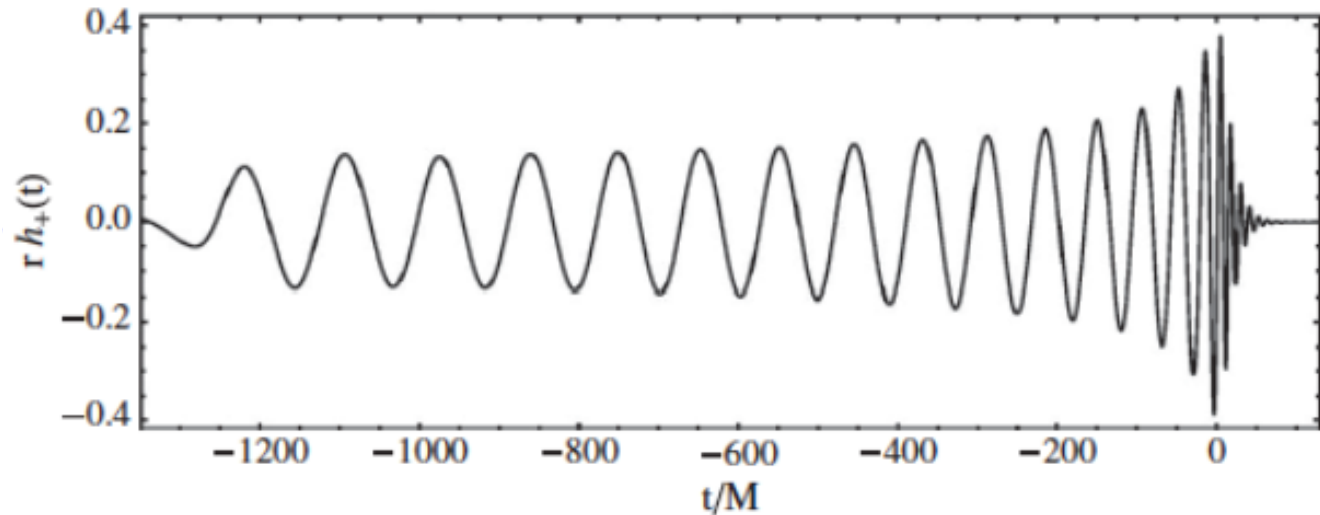
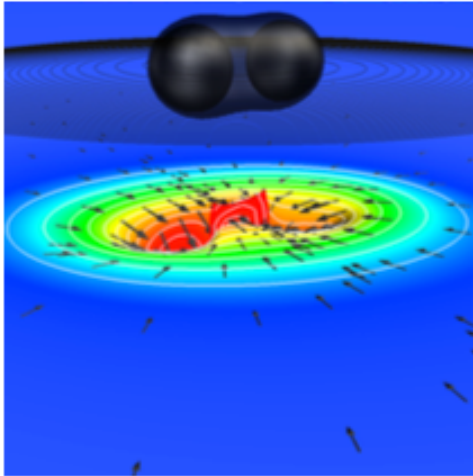


중력파 파형 & 모델

Hee-Suk Cho (Pusan National University)

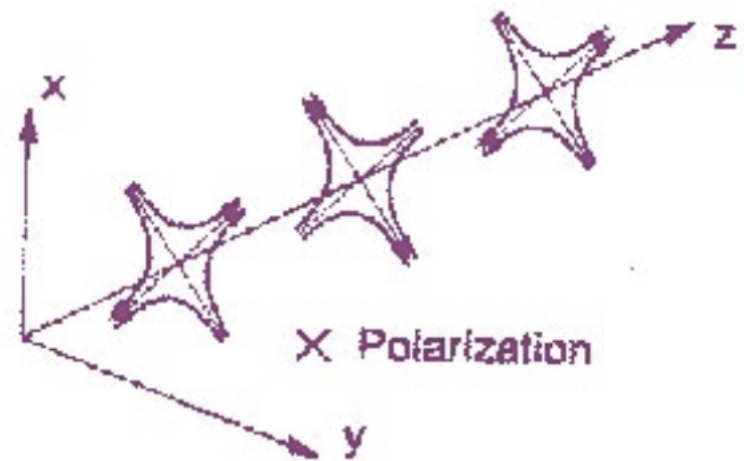
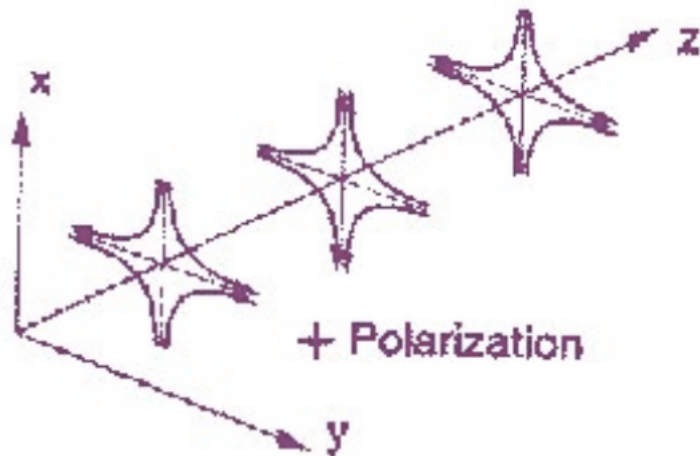
2025 GWNR 겨울학교 2025.1.22



- Binary Black Hole (BBH) waveforms
 - nonspinning
 - precessing (amplitude modulation)
 - PN phase evolution
 - waveform models

GW polarizations

$$h_+ = -\frac{2G\mu}{c^2 R} x(1 + \cos^2 \Theta) \cos 2\psi$$
$$h_\times = -\frac{2G\mu}{c^2 R} x(2 \cos \Theta) \sin 2\psi$$



GW polarizations: nonspinning

$$h_+ = -\frac{2G\mu}{c^2 R} x (1 + \cos^2 \Theta) \cos 2\psi$$
$$h_\times = -\frac{2G\mu}{c^2 R} x (2 \cos \Theta) \sin 2\psi$$

Diagram illustrating the parameters in the GW polarization equations:

- PN parameter** (blue text) points to the x term in both equations.
- Inclination** (red text) points to the $\cos \Theta$ term in both equations.
- Orbital phase** (purple text) points to the $\cos 2\psi$ and $\sin 2\psi$ terms.

$$x \equiv \left(\frac{G m \omega}{c^3} \right)^{2/3}$$

Waveform: nonspinning

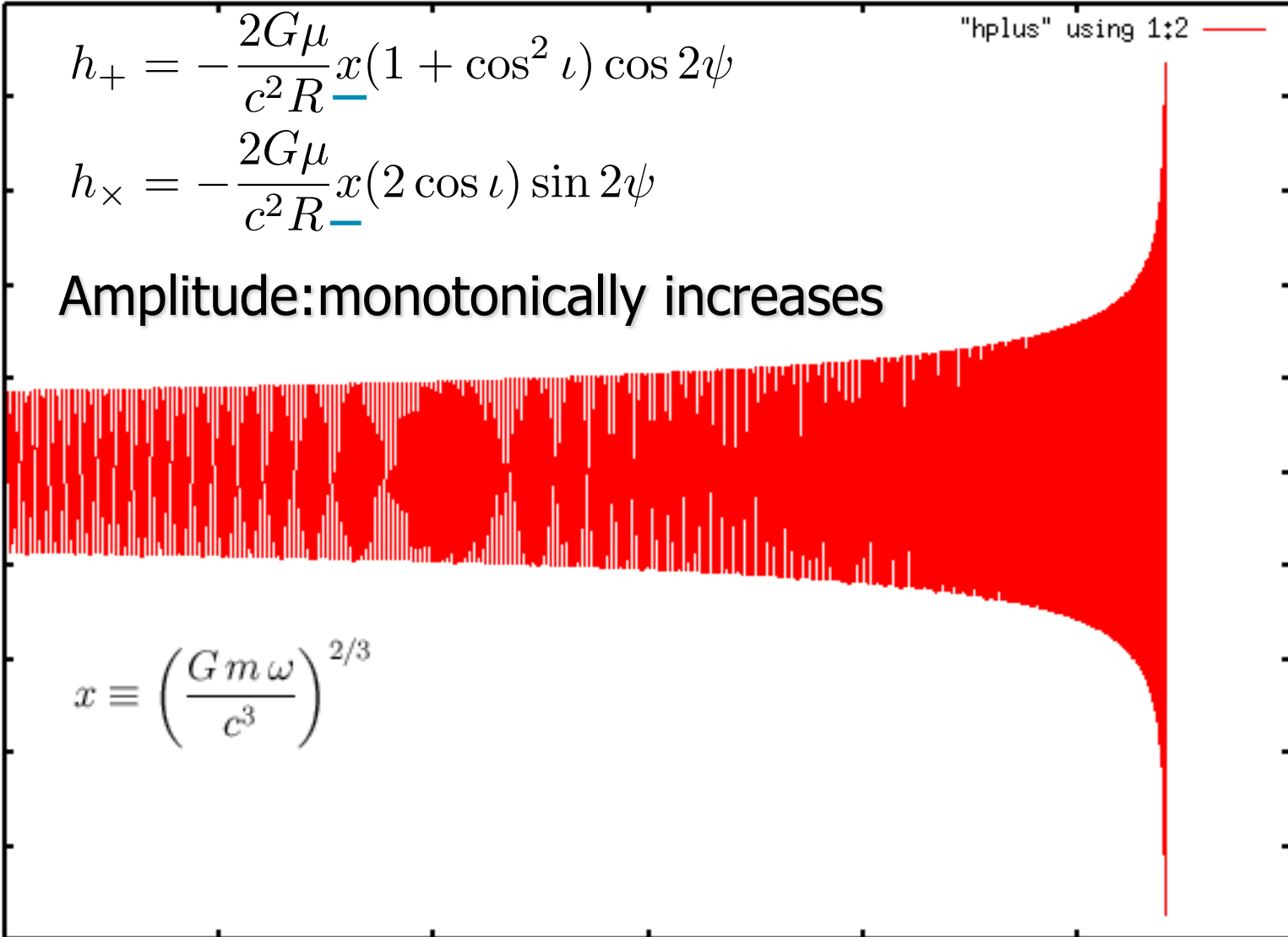
$$h_+ = -\frac{2G\mu}{c^2 R} x (1 + \cos^2 \iota) \cos 2\psi$$

$$h_\times = -\frac{2G\mu}{c^2 R} x (2 \cos \iota) \sin 2\psi$$

Amplitude: monotonically increases

$$x \equiv \left(\frac{G m \omega}{c^3} \right)^{2/3}$$

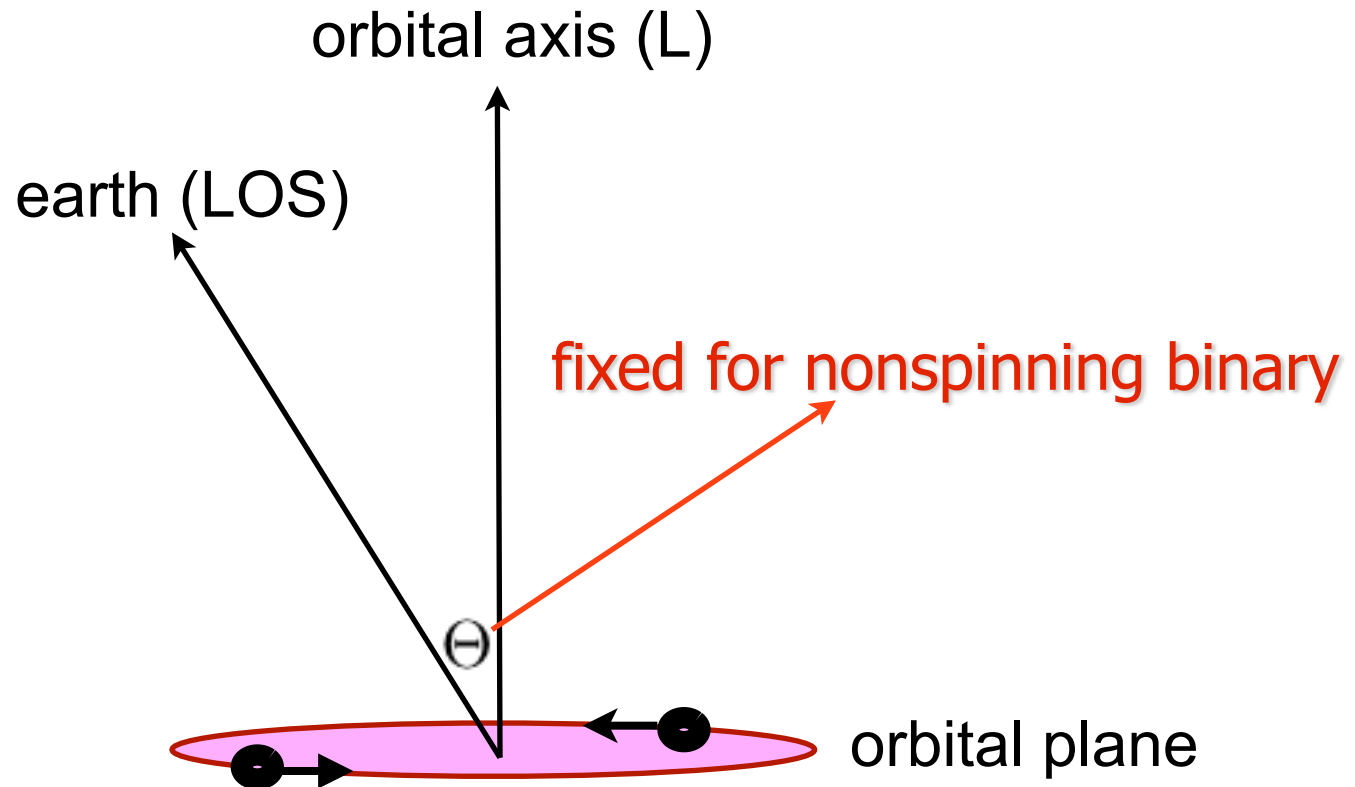
"hplus" using 1:2 —



Inclination (Θ)

