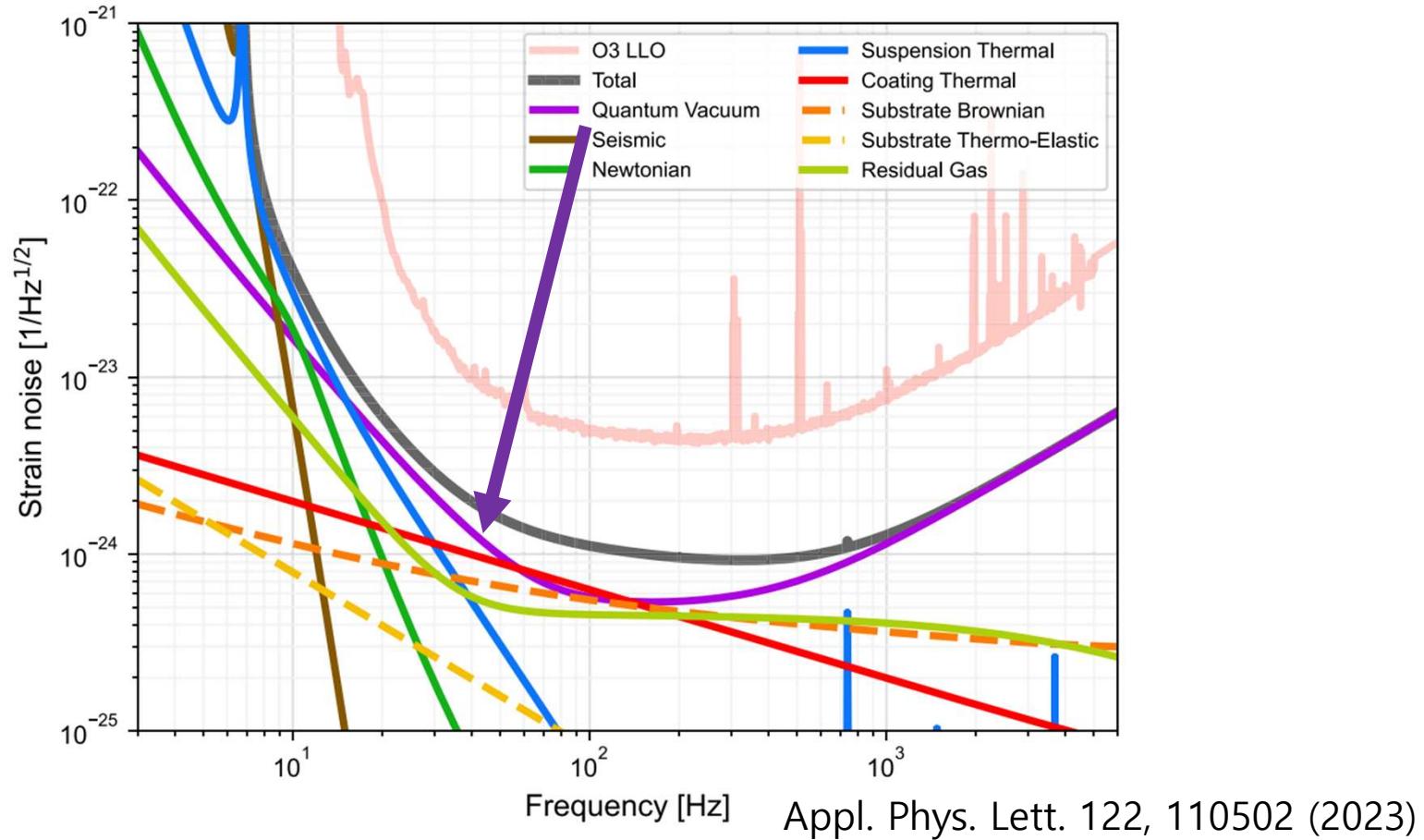
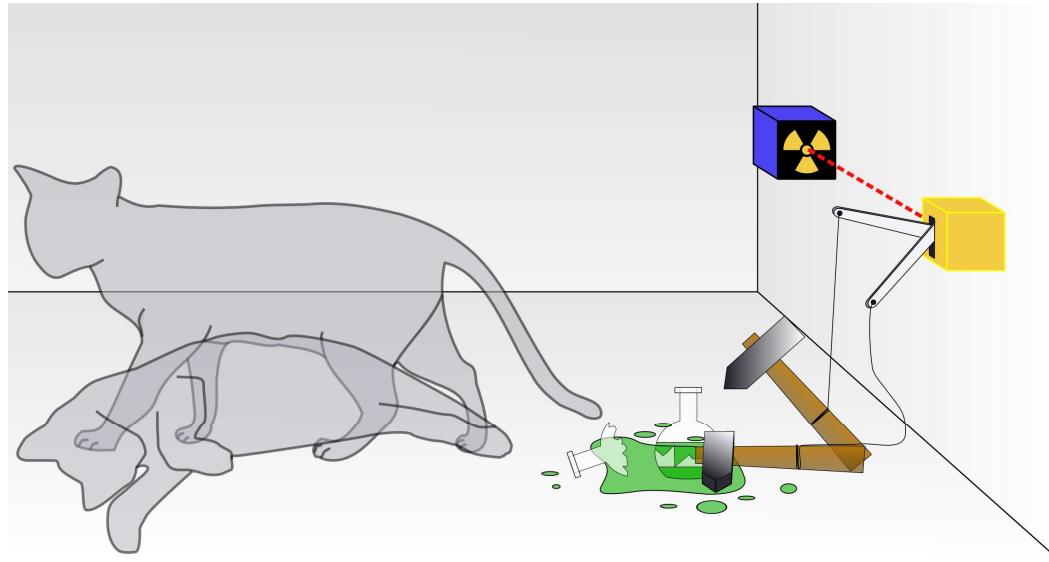