$$h_+ = -\frac{2G\mu}{c^2 R} x (1 + \cos^2 \underline{\Theta}) \cos 2\psi$$

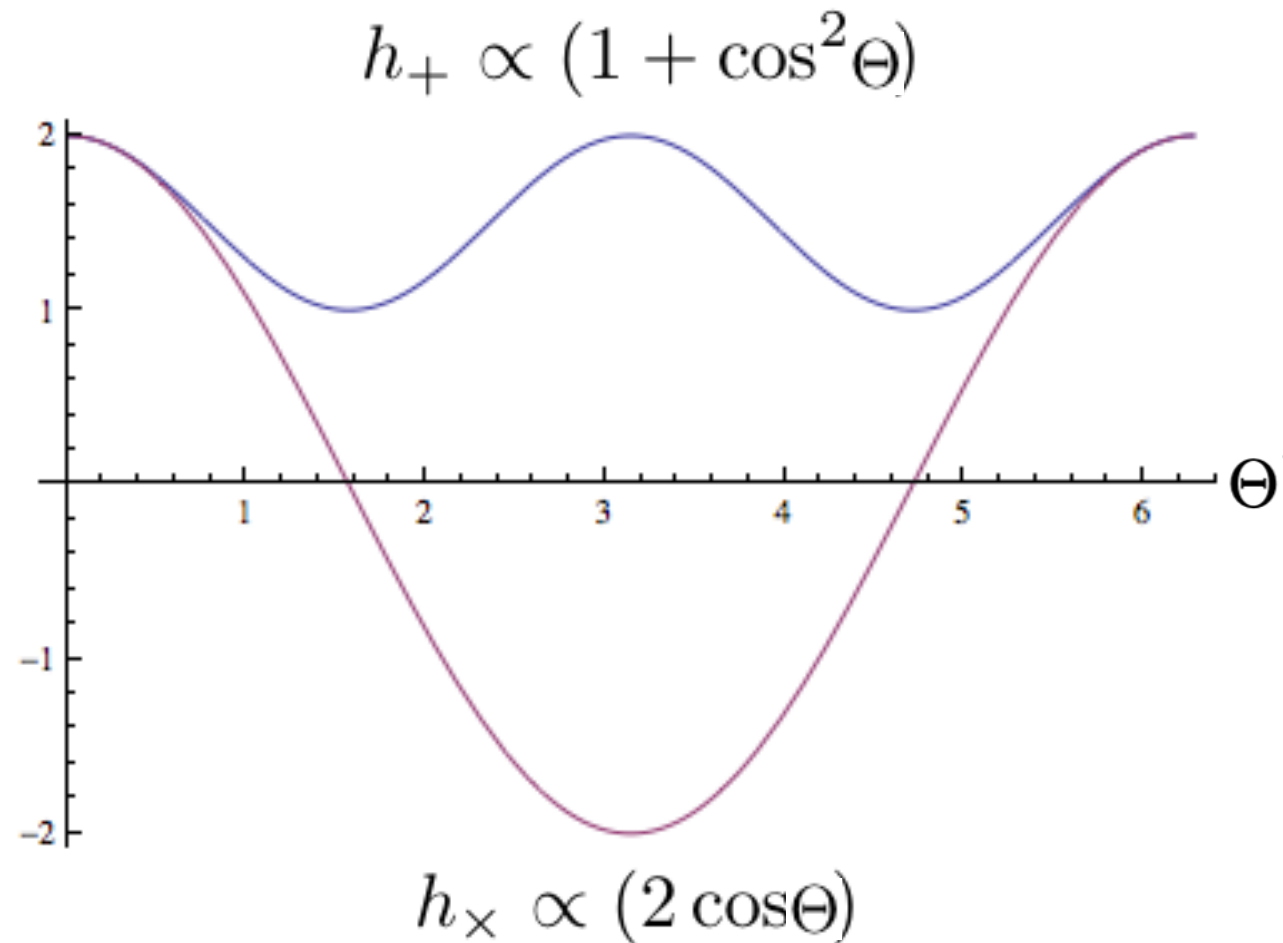
$$h_\times = -\frac{2G\mu}{c^2 R} x (2 \cos \underline{\Theta}) \sin 2\psi$$



Dependence on inclination

$$h_+ = -\frac{2G\mu}{c^2 R} x \underline{(1 + \cos^2 \Theta)} \cos 2\psi$$

$$h_\times = -\frac{2G\mu}{c^2 R} x \underline{(2 \cos \Theta)} \sin 2\psi$$

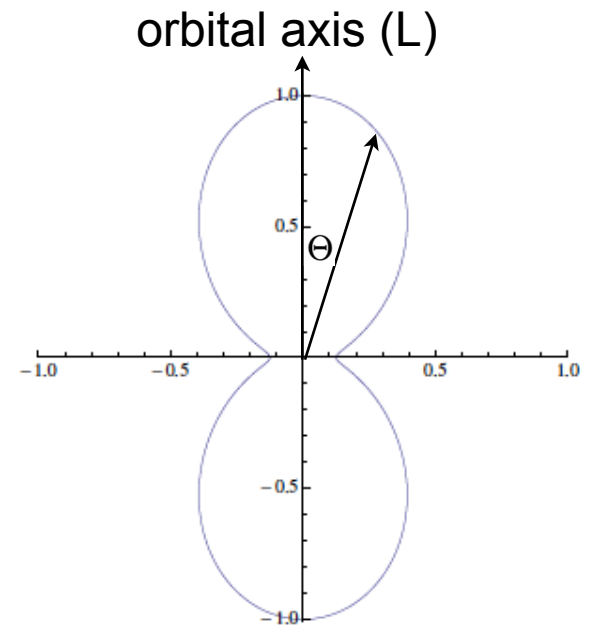
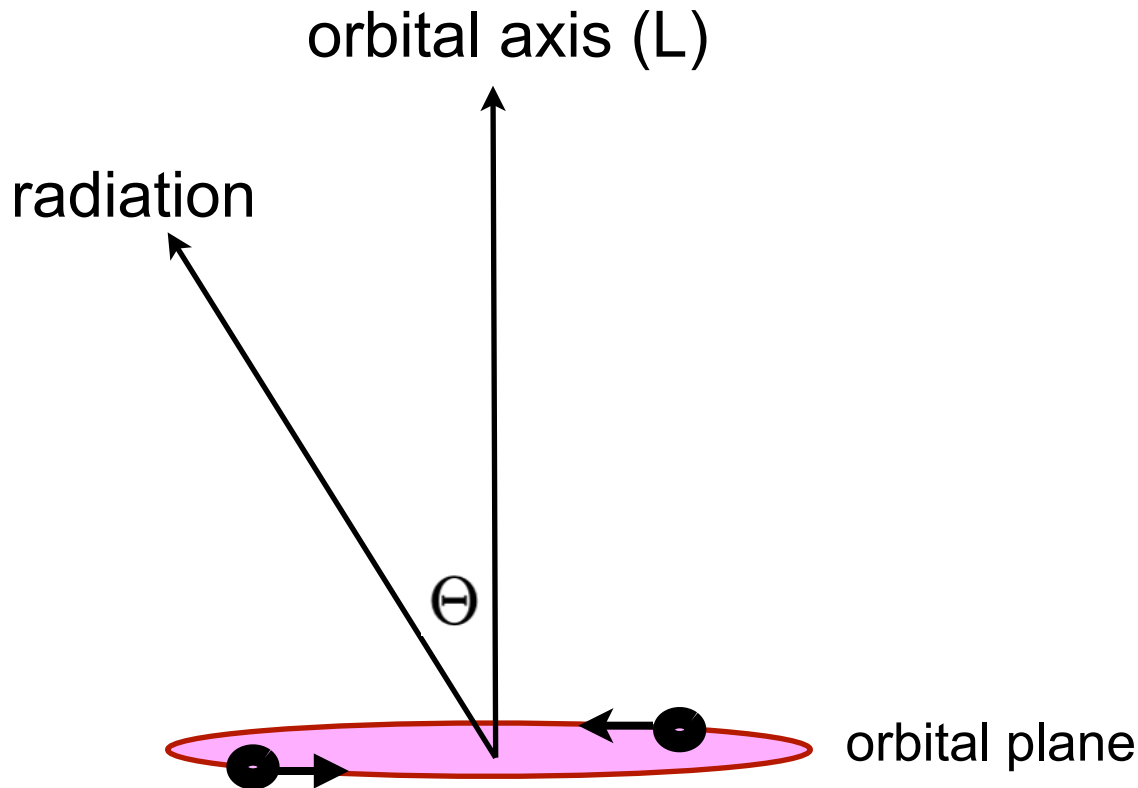


Radiation power (source frame)

$$h_+ = -\frac{2G\mu}{c^2 R} x(1 + \cos^2 \Theta) \cos 2\psi$$

$$h_\times = -\frac{2G\mu}{c^2 R} x(2 \cos \Theta) \sin 2\psi$$

$$\left(\frac{dP}{d\Omega}\right) = \frac{2G\mu^2 a^4 \omega^6}{\pi c^5} g(\Theta),$$



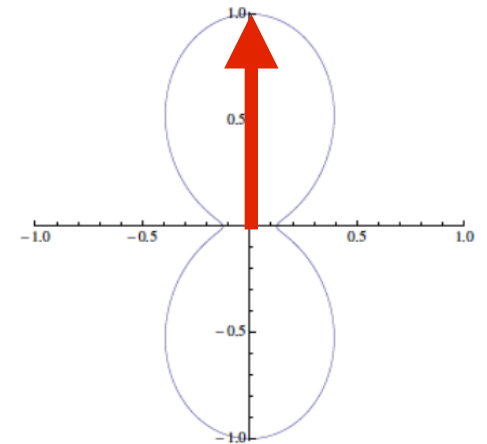
Non spinning GW polarizations (theta=0)

maximum & same amplitudes

"hplus" using 1:2
"hcross" using 1:3

hplus

hcross

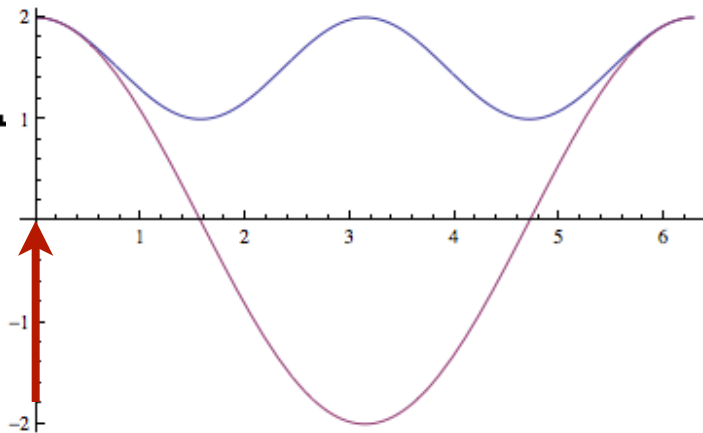


orbital axis (L)

LOS

$\Theta=0$

circular polarization



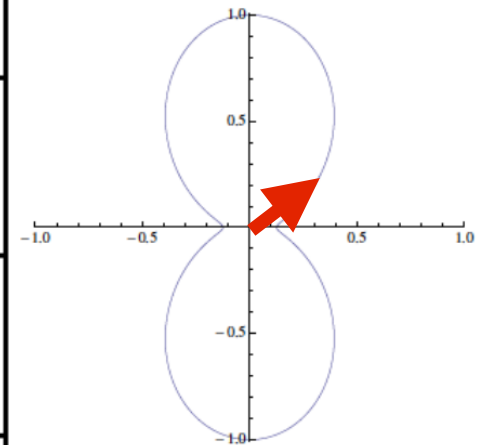
Non spinning polarizations (theta=pi/3)

reduced & different amplitudes

"hplus" using 1:2
"hcross" using 1:3

hplus

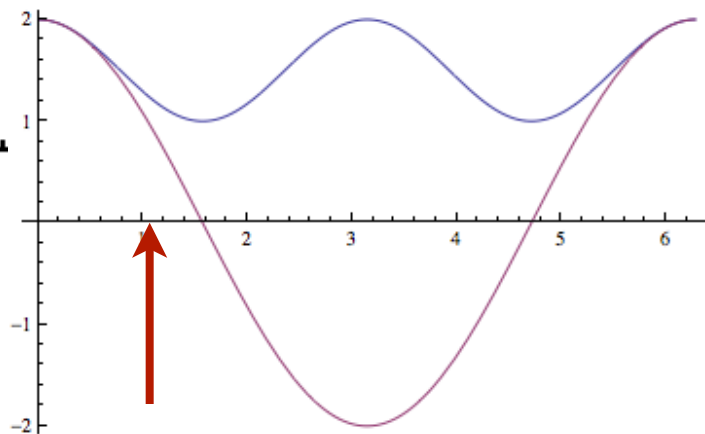
hcross



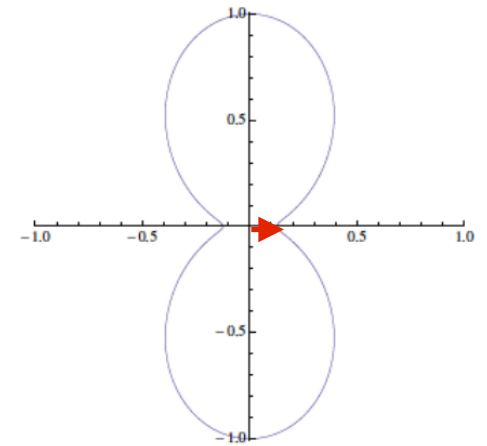
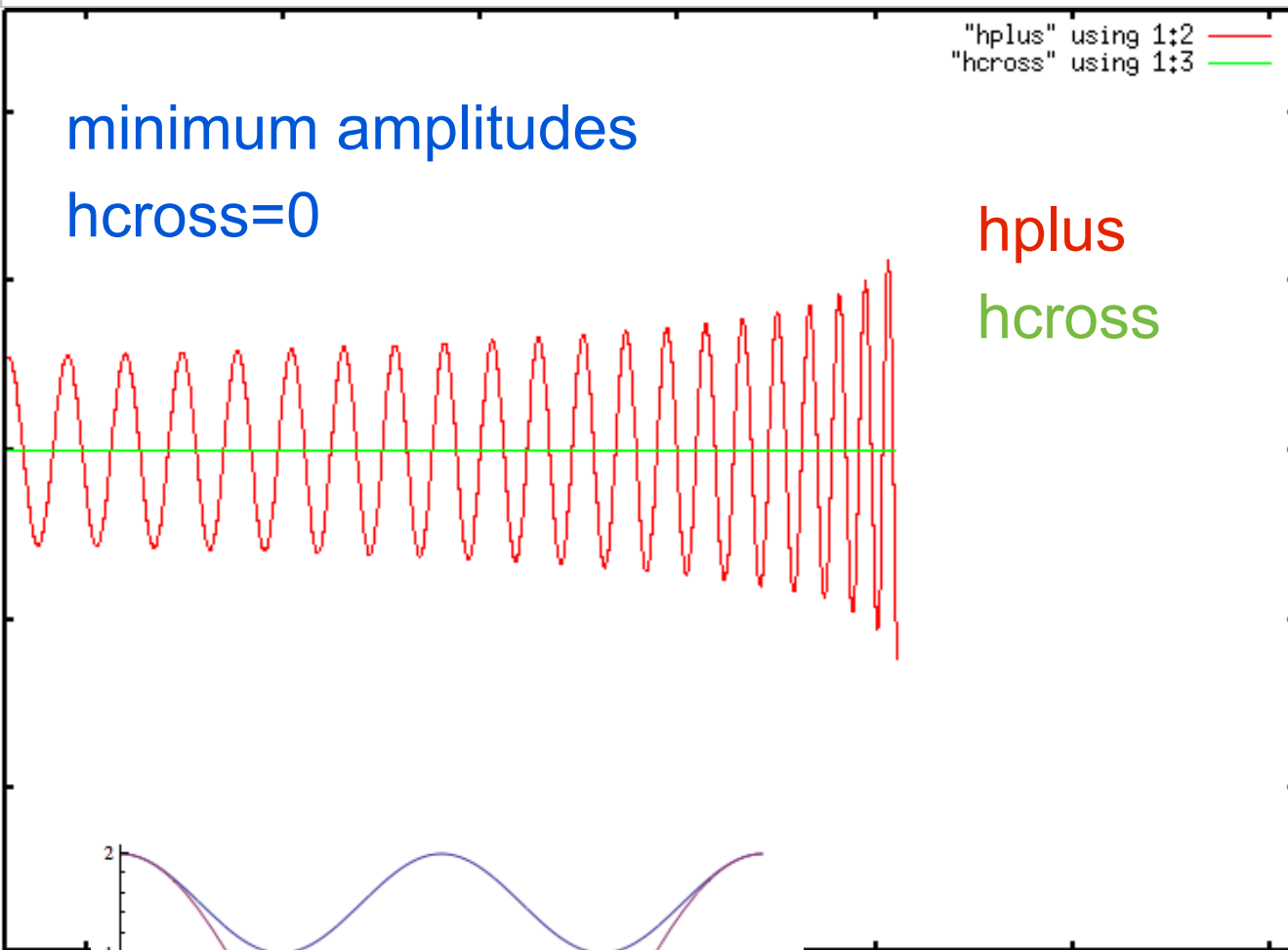
orbital axis (L)

LOS

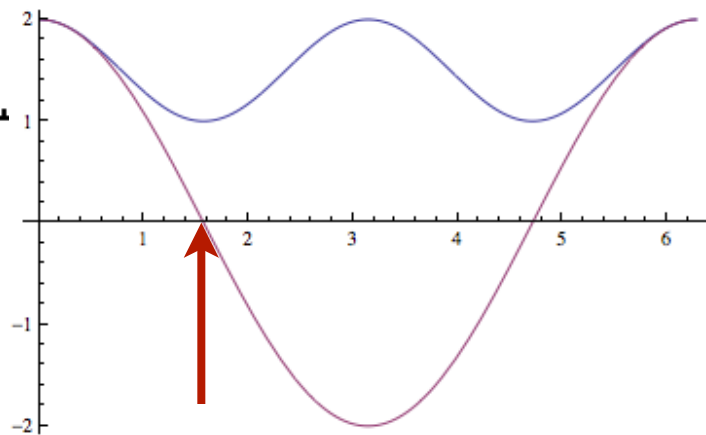
$\Theta = \pi/3$



Non spinning polarizations (theta=pi/2)



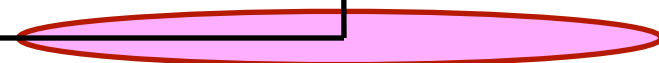
orbital axis (L)



linear polarization

LOS

$\Theta = \pi/2$



distance-inclination degeneracy

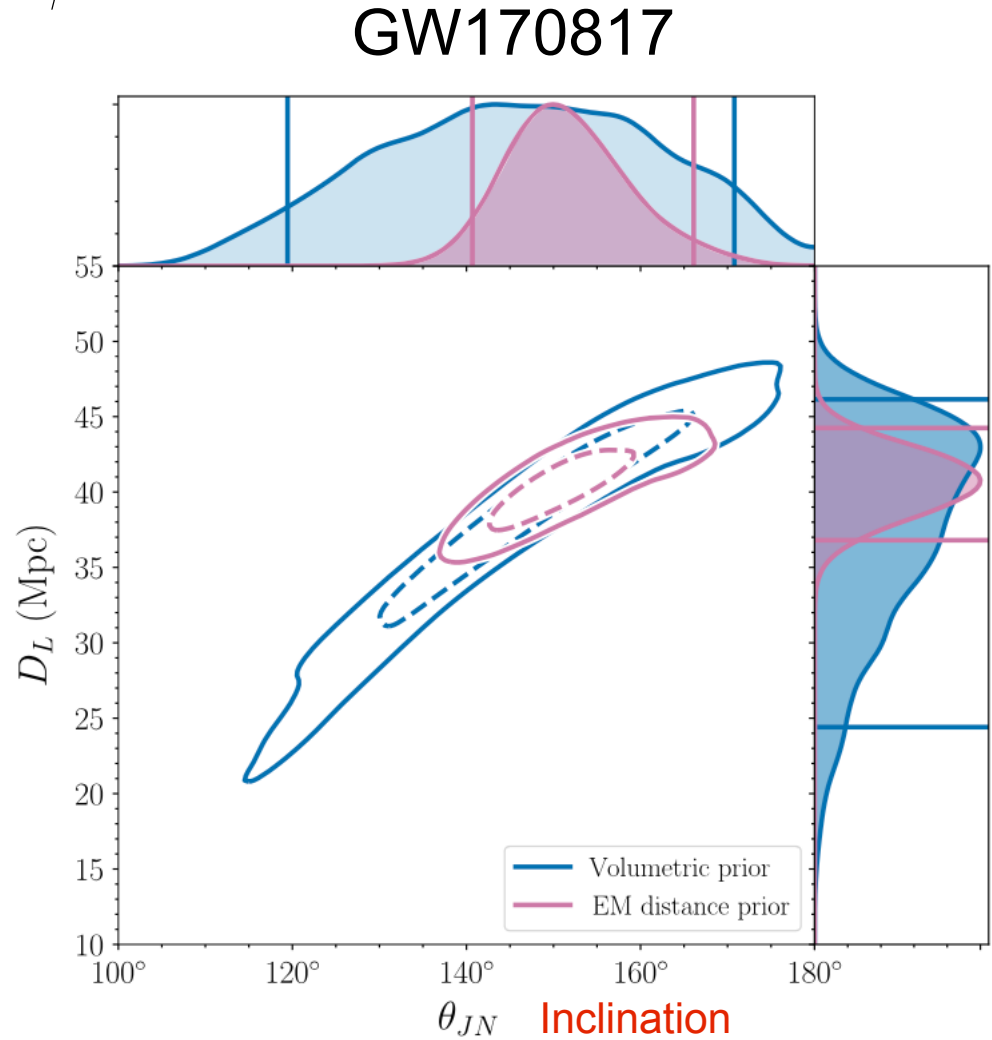
$$h_+ = -\frac{2G\mu}{c^2 R} x(1 + \cos^2 \Theta) \cos 2\psi$$

$$h_\times = -\frac{2G\mu}{c^2 R} x(2 \cos \Theta) \sin 2\psi$$

Distance

Inclination

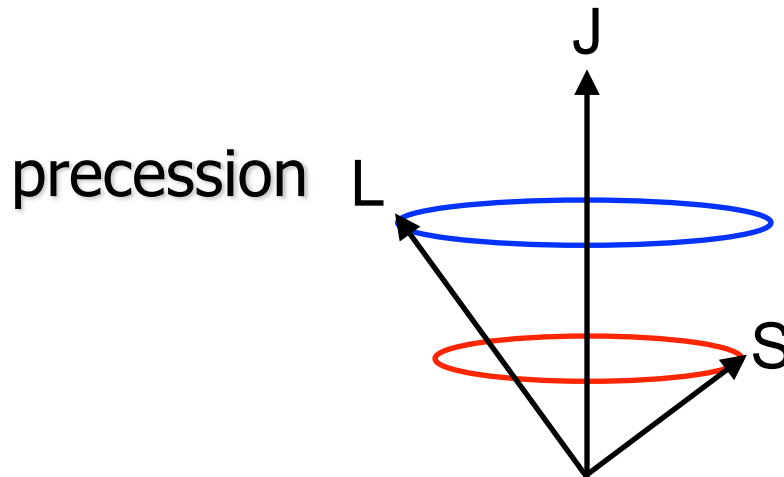
Distance



Precession of a spinning binary

Spin-Orbit coupling, Spin-Spin coupling

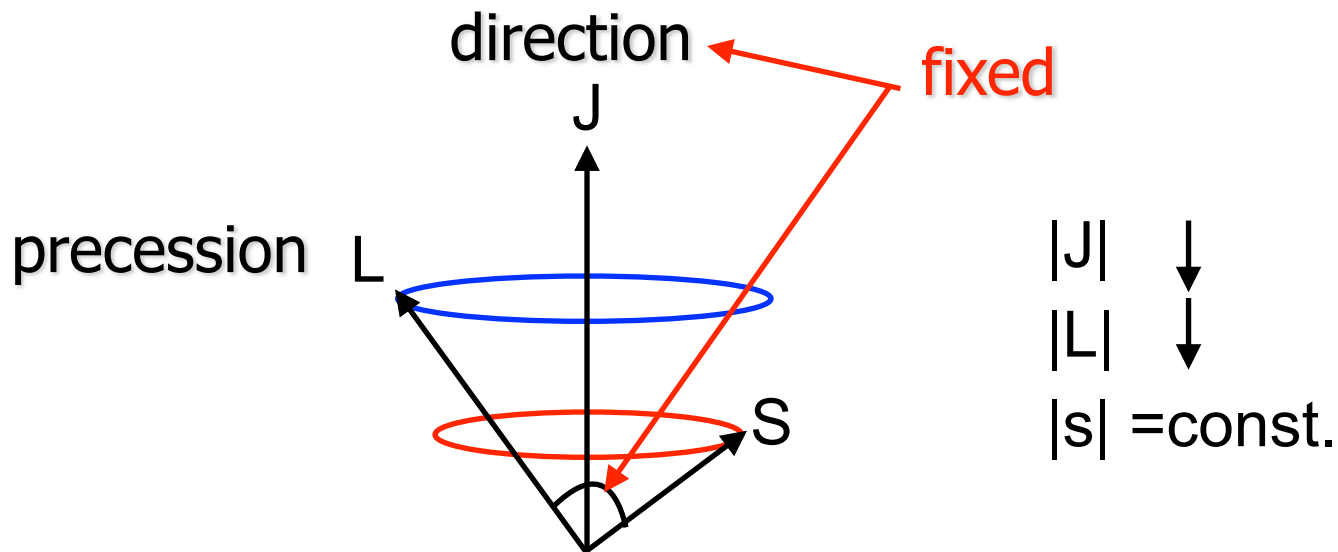
$$\begin{aligned}\dot{\mathbf{S}}_1 &= \frac{(M\omega)^2}{2M} \left\{ \eta (M\omega)^{-1/3} \left(4 + 3\frac{m_2}{m_1} \right) \hat{\mathbf{L}}_N + \frac{1}{M^2} \left[\mathbf{S}_2 - 3(\mathbf{S}_2 \cdot \hat{\mathbf{L}}_N) \hat{\mathbf{L}}_N \right] \right\} \times \mathbf{S}_1 \\ \dot{\mathbf{S}}_2 &= \frac{(M\omega)^2}{2M} \left\{ \eta (M\omega)^{-1/3} \left(4 + 3\frac{m_1}{m_2} \right) \hat{\mathbf{L}}_N + \frac{1}{M^2} \left[\mathbf{S}_1 - 3(\mathbf{S}_1 \cdot \hat{\mathbf{L}}_N) \hat{\mathbf{L}}_N \right] \right\} \times \mathbf{S}_2 \\ \dot{\hat{\mathbf{L}}}_N &= -\frac{(M\omega)^{1/3}}{\eta M^2} \dot{\mathbf{S}} = \frac{\omega^2}{2M} \left\{ \left[\left(4 + 3\frac{m_2}{m_1} \right) \mathbf{S}_1 + \left(4 + 3\frac{m_1}{m_2} \right) \mathbf{S}_2 \right] \times \hat{\mathbf{L}}_N \right. \\ &\quad \left. - \frac{3\omega^{1/3}}{\eta M^{5/3}} \left[(\mathbf{S}_2 \cdot \hat{\mathbf{L}}_N) \mathbf{S}_1 + (\mathbf{S}_1 \cdot \hat{\mathbf{L}}_N) \mathbf{S}_2 \right] \times \hat{\mathbf{L}}_N \right\}\end{aligned}$$



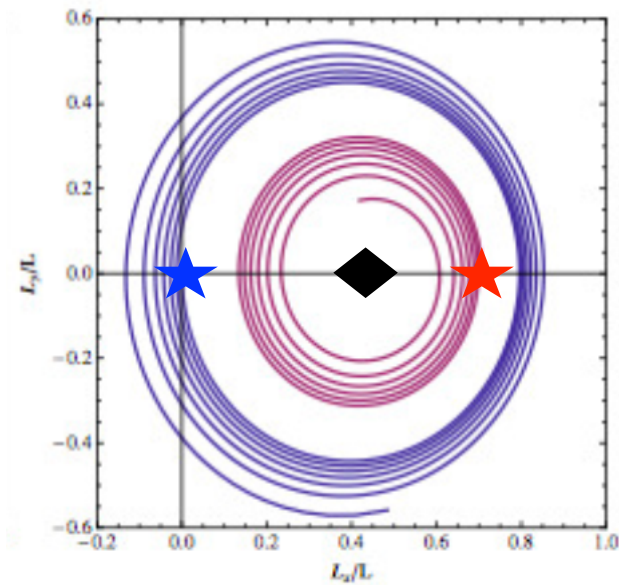
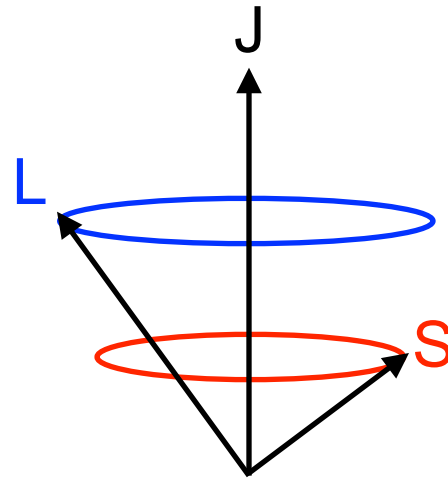
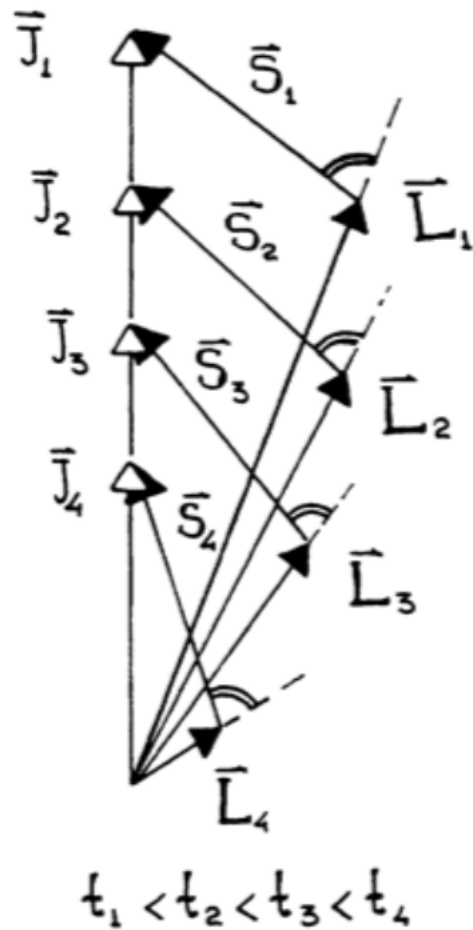
Precession of a spinning binary