3. Quantum noise of gravitational wave detector

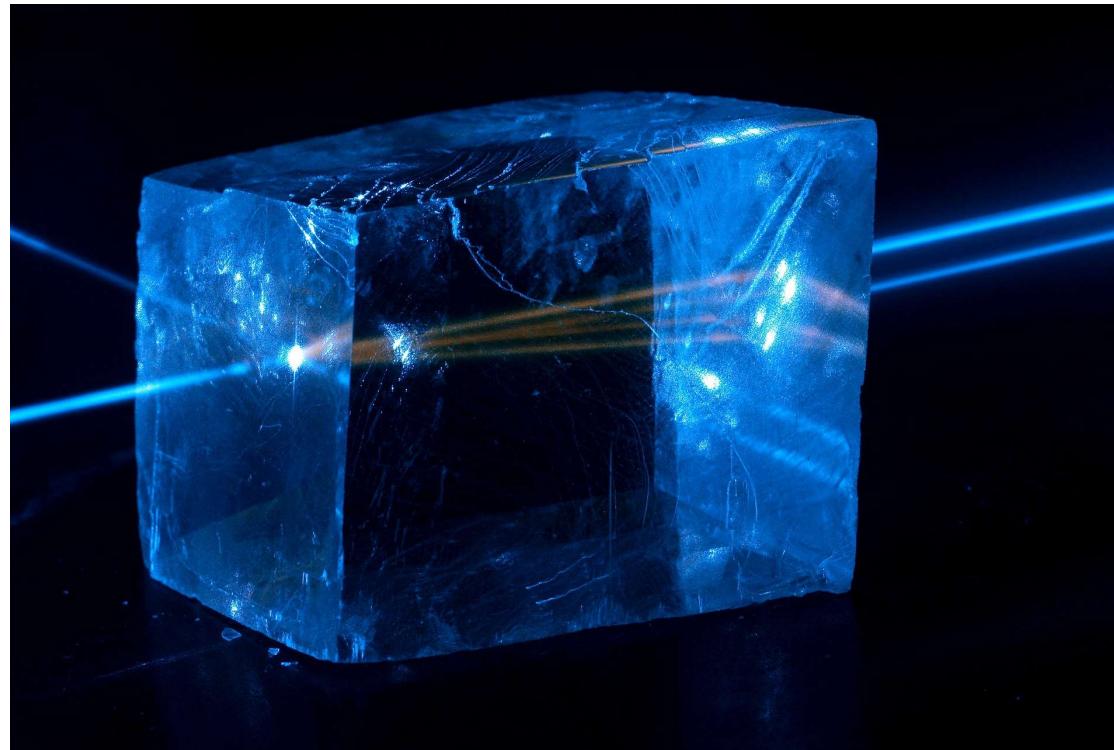
LIGO sensitivity



Quantum noise



Wave-particle duality



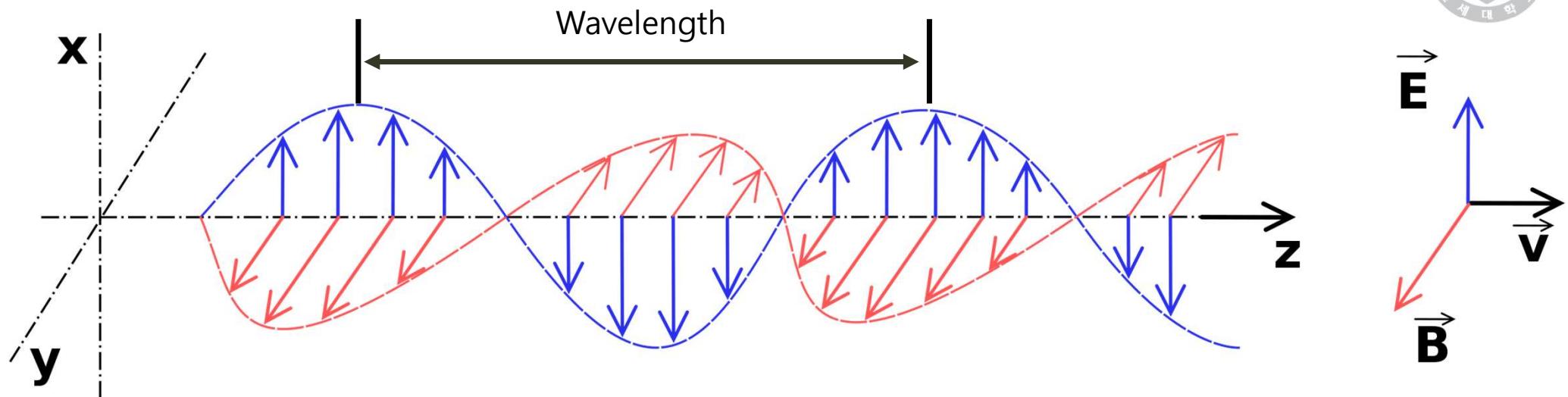
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Wave-particle duality

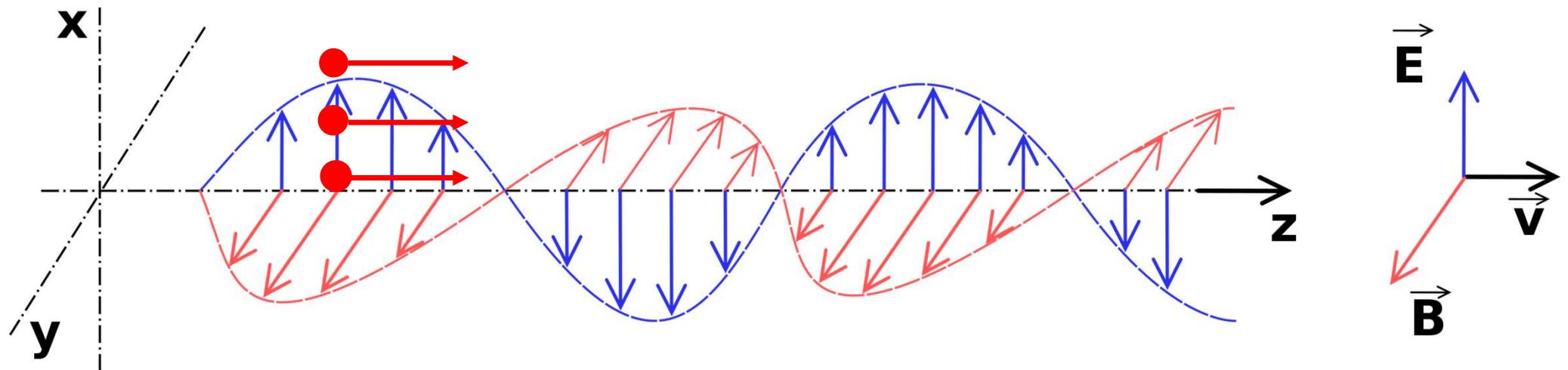


Light(Electromagnetic wave)

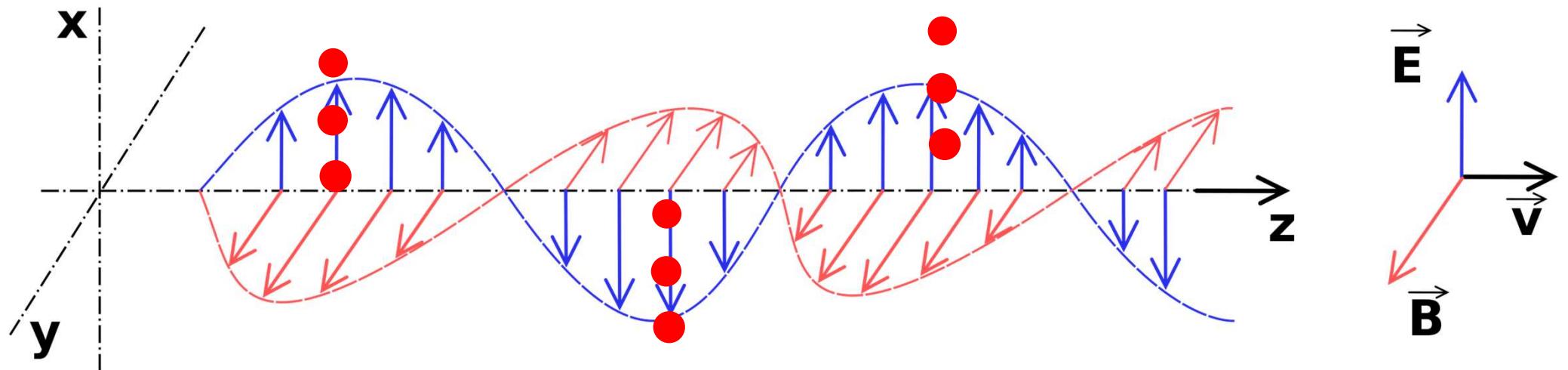


$$|\vec{E}| \propto \text{Intensity of Light}$$

Wave-particle duality



Wave-particle duality



Observed Photon

Photoelectric effect

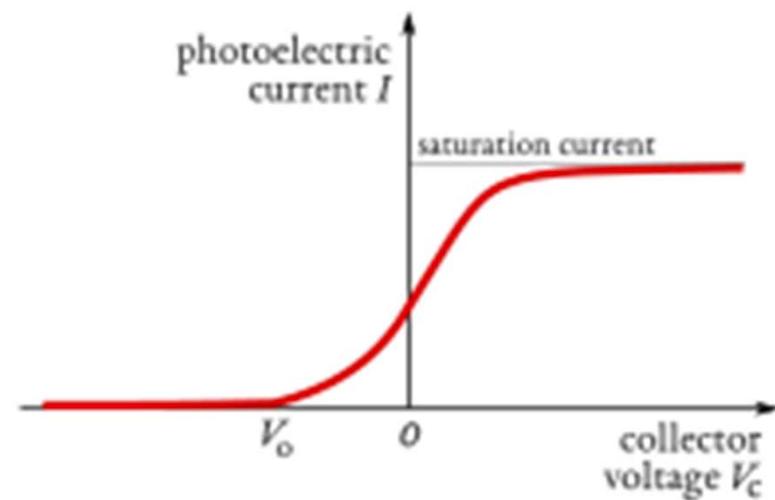
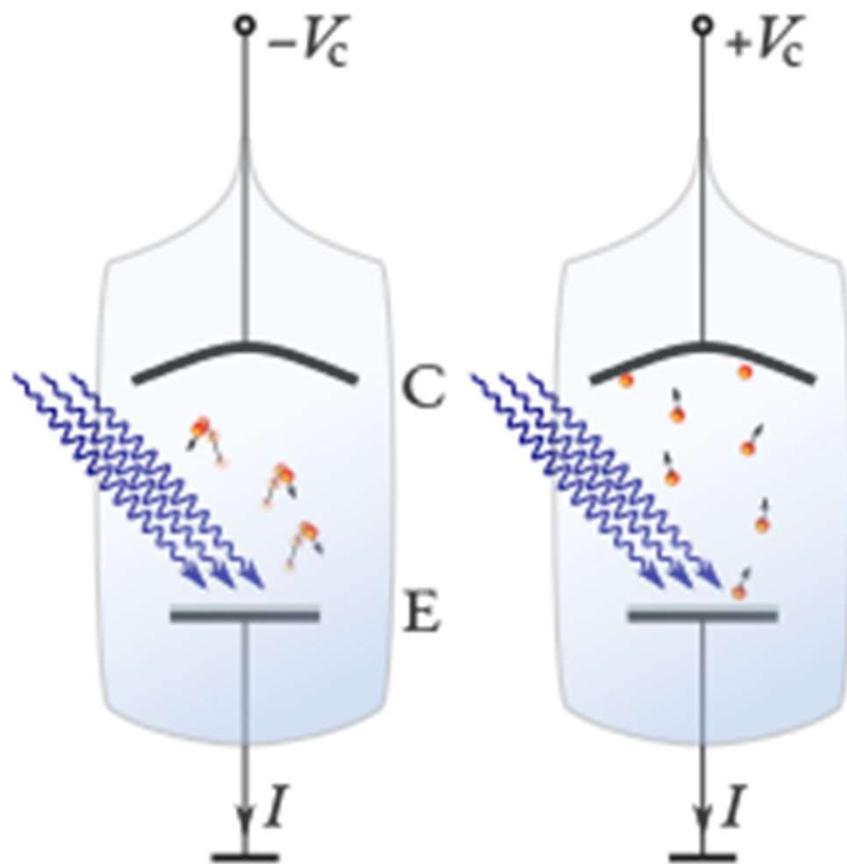