Spin-Orbit coupling, Spin-Spin coupling

$$\begin{aligned}\dot{\mathbf{S}}_1 &= \frac{(M\omega)^2}{2M} \left\{ \eta (M\omega)^{-1/3} \left(4 + 3\frac{m_2}{m_1} \right) \hat{\mathbf{L}}_N + \frac{1}{M^2} \left[\mathbf{S}_2 - 3(\mathbf{S}_2 \cdot \hat{\mathbf{L}}_N) \hat{\mathbf{L}}_N \right] \right\} \times \mathbf{S}_1 \\ \dot{\mathbf{S}}_2 &= \frac{(M\omega)^2}{2M} \left\{ \eta (M\omega)^{-1/3} \left(4 + 3\frac{m_1}{m_2} \right) \hat{\mathbf{L}}_N + \frac{1}{M^2} \left[\mathbf{S}_1 - 3(\mathbf{S}_1 \cdot \hat{\mathbf{L}}_N) \hat{\mathbf{L}}_N \right] \right\} \times \mathbf{S}_2 \\ \dot{\hat{\mathbf{L}}}_N &= -\frac{(M\omega)^{1/3}}{\eta M^2} \dot{\mathbf{S}} = \frac{\omega^2}{2M} \left\{ \left[\left(4 + 3\frac{m_2}{m_1} \right) \mathbf{S}_1 + \left(4 + 3\frac{m_1}{m_2} \right) \mathbf{S}_2 \right] \times \hat{\mathbf{L}}_N \right. \\ &\quad \left. - \frac{3\omega^{1/3}}{\eta M^{5/3}} \left[(\mathbf{S}_2 \cdot \hat{\mathbf{L}}_N) \mathbf{S}_1 + (\mathbf{S}_1 \cdot \hat{\mathbf{L}}_N) \mathbf{S}_2 \right] \times \hat{\mathbf{L}}_N \right\}\end{aligned}$$



Precession of a spinning binary



GW Polarization of spinning binary

$$h_+ = -\frac{2G\mu}{c^2 R} x [C_+ \cos 2\psi + S_+ \sin 2\psi]$$

$$h_\times = -\frac{2G\mu}{c^2 R} x [C_\times \cos 2\psi + S_\times \sin 2\psi]$$

$$C_+ = \frac{1}{2} \cos^2 \Theta (\sin^2 \alpha - \cos^2 i \cos^2 \alpha) + \frac{1}{2} (\cos^2 i \sin^2 \alpha - \cos^2 \alpha) - \frac{1}{2} \sin^2 \Theta \sin^2 i - \frac{1}{4} \sin 2\Theta \sin 2i \cos \alpha,$$

$$S_+ = \frac{1}{2} (1 + \cos^2 \Theta) \cos i \sin 2\alpha + \frac{1}{2} \sin 2\Theta \sin i \sin \alpha,$$

$$C_\times = -\frac{1}{2} \cos \Theta \sin 2\alpha (1 + \cos^2 i) - \frac{1}{2} \sin \Theta \sin 2i \sin \alpha,$$

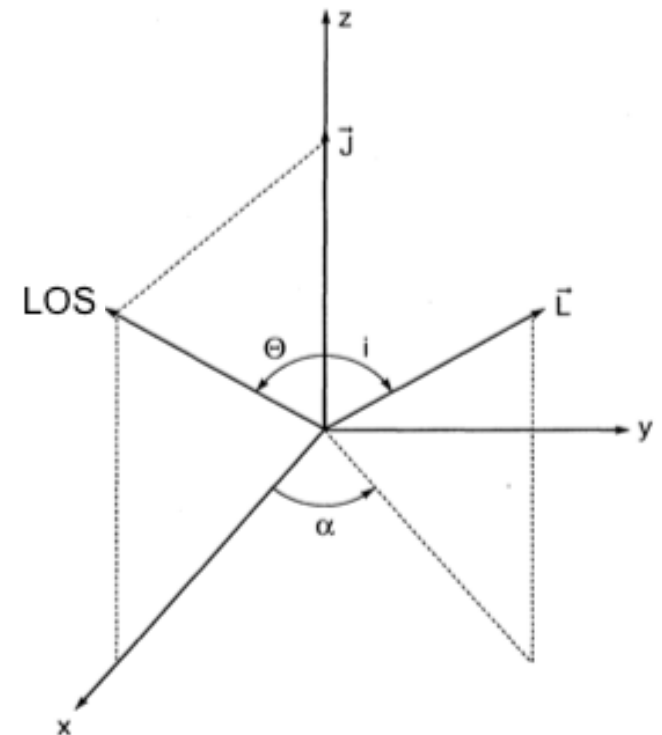
$$S_\times = -\cos \Theta \cos i \cos 2\alpha - \sin \Theta \sin i \cos \alpha,$$

ι, α, Θ : vary in time
--> precessional motion

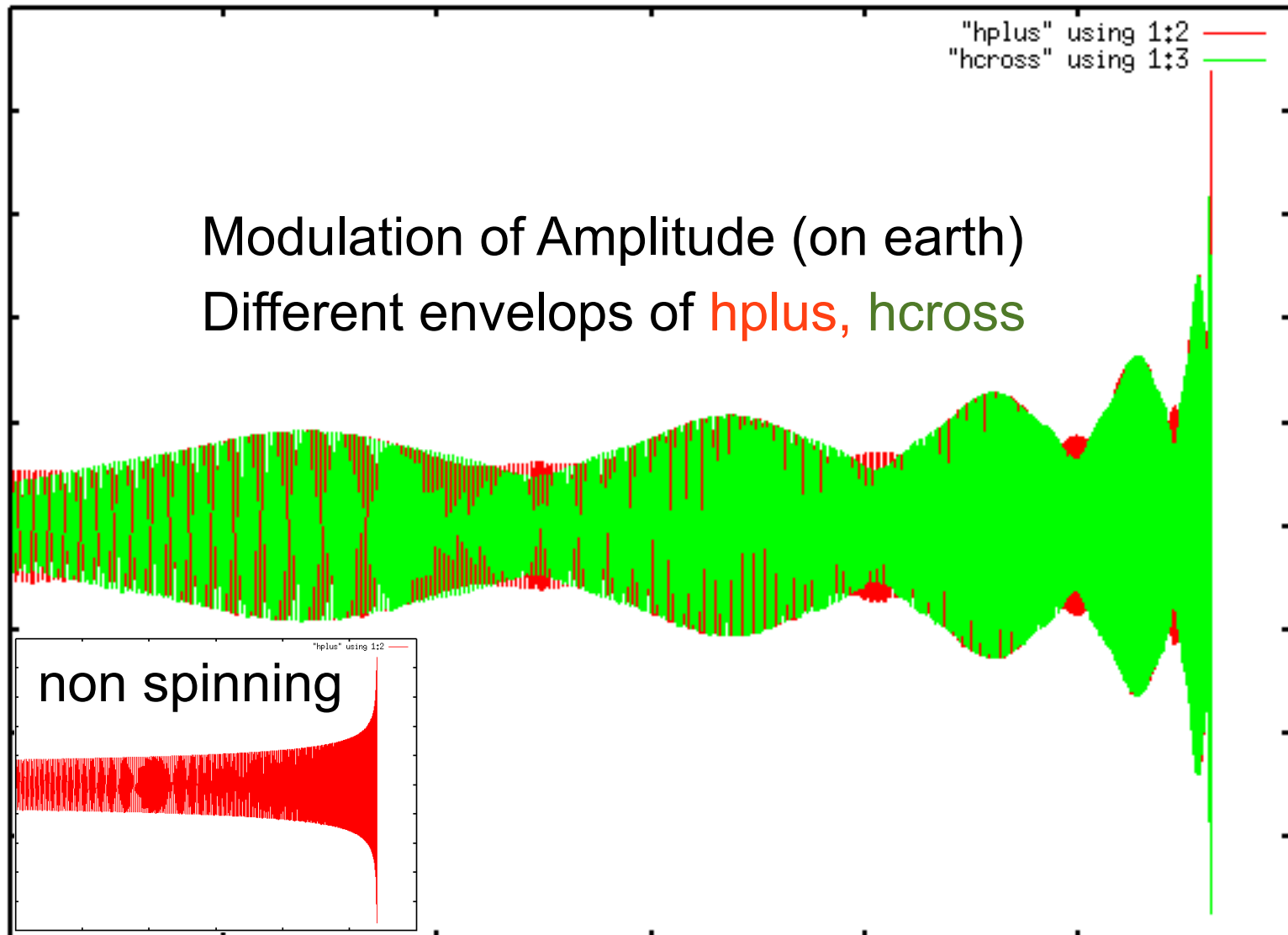
non spinning

$$h_+ = -\frac{2G\mu}{c^2 R} x (1 + \cos^2 \Theta) \cos 2\psi$$

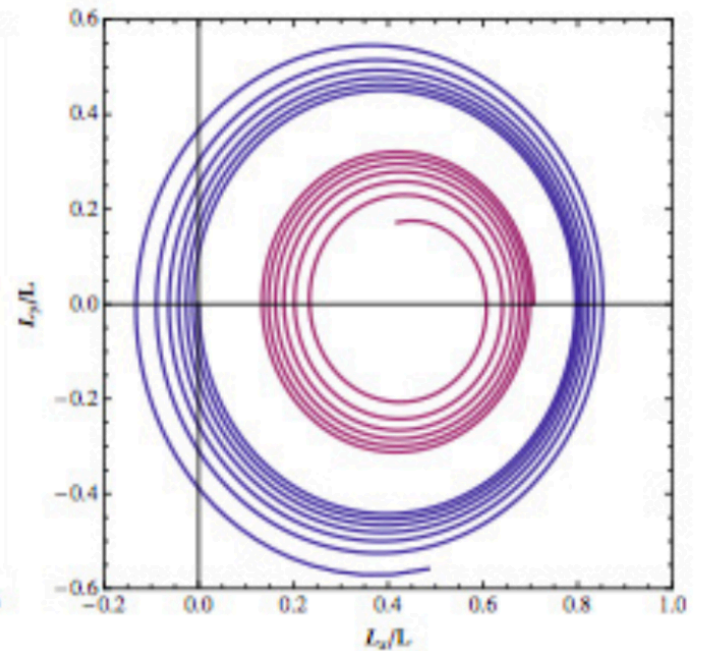
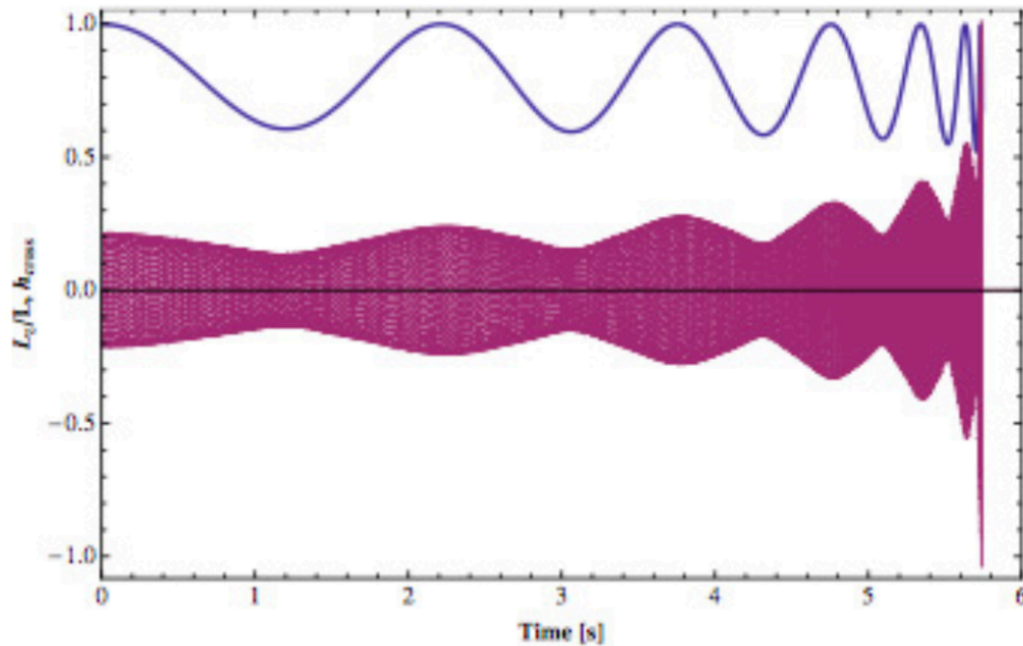
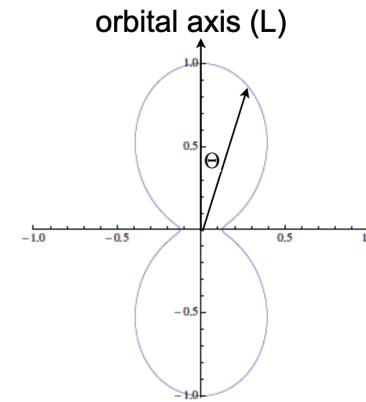
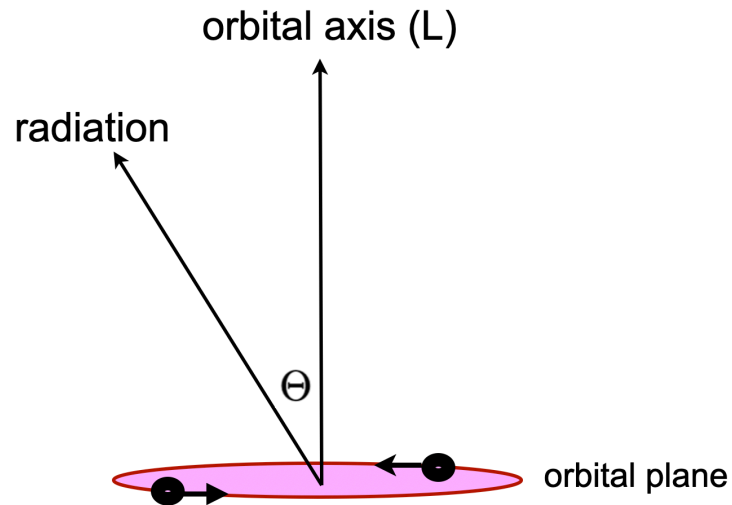
$$h_\times = -\frac{2G\mu}{c^2 R} x (2 \cos \Theta) \sin 2\psi$$



Waveform of spinning binary



Radiation power (source frame)



Modulation magnitude

Variation of L_N depends on S , $S \cdot L_N$, $S \times L_N$

---> Precession effect depends on

- 1) Spin magnitude
- 2) Angle between L_N and S

$$\dot{\underline{\mathbf{L}}}_N = -\frac{(M\omega)^{1/3}}{\eta M^2} \dot{\underline{\mathbf{S}}} = \frac{\omega^2}{2M} \left\{ \left[\left(4 + 3\frac{m_2}{m_1}\right) \underline{\mathbf{S}}_1 + \left(4 + 3\frac{m_1}{m_2}\right) \underline{\mathbf{S}}_2 \right] \times \hat{\underline{\mathbf{L}}}_N - \frac{3\omega^{1/3}}{\eta M^{5/3}} \left[(\underline{\mathbf{S}}_2 \cdot \hat{\underline{\mathbf{L}}}_N) \underline{\mathbf{S}}_1 + (\underline{\mathbf{S}}_1 \cdot \hat{\underline{\mathbf{L}}}_N) \underline{\mathbf{S}}_2 \right] \times \hat{\underline{\mathbf{L}}}_N \right\}$$