Photo detector

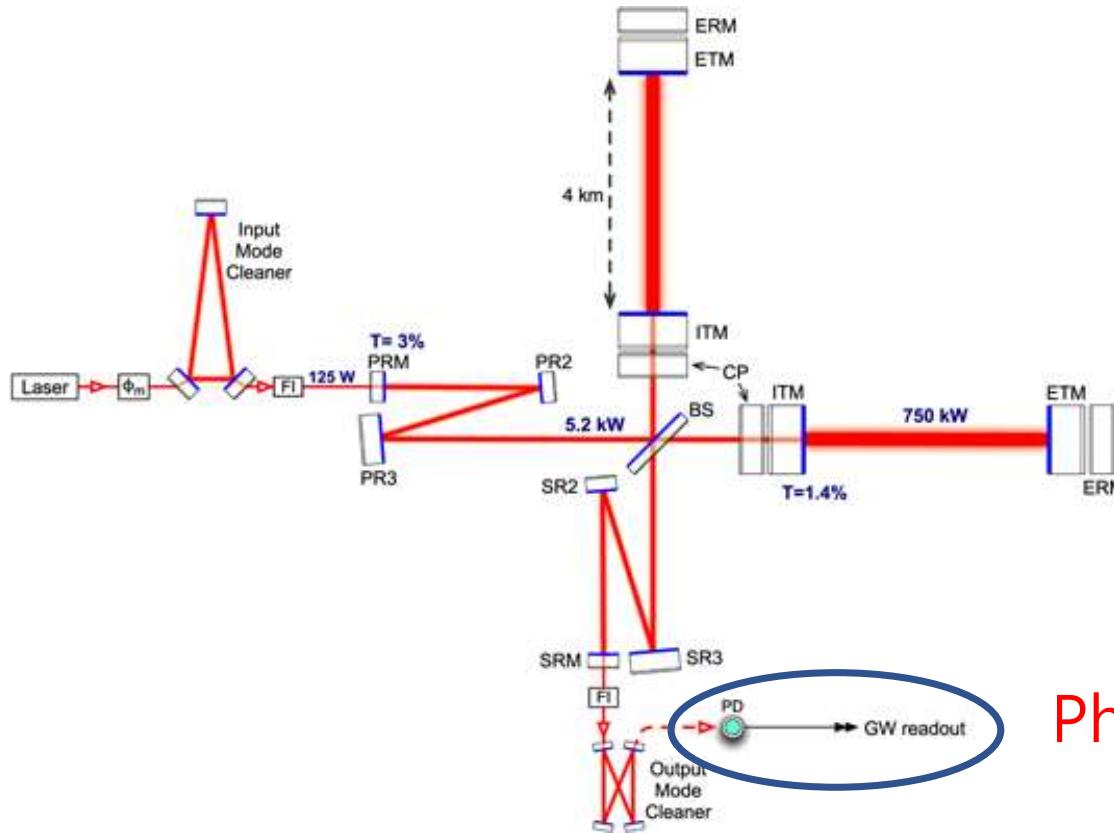


Amplitude of electricfield
 \propto Number of photon

Measuring the intensity of light
= Counting the number of photon

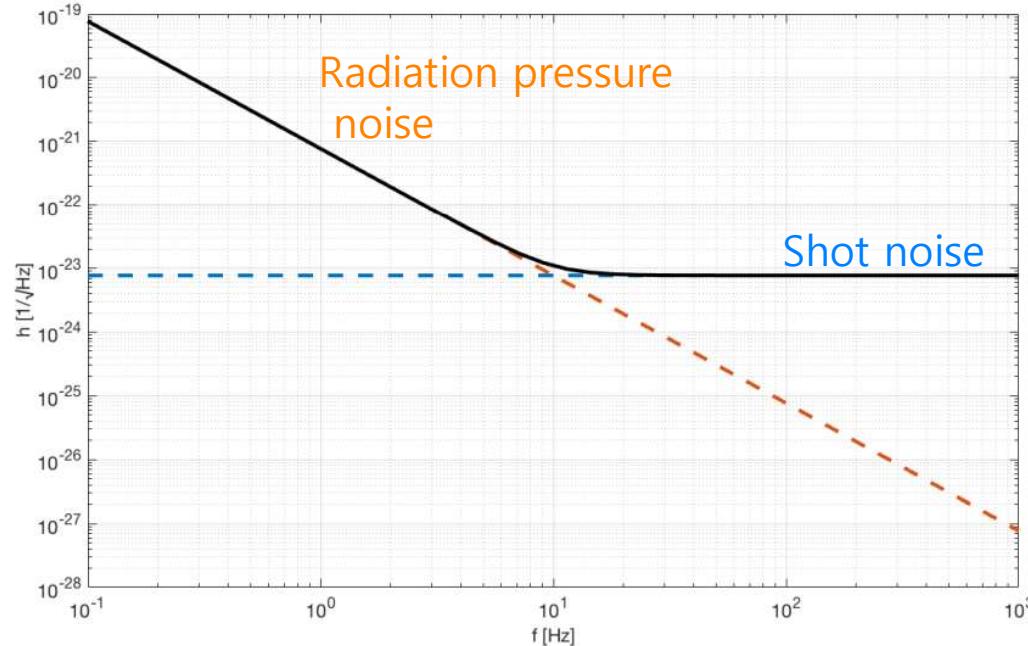


Gravitational wave detector



Photon counting

Standard quantum limit of GW detector



Standard quantum limit of gravitational wave detector
Shot noise + Radiation pressure noise

Shot noise



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Shot noise



Shot noise of interferometer

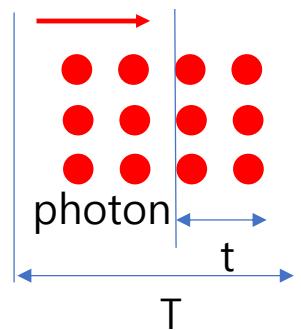


$$E = h\nu$$

Photon energy

Laser power(W) = the number of photon(n) / time(s)

$$nh\nu = \text{ / s}$$



Shot noise

Shot noise of interferometer

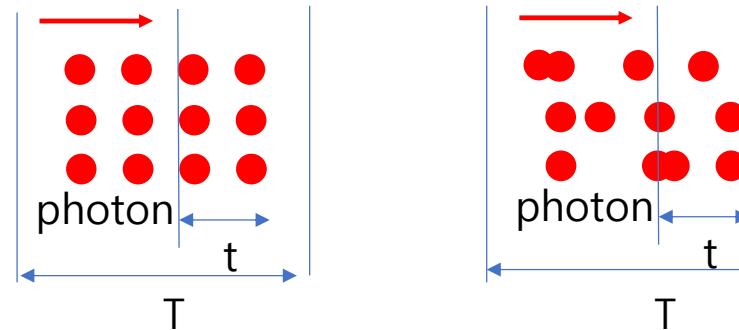


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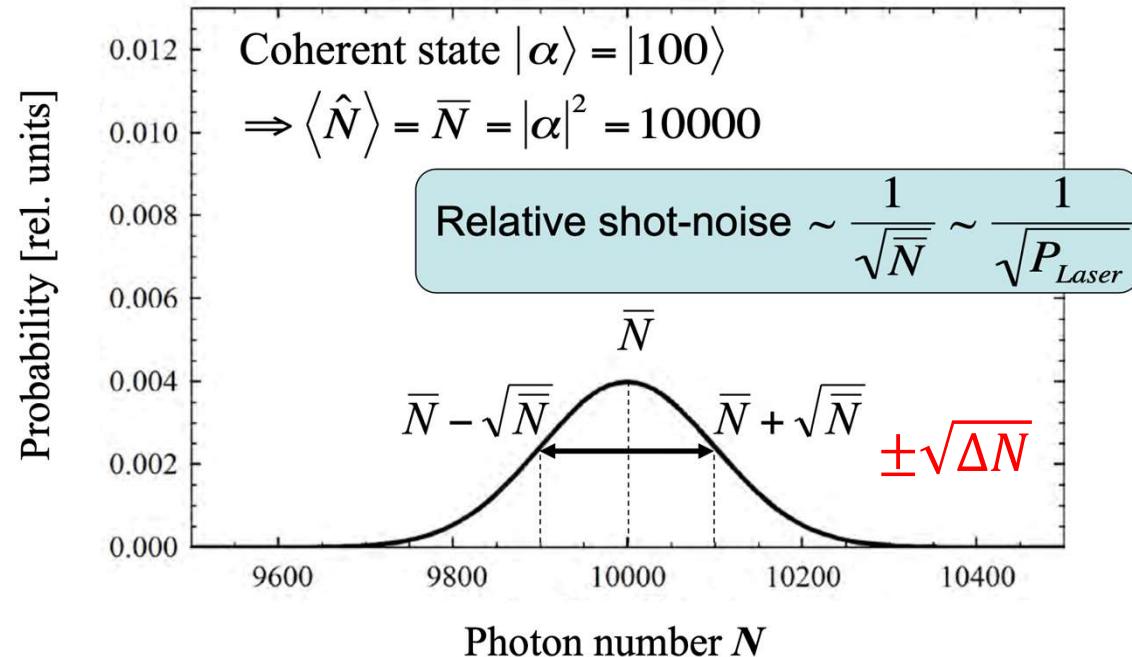
Shot noise

Shot noise of interferometer



Photon Counting Statistics

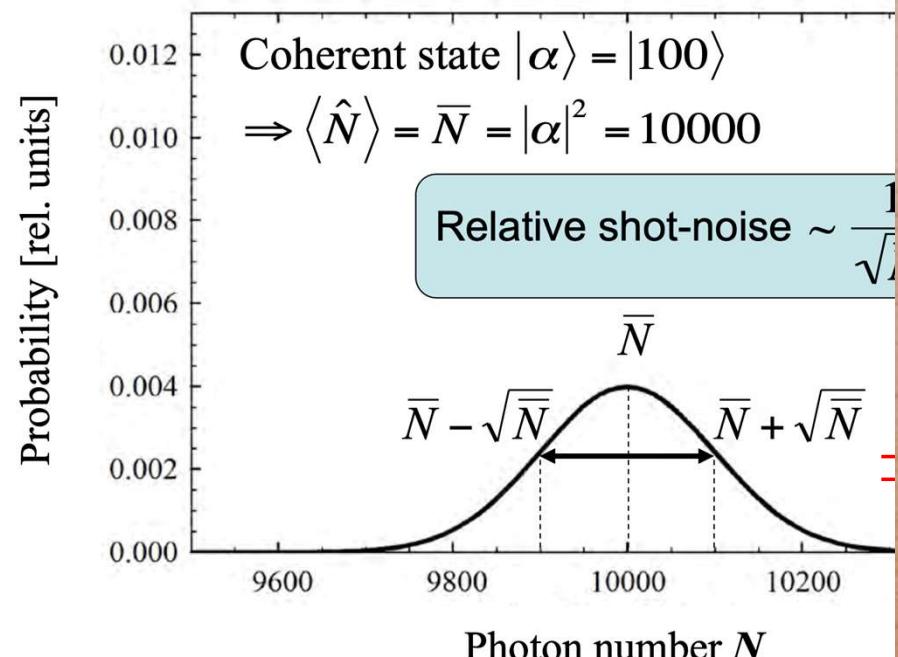
$$h = 6.62607015 \times 10^{-34} \text{ J}\cdot\text{Hz}^{-1}$$