Modulation magnitude

Variation of L_N depends on S , $S \cdot L_N$, $S \times L_N$

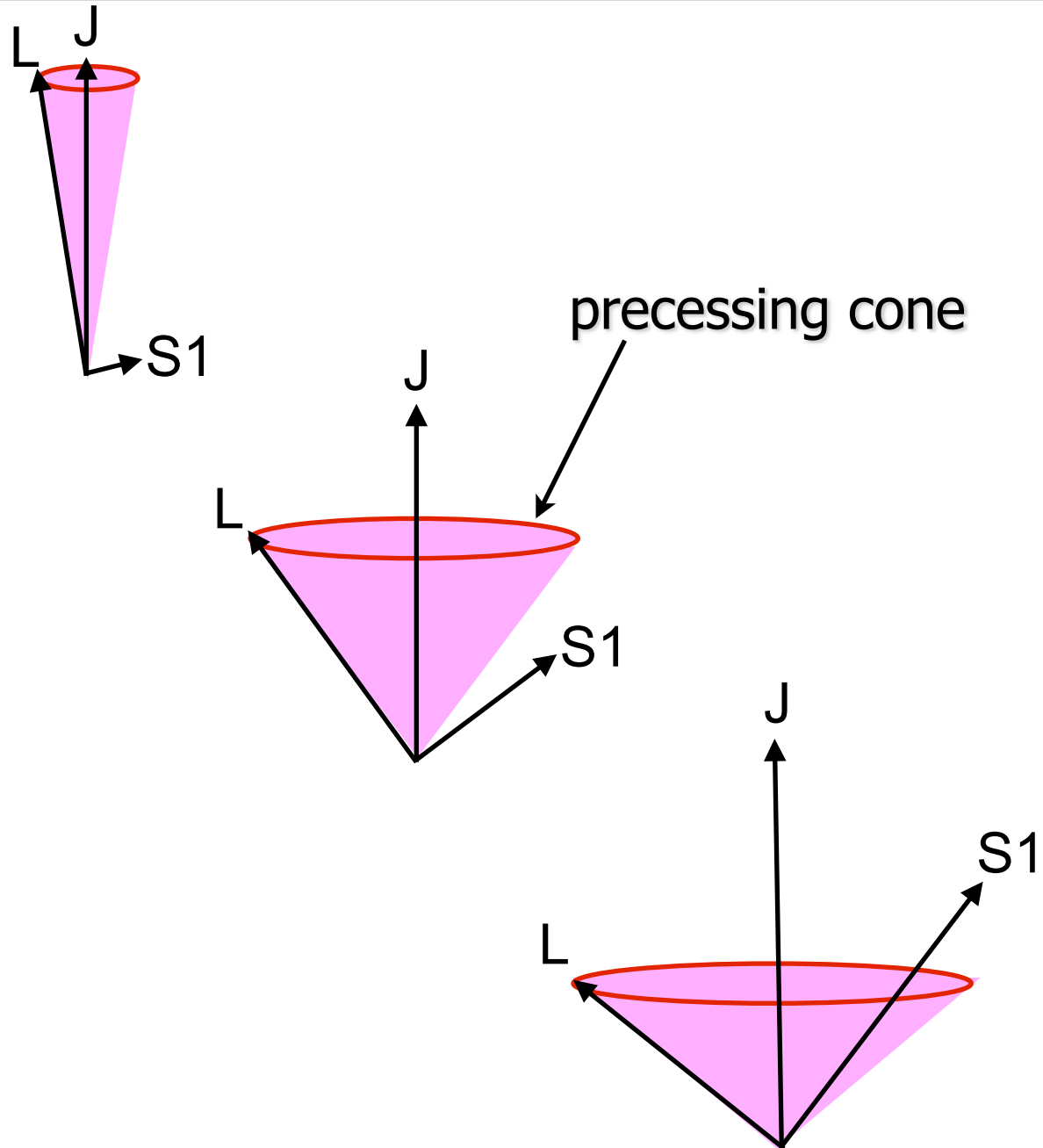
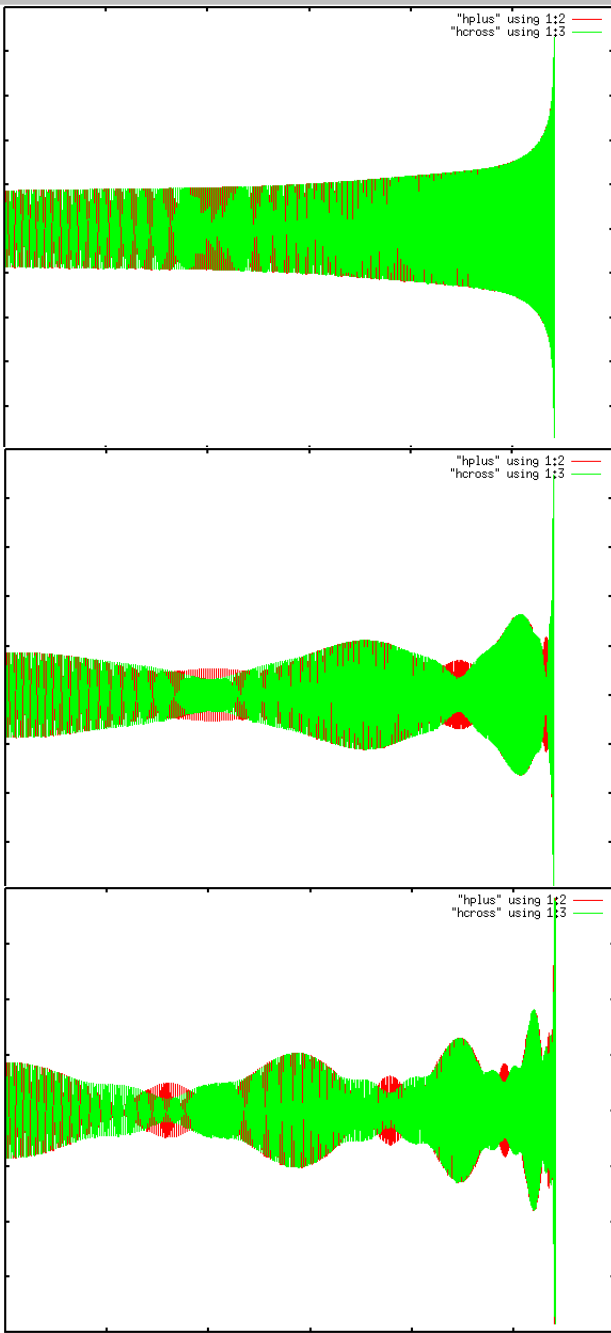
---> Precession effect depends on

- 1) Spin magnitude
- 2) Angle between L_N and S

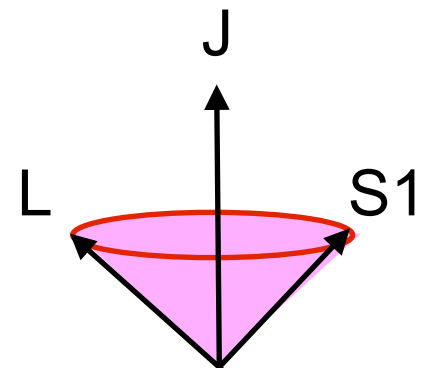
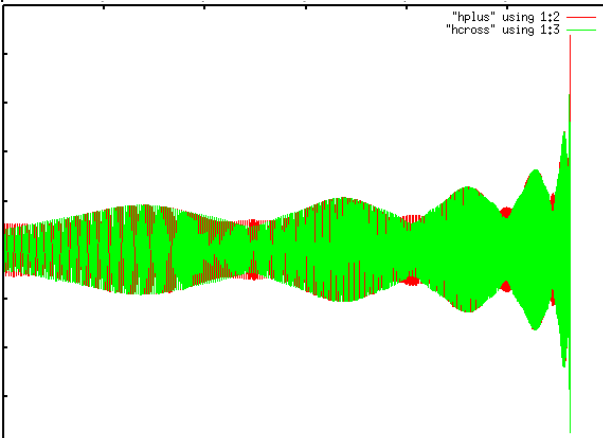
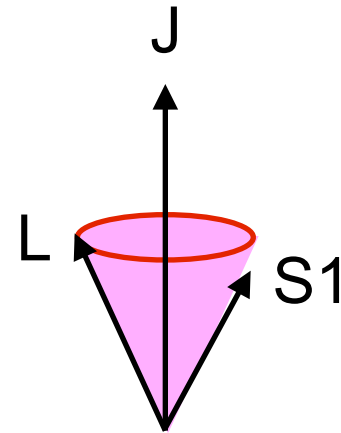
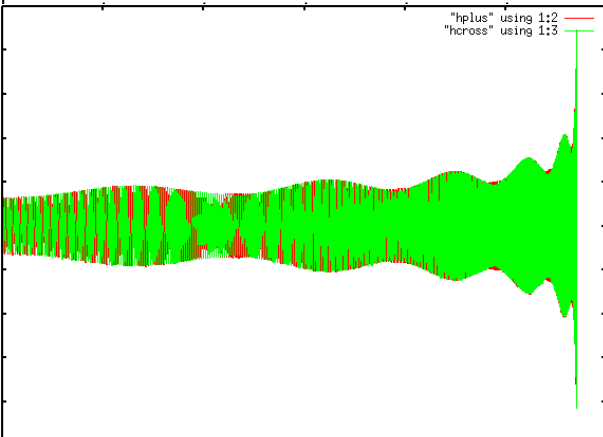
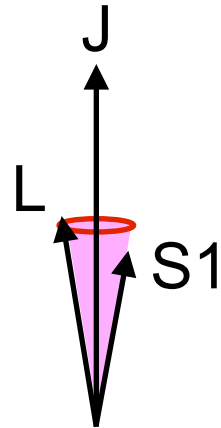
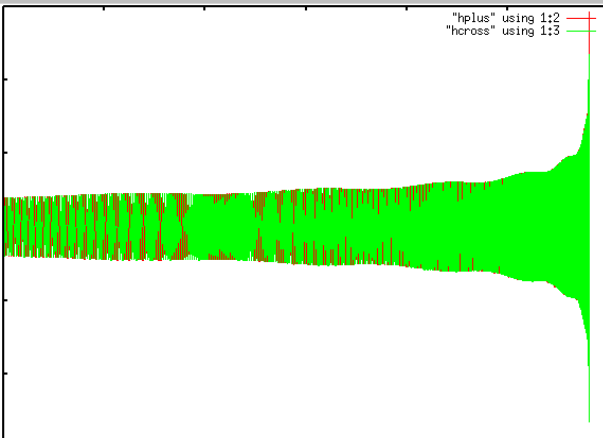
$$\dot{\hat{\mathbf{L}}}_N = -\frac{(M\omega)^{1/3}}{\eta M^2} \dot{\mathbf{S}} = \frac{\omega^2}{2M} \left\{ \left[\left(4 + 3\frac{m_2}{m_1}\right) \mathbf{S}_1 + \left(4 + 3\frac{m_1}{m_2}\right) \mathbf{S}_2 \right] \times \hat{\mathbf{L}}_N - \frac{3\omega^{1/3}}{\eta M^{5/3}} \left[(\mathbf{S}_2 \cdot \hat{\mathbf{L}}_N) \mathbf{S}_1 + (\mathbf{S}_1 \cdot \hat{\mathbf{L}}_N) \mathbf{S}_2 \right] \times \hat{\mathbf{L}}_N \right\}$$

aligned-spin: no precession

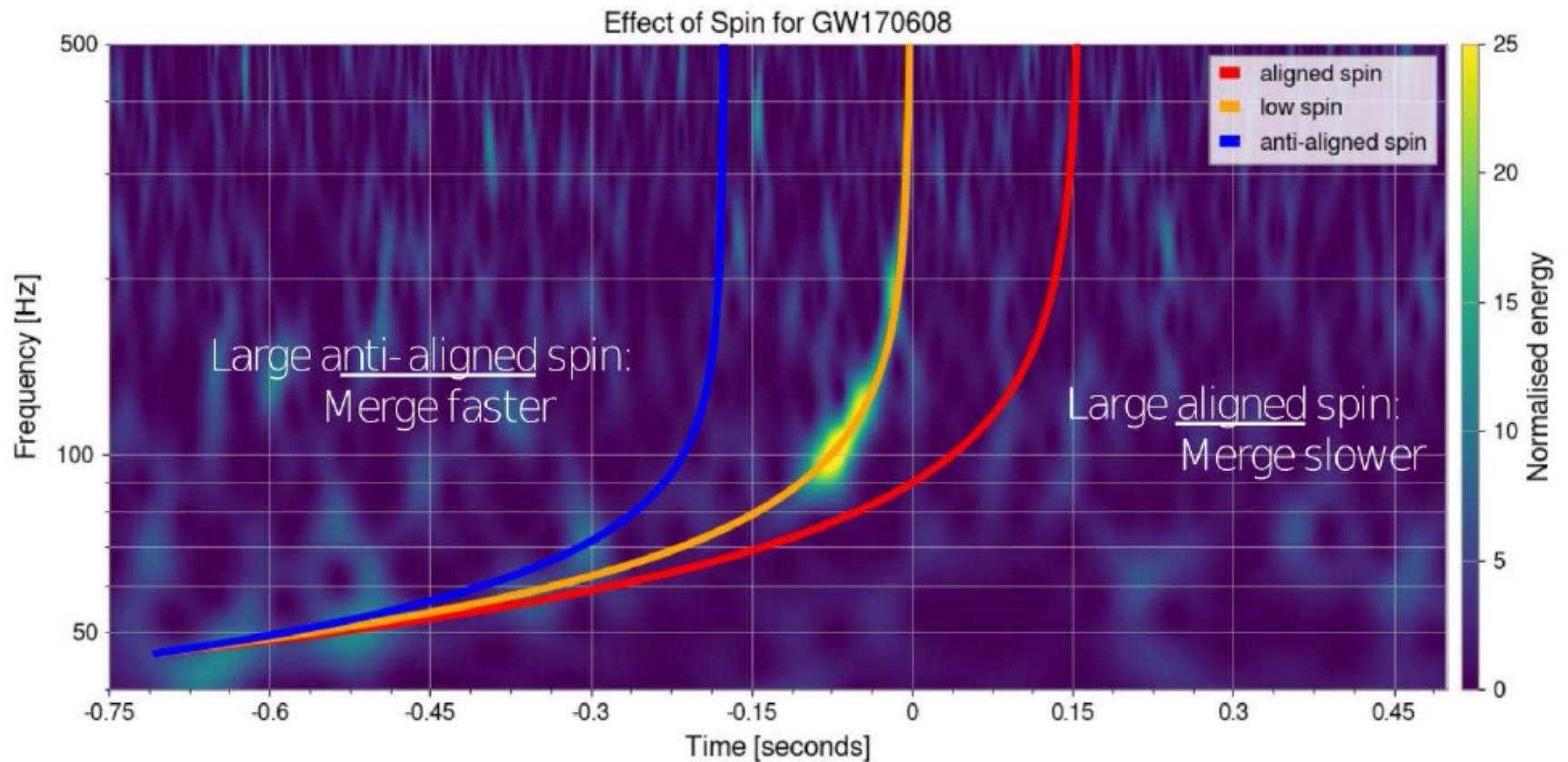
Amplitude modulation with **Spin** (Angle= $\pi/2$, $S2=0$)