Shot noise of interferometer



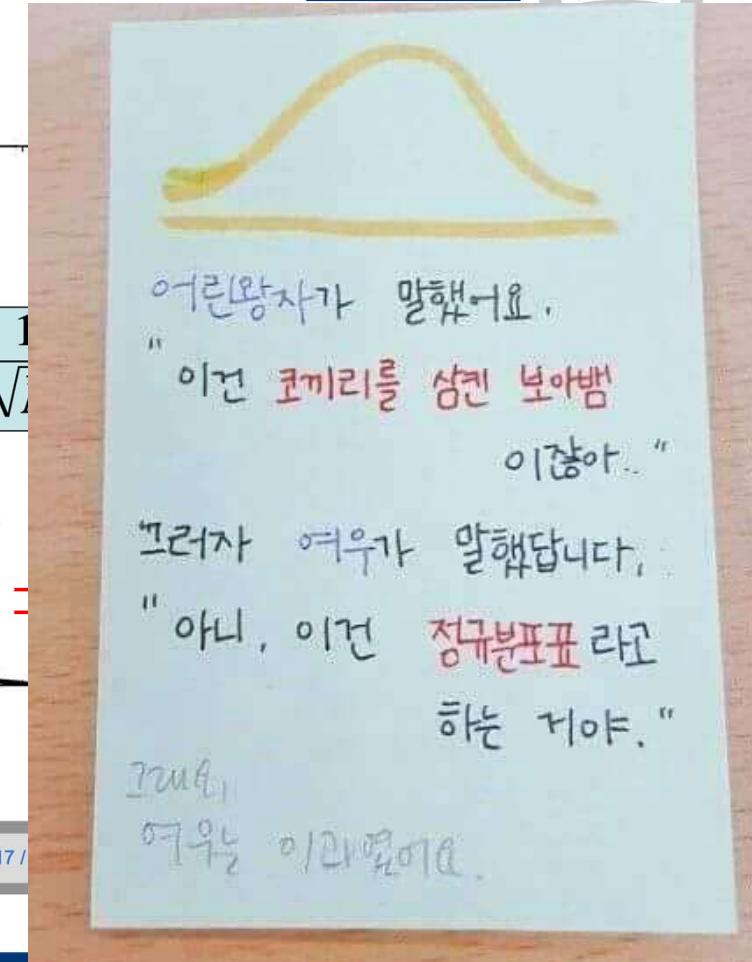
Photon Counting Statistics



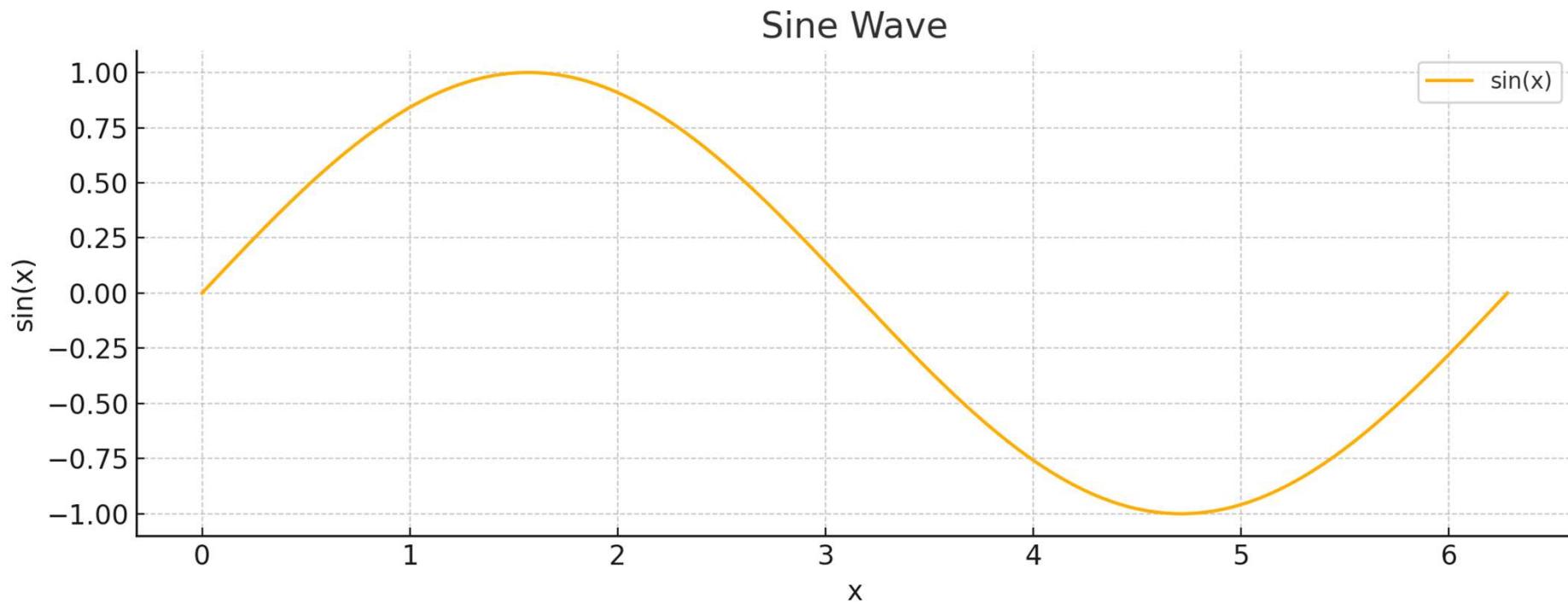
Albert-Einstein-Institut



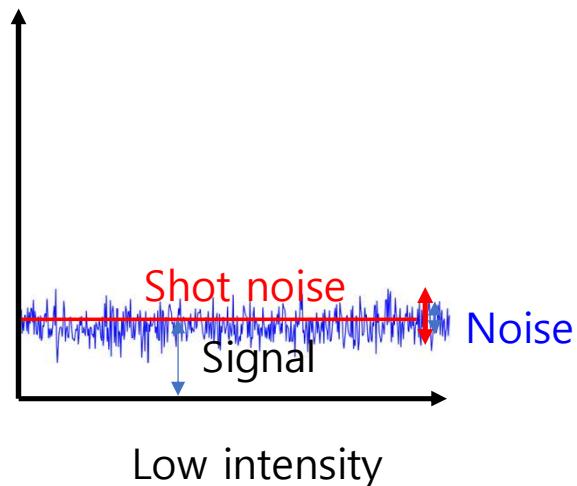