Amplitude modulation with **Angle** ($S1=0.9$, $S2=0$)



aligned-spin binaries

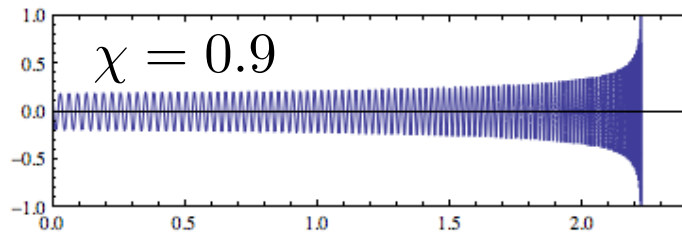
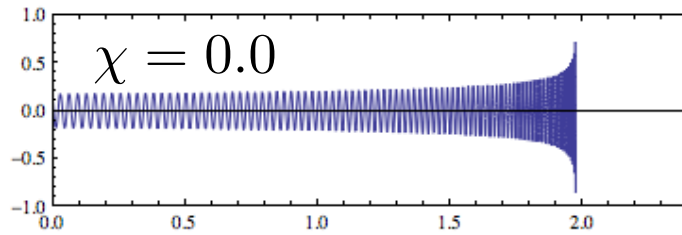
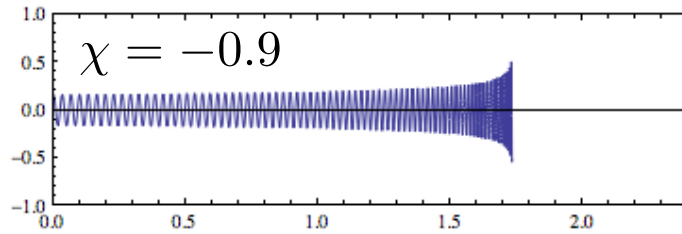


(see the pyCBC tutorial:

<http://pycbc.org/pycbc/latest/html/waveform.html#plotting-frequency-evolution-of-td-waveform>)

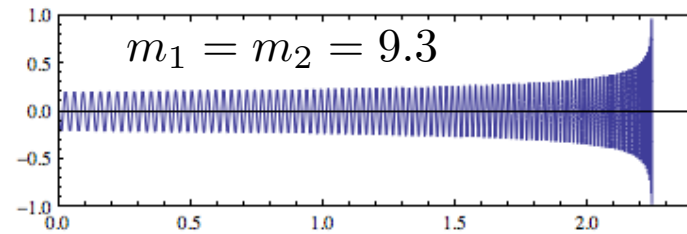
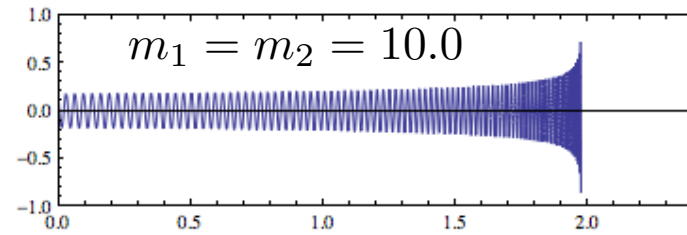
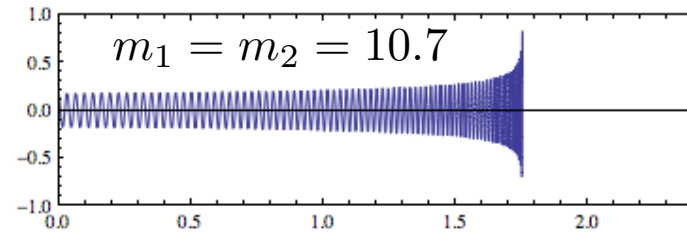
spin-mass degeneracy

$$m_1 = m_2 = 10M_\odot$$



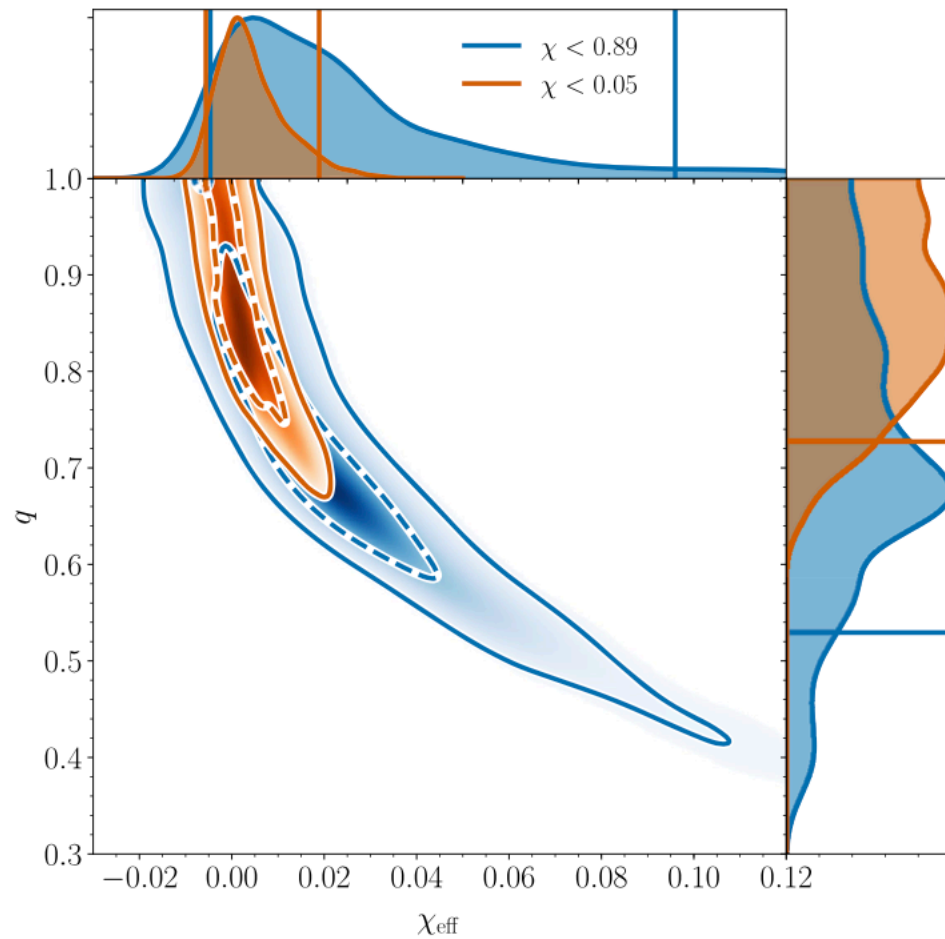
Time [s]

$$\chi = 0.0$$



Time [s]

GW170817



GW phase: nonspinning

$$h_{+} = -\frac{2G\mu}{c^2 R} x (1 + \cos^2 \Theta) \cos \underline{2\psi}$$

$$h_{\times} = -\frac{2G\mu}{c^2 R} x (2 \cos \Theta) \sin \underline{2\psi}$$

Phase evolution from
post Newtonian (PN)

Post Newtonian Energy & Flux

$$E = -\frac{\mu c^2 x}{2} \left\{ 1 + \left(-\frac{3}{4} - \frac{1}{12} \nu \right) x + \left(-\frac{27}{8} + \frac{19}{8} \nu - \frac{1}{24} \nu^2 \right) x^2 \right. \\ \left. + \left(-\frac{675}{64} + \left[\frac{34445}{576} - \frac{205}{96} \pi^2 \right] \nu - \frac{155}{96} \nu^2 - \frac{35}{5184} \nu^3 \right) x^3 \right\} \\ + \mathcal{O} \left(\frac{1}{c^8} \right).$$

Newtonian binding energy of a binary

$$\mathcal{L} = \frac{32c^5}{5G} \nu^2 x^5 \left\{ 1 + \left(-\frac{1247}{336} - \frac{35}{12} \nu \right) x + 4\pi x^{3/2} + \left(-\frac{44711}{9072} + \frac{9271}{504} \nu + \frac{65}{18} \nu^2 \right) x^2 \right. \\ + \left(-\frac{8191}{672} - \frac{583}{24} \nu \right) \pi x^{5/2} \\ + \left[\frac{6643739519}{69854400} + \frac{16}{3} \pi^2 - \frac{1712}{105} C - \frac{856}{105} \ln(16x) \right. \\ \left. + \left(-\frac{134543}{7776} + \frac{41}{48} \pi^2 \right) \nu - \frac{94403}{3024} \nu^2 - \frac{775}{324} \nu^3 \right] x^3 \\ \left. + \left(-\frac{16285}{504} + \frac{214745}{1728} \nu + \frac{193385}{3024} \nu^2 \right) \pi x^{7/2} + \mathcal{O} \left(\frac{1}{c^8} \right) \right\}.$$

$x \equiv \left(\frac{G m \omega}{c^3} \right)^{2/3}$	$m = m_1 + m_2$	$\mu = m_1 m_2 / m$	$\nu \equiv \frac{\mu}{m} \equiv \frac{m_1 m_2}{(m_1 + m_2)^2}.$
--	-----------------	---------------------	--

Phase evolution

Energy balance equation : orbital binding energy loss = GW emission energy

$$\frac{dE}{dt} = -L$$

$$\frac{dE}{dt} = \frac{dE}{dx} \frac{dx}{dt}$$

$$\frac{dx}{dt} = -\frac{L}{(dE/dx)}$$

expand with x , $x \equiv \left(\frac{G m \omega}{c^3} \right)^{2/3}$

$$\begin{aligned} \frac{\dot{\omega}}{\omega^2} = \frac{96}{5} \eta (M \omega)^{5/3} & \left(1 - \frac{743+924\eta}{336} (M \omega)^{2/3} + \left(\frac{34\,103}{18\,144} + \frac{13\,661}{2016} \eta + \frac{59}{18} \eta^2 \right) (M \omega)^{4/3} - \frac{1}{672} (4159 + 14\,532 \eta) \pi (M \omega)^{5/3} \eta \right. \\ & + \left[\left(\frac{16\,447\,322\,263}{139\,708\,800} - \frac{1712}{105} \gamma_E + \frac{16}{3} \pi^2 \right) + \left(-\frac{273\,811\,877}{1\,088\,640} + \frac{451}{48} \pi^2 - \frac{88}{3} \hat{\theta} \right) \eta + \frac{541}{896} \eta^2 - \frac{5605}{2592} \eta^3 \right. \\ & \left. \left. - \frac{856}{105} \log[16(M \omega)^{2/3}] \right] (M \omega)^2 + \left(-\frac{4\,415}{4\,032} + \frac{661\,775}{12\,096} \eta + \frac{149\,789}{3\,024} \eta^2 \right) \pi (M \omega)^{7/3} \right), \end{aligned}$$

$$w(t) = \int_0^t \dot{w}(w) dt, \quad \psi(t) = \int_0^t w(t) dt \quad \text{---> Numerical integration}$$

TaylorT4, T1,2,3,.. : time domain models

TaylorF2

$$h(f) = \frac{M_c^{5/6}}{\pi^{2/3} D_{\text{eff}}} \sqrt{\frac{5}{24}} \underline{f^{-7/6}} e^{i\Psi(f)},$$

stationary phase approx.

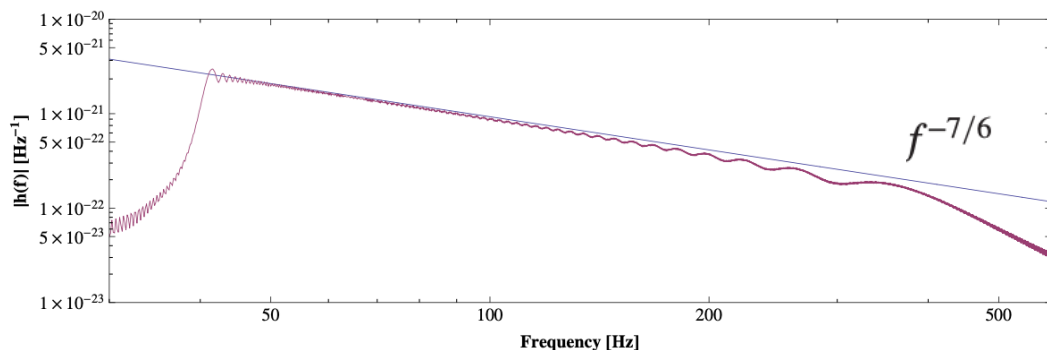


Figure 3.2: **Fourier domain TaylorT4 and SPA waveforms from a non-spinning binary.** We assume the same binary model as in figure 3.1. TaylorT4 (red) and SPA (blue) waveforms coincide at 40 Hz.

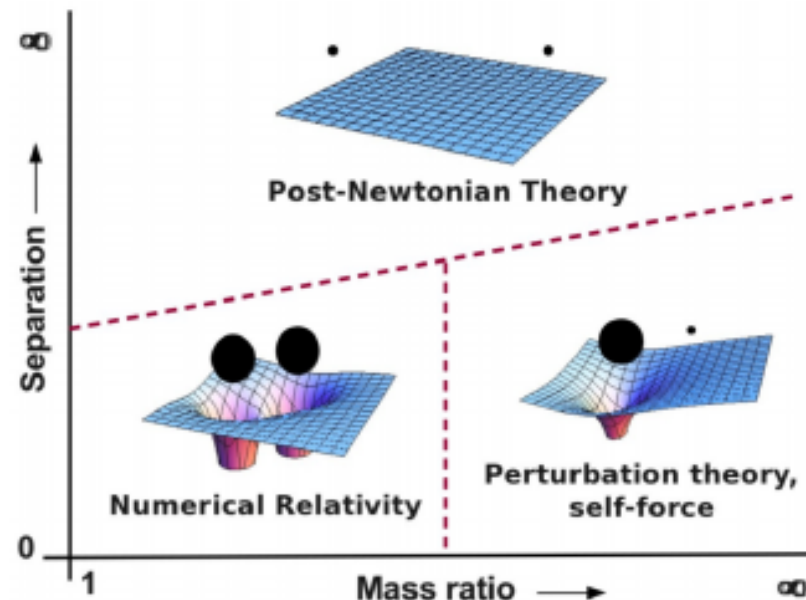
$$\begin{aligned} \Phi_{\text{SPA}}(f) = & 2\pi f t_c - \Phi_c - \frac{\pi}{4} + \frac{3}{128\nu v^5} \left\{ 1 + \frac{20}{9} \left(\frac{743}{336} + \frac{11}{4} \nu \right) v^2 - 16\pi v^3 \right. \\ & + 10 \left(\frac{3058673}{1016064} + \frac{5429}{1008} \nu + \frac{617}{144} \nu^2 \right) v^4 \\ & + \left(\frac{38645}{756} - \frac{65}{9} \nu \right) \left[1 + 3 \log \left(\frac{v}{v_{\text{iso}}} \right) \right] \pi v^5 \left[\left(\frac{11583231236531}{4694215680} \right. \right. \\ & - \frac{640}{3} \pi^2 - \frac{6848}{21} \gamma_E - \frac{6848}{21} \log(4v) \Big) + \left(\frac{2255}{12} \pi^2 - \frac{15737765635}{3048192} \right) \nu \\ & \left. \left. + \frac{76055}{1728} \nu^2 - \frac{127825}{1296} \nu^3 \right] v^6 + \left(\frac{77096675}{254016} + \frac{378515}{1512} \nu - \frac{74045}{756} \nu^2 \right) \pi v^7 \right\} \end{aligned} \quad (3.20)$$

$$v \equiv [\pi f (m_1 + m_2)]^{1/3}$$

Waveform modeling

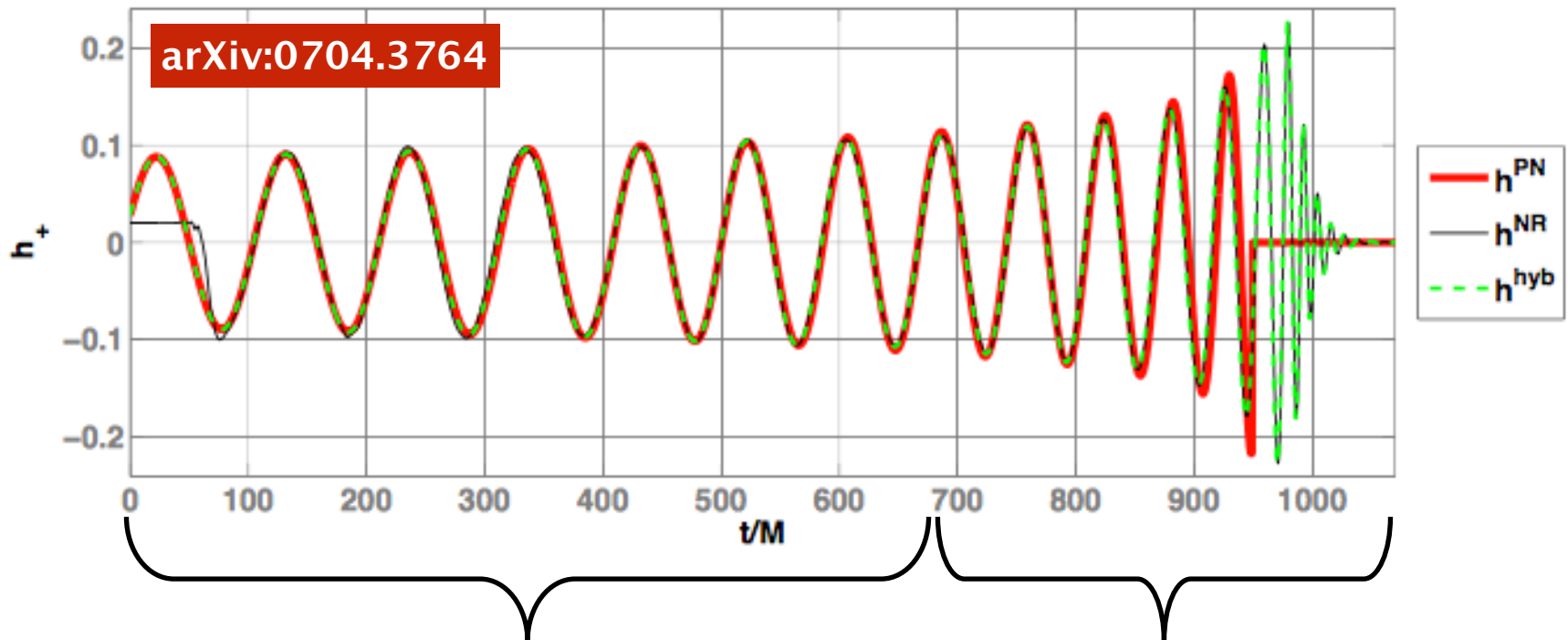
- **Post-Newtonian (PN) approximation**: slow motion approximation ($v/c \ll 1$)
 - **Perturbative theory**: small mass ratio ($m/M \ll 1$)
 - **Effective-One-Body approach**: combination of PN, Perturbative approach and NR
- time-domain model !!

Binary parameter space



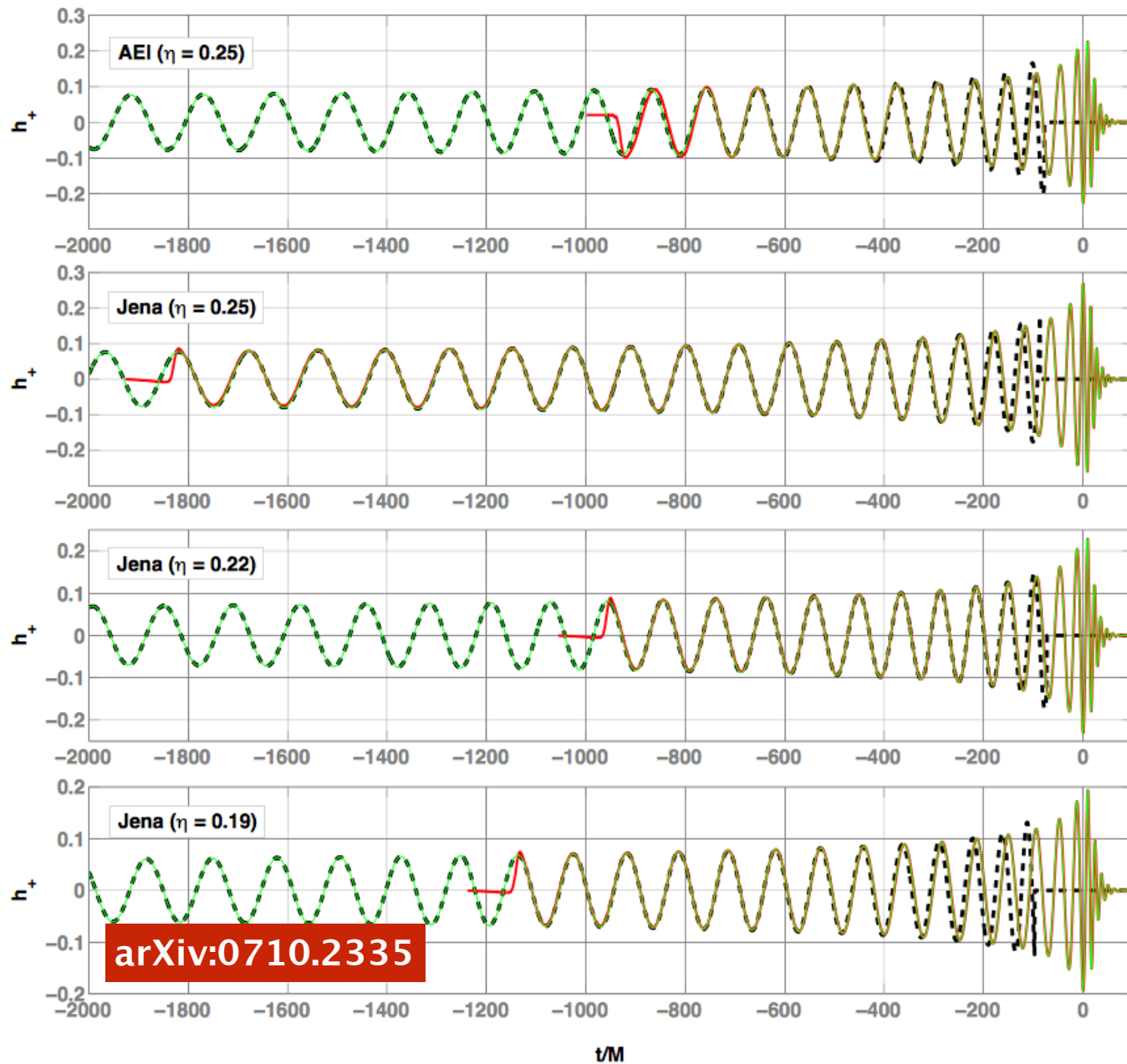
Phenomenological model : IMRPhenom

frequency-domain model !!