Roman Schnabel, 17 /



Sine wave

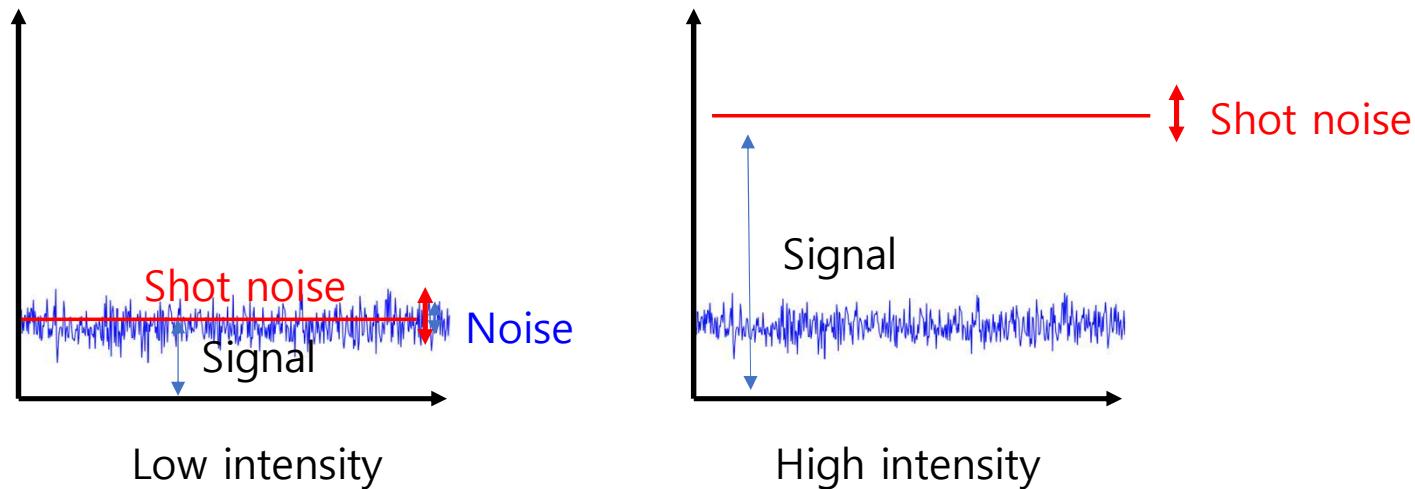


Shot noise of interferometer



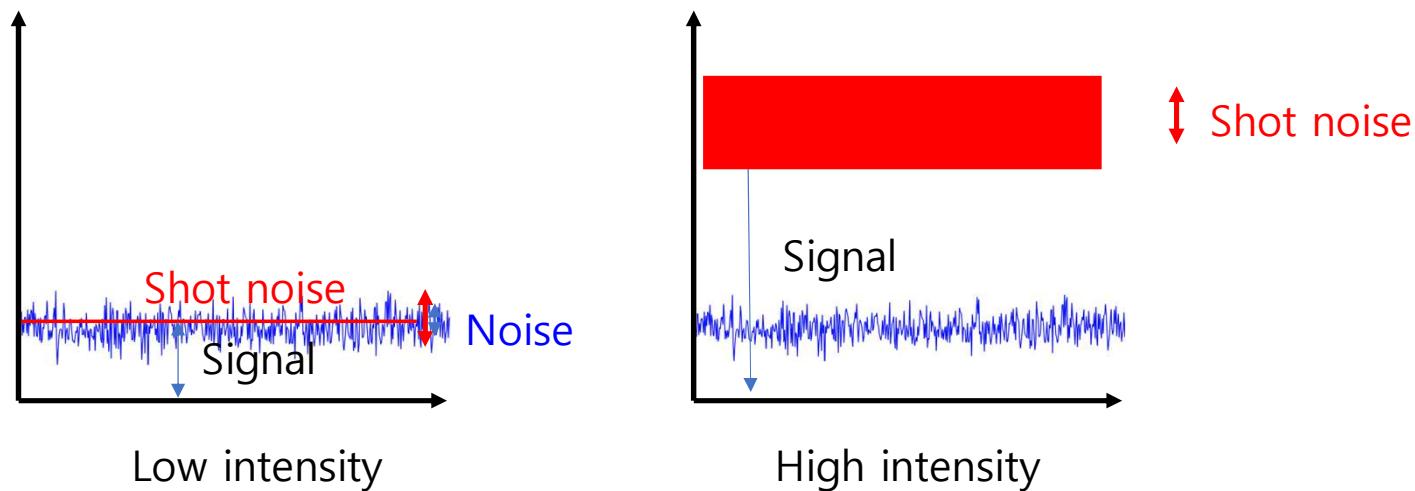
$$\text{Signal to Noise ratio} = \frac{N}{\sqrt{N}}$$