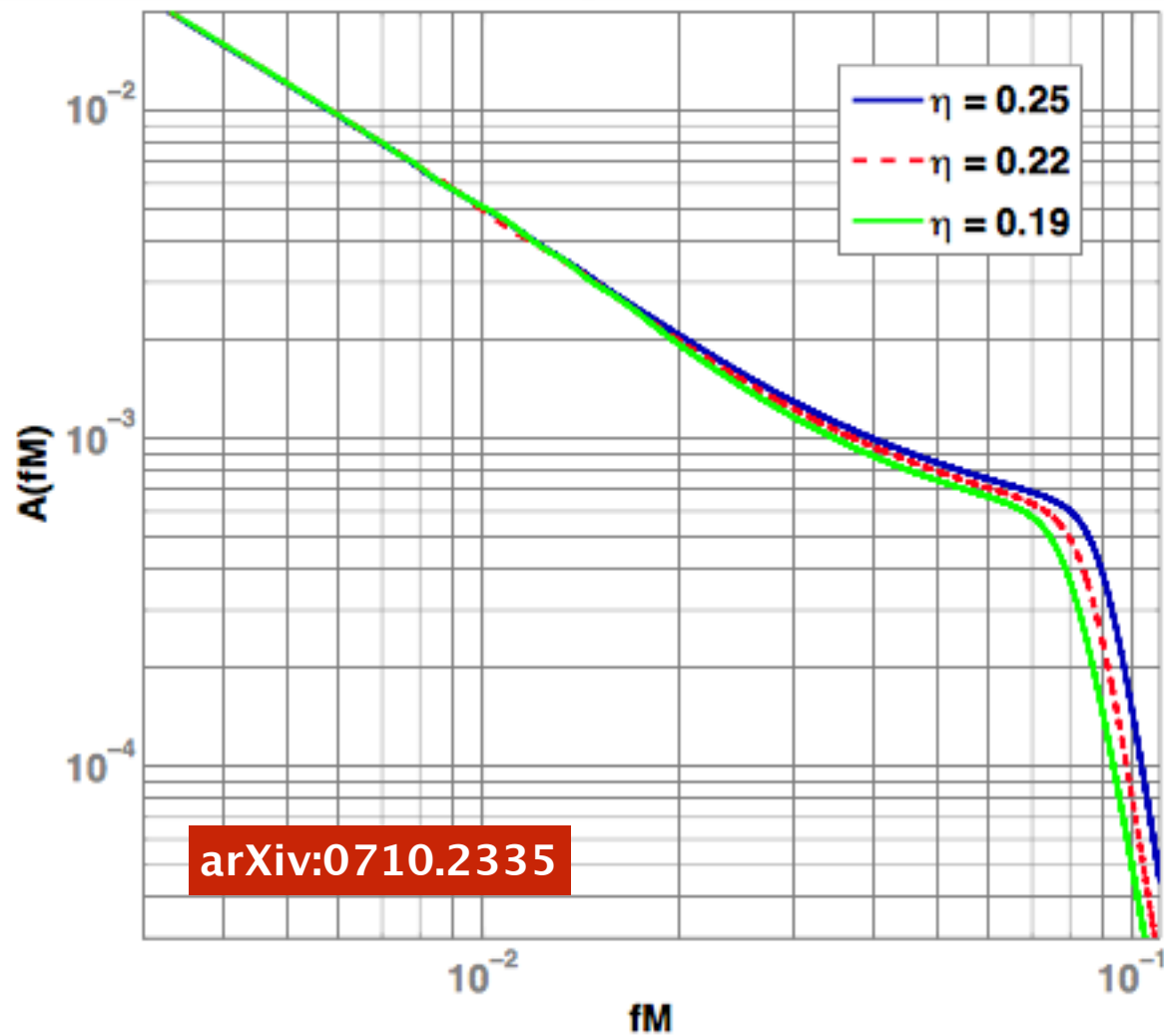


Hybrid = early inspiral:PN + late inspiral-M-R:NR

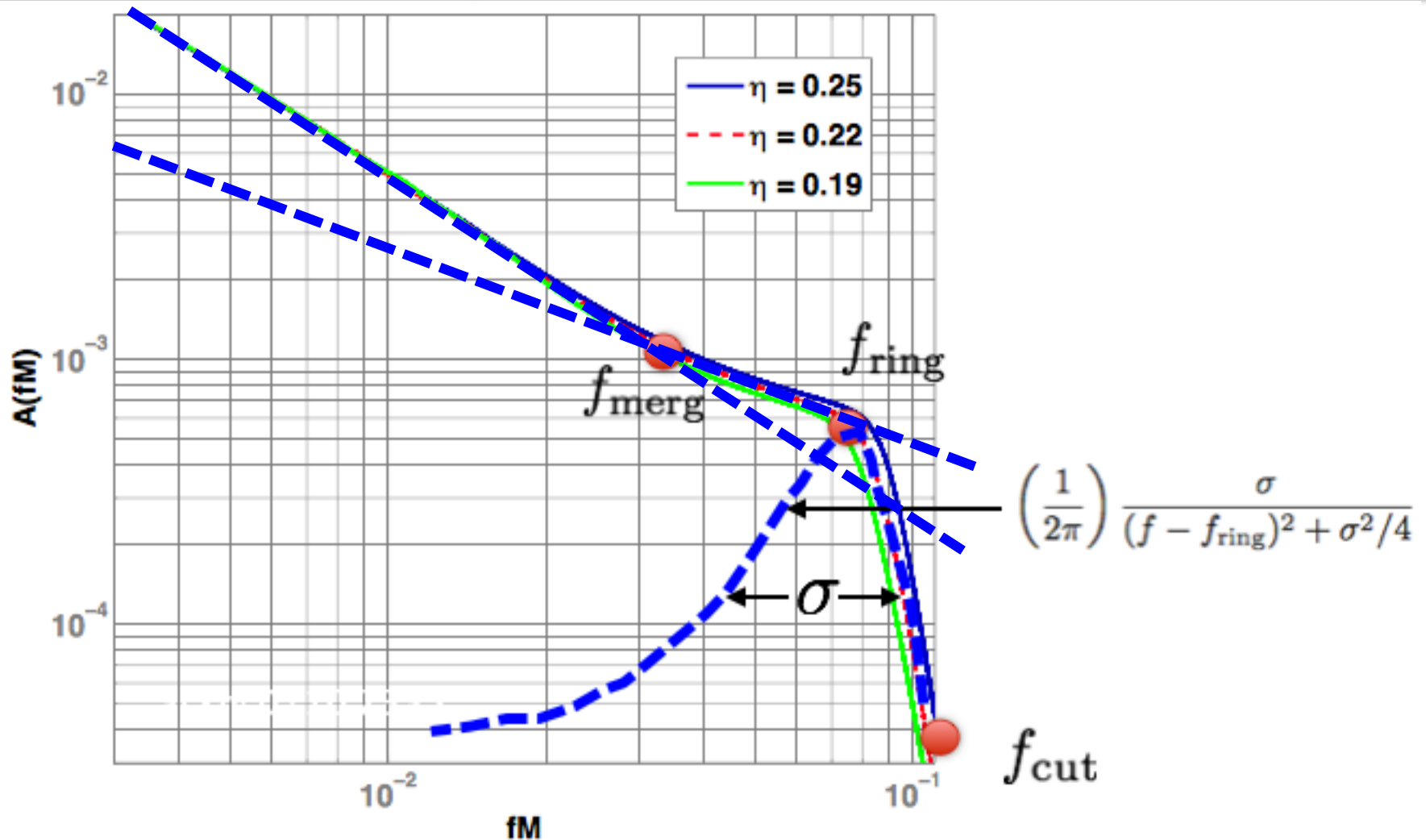
Constructing Hybrid waveforms



Fourier amplitudes of hybrid waveforms



Fourier amplitudes of hybrid waveforms

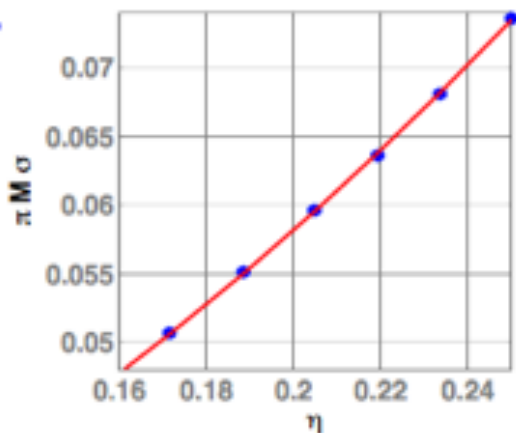
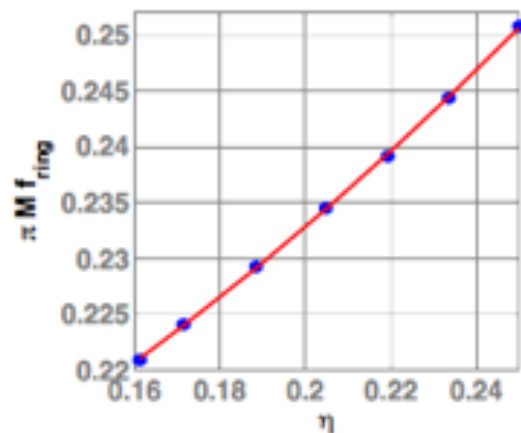
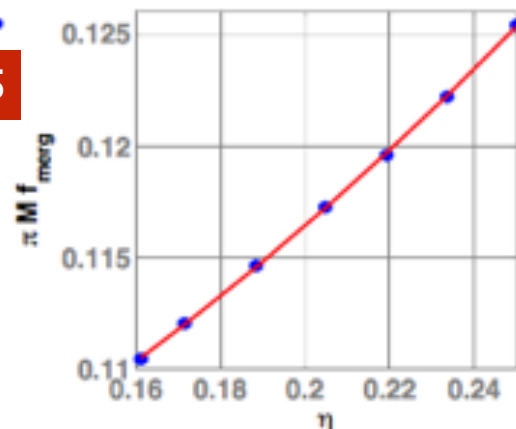
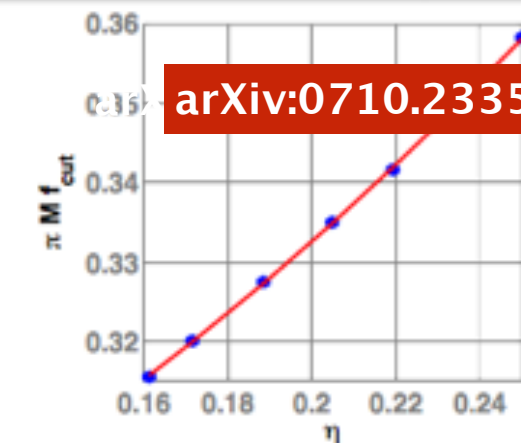


$$A_{\text{eff}}(f) \equiv C \begin{cases} (f/f_{\text{merg}})^{-7/6} & \text{if } f < f_{\text{merg}} \\ (f/f_{\text{merg}})^{-2/3} & \text{if } f_{\text{merg}} \leq f < f_{\text{ring}} \\ w \mathcal{L}(f, f_{\text{ring}}, \sigma) & \text{if } f_{\text{ring}} \leq f < f_{\text{cut}} \end{cases}$$

4 parameters

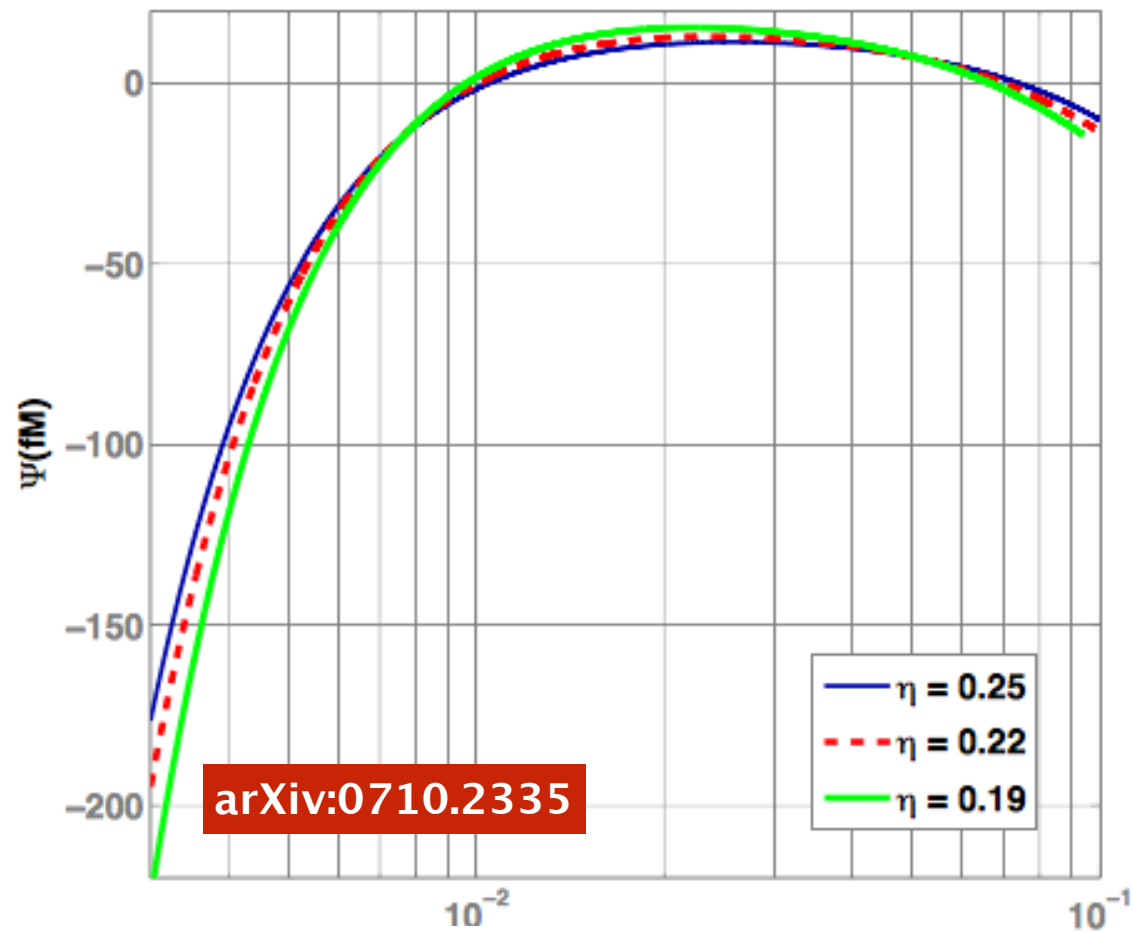
Best-match amplitude parameters

arXiv:0710.2335



$$\alpha_{j \text{ int}} = \frac{a_j \eta^2 + b_j \eta + c_j}{\pi M}$$

Parameter	a_k	b_k	c_k
f_{merg}	2.9740×10^{-1}	4.4810×10^{-2}	9.5560×10^{-2}
f_{ring}	5.9411×10^{-1}	8.9794×10^{-2}	1.9111×10^{-1}
σ	5.0801×10^{-1}	7.7515×10^{-2}	2.2369×10^{-2}
f_{cut}	8.4845×10^{-1}	1.2848×10^{-1}	2.7299×10^{-1}

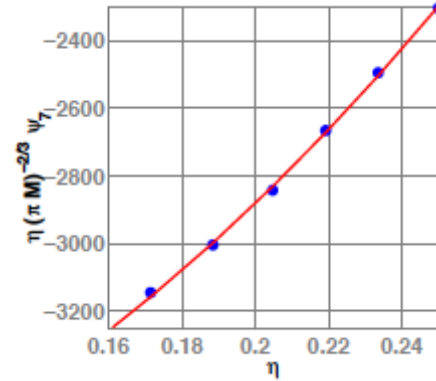
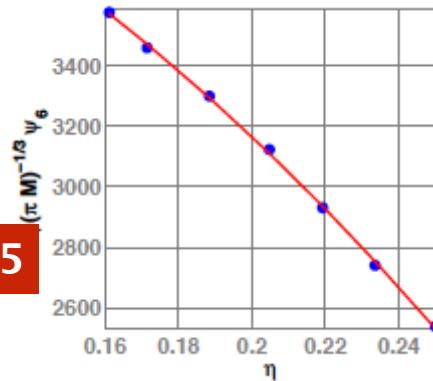
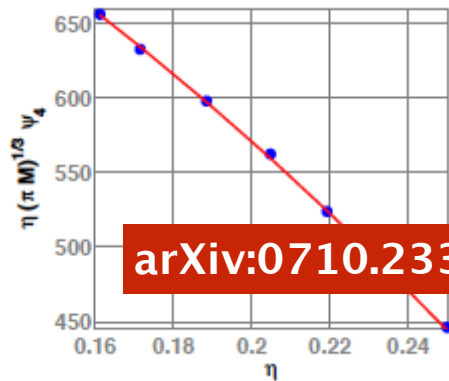
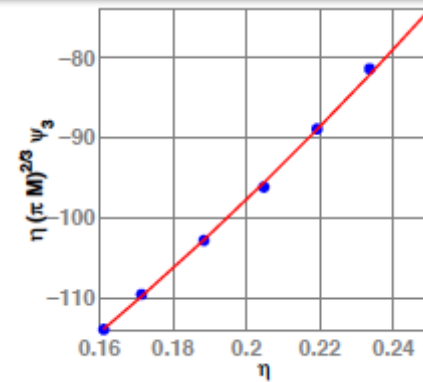
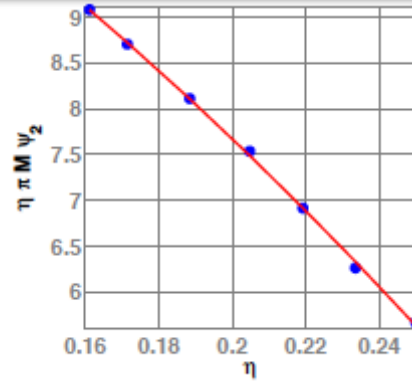
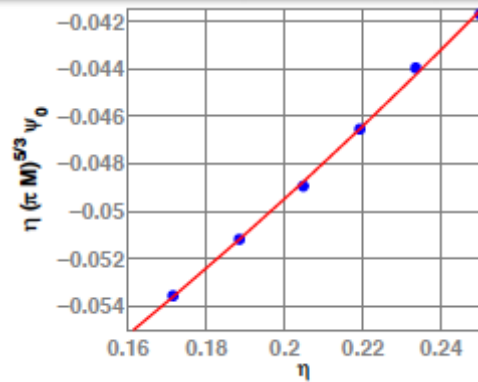


$$\Psi_{\text{eff}}(f) = 2\pi f t_0 + \varphi_0 + \sum_{k=0}^7 \psi_k f^{(k-5)/3},$$

$$\psi = \{\psi_0, \psi_2, \psi_3, \psi_4, \psi_6, \psi_7\}$$

6 parameters

Best-match phase parameters



arXiv:0710.2335

$$\psi_{k \text{ int}} = \frac{x_k \eta^2 + y_k \eta + z_k}{\eta (\pi M)^{(5-k)/3}}$$

Parameter	x_k	y_k	z_k
ψ_0	1.7516×10^{-1}	7.9483×10^{-2}	-7.2390×10^{-2}
ψ_2	-5.1571×10^1	-1.7595×10^1	1.3253×10^1
ψ_3	6.5866×10^2	1.7803×10^2	-1.5972×10^2
ψ_4	-3.9031×10^3	-7.7493×10^2	8.8195×10^2
ψ_6	-2.4874×10^4	-1.4892×10^3	4.4588×10^3
ψ_7	2.5196×10^4	3.3970×10^2	-3.9573×10^3

Early inspiral amplitude : TaylorF2

TaylorF2

$$h(f) = \frac{M_c^{5/6}}{\pi^{2/3} D_{\text{eff}}} \sqrt{\frac{5}{24}} \underline{f^{-7/6}} e^{i\Psi(f)},$$

IMRPhenom

$$u(f) \equiv \mathcal{A}_{\text{eff}}(f) e^{i\Psi_{\text{eff}}(f)}, \quad \mathcal{A}_{\text{eff}}(f) \equiv \begin{cases} (f/f_{\text{merg}})^{-7/6} & \text{if } f < f_{\text{merg}} \\ \underline{(f/f_{\text{merg}})^{-2/3}} & \text{if } f_{\text{merg}} \leq f < f_{\text{ring}} \\ w\mathcal{L}(f, f_{\text{ring}}, \sigma) & \text{if } f_{\text{ring}} \leq f < f_{\text{cut}}. \end{cases}$$

$$\Psi_{\text{eff}}(f) = 2\pi f t_0 + \phi_0 + \psi_0 f^{-5/3} + \psi_2 f^{-1} + \psi_3 f^{-2/3} + \psi_4 f^{-1/3} + \psi_6 f^{1/3}$$

Model	PhenomA	PhenomB	PhenomC
Mass range [$M_{\odot}$]	$50 \leq M \leq 200$	$M \leq 400$	$M \leq 350$
Mass ratio range	$q \leq 4$	$q \leq 10$	$q \leq 4$
Detector	initial LIGO, Virgo, advanced LIGO	initial LIGO	advanced LIGO

BBH Models

Family	Short name	Full name	<u>Precession</u>	Multipoles ($\ell, m $)	Reference
EOBNR	EOBNR	SEOBNRv4_ROM	$\times$	(2, 2)	[57]
	EOBNR HM	SEOBNRv4HM_ROM	$\times$	(2, 2), (2, 1), (3, 3), (4, 4), (5, 5)	[26,32]
	EOBNR P	SEOBNRv4P	$\checkmark$	(2, 2), (2, 1)	[33,118,119]
	EOBNR PHM	SEOBNRv4PHM	$\checkmark$	(2, 2), (2, 1), (3, 3), (4, 4), (5, 5)	[33,118,119]
Phenom	Phenom	IMRPhenomD	$\times$	(2, 2)	[120,121]
	Phenom HM	IMRPhenomHM	$\times$	(2, 2), (2, 1), (3, 3), (3, 2), (4, 4), (4, 3)	[22]
	Phenom P	IMRPhenomPv2/v3 ^a	$\checkmark$	(2, 2)	[23,122]
	Phenom PHM	IMRPhenomPv3HM	$\checkmark$	(2, 2), (2, 1), (3, 3), (3, 2), (4, 4), (4, 3)	[24]

BH-NS Models

Full name (implemented in LAL)	References	
Short label (used in this work)	Base model	Corrections
SEOBNRv4_ROM_NRTidalv2	[14–16]	[17, 18]
SEOBNR_T		
SEOBNRv4_ROM_NRTidalv2_NSBH [19]	[14–16]	[17, 18, 20]
SEOBNR_NSBH		
IMRPhenomPv2_NRTidalv2	[21–23]	[17, 18]
IMRPhenomP_T		
IMRPhenomNSBH [24]	[25]	[18, 20]
IMRPhenom_NSBH		

TABLE I: Waveform models used in our analysis for NSBH systems.

Thanks !