Shot noise of interferometer



$$\text{Signal to Noise ratio} = \frac{N}{\sqrt{N}}$$

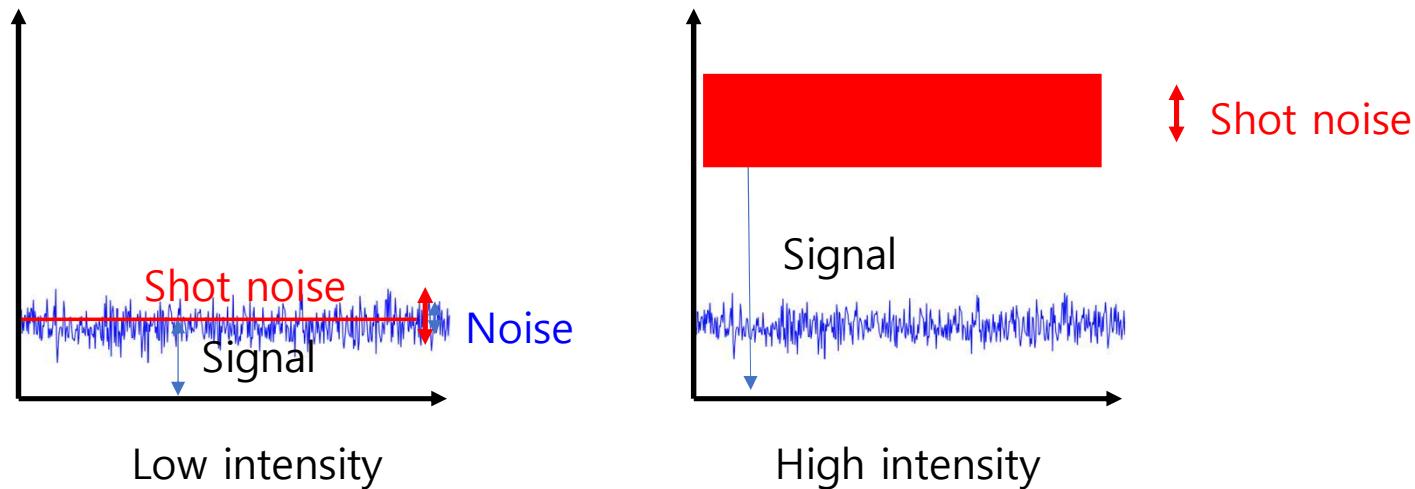
Shot noise of interferometer



$$\text{Signal to Noise ratio} = \frac{N}{\sqrt{N}}$$

If Shot noise is relatively larger than other noise(Thermal, Electric.. etc)
We say that it has **shot noise limit sensitivity**

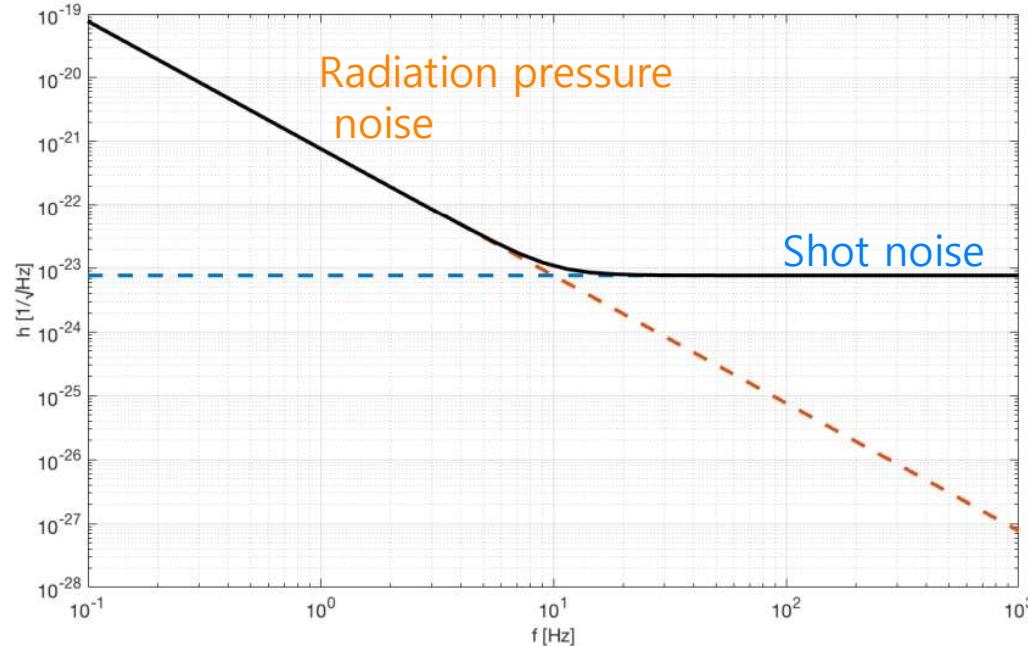
Shot noise of interferometer



$$\text{Signal to Noise ratio} = \frac{N}{\sqrt{N}} \quad \begin{aligned} N=100, \text{ SNR}=10 \\ N=10000, \text{ SNR}=100 \end{aligned}$$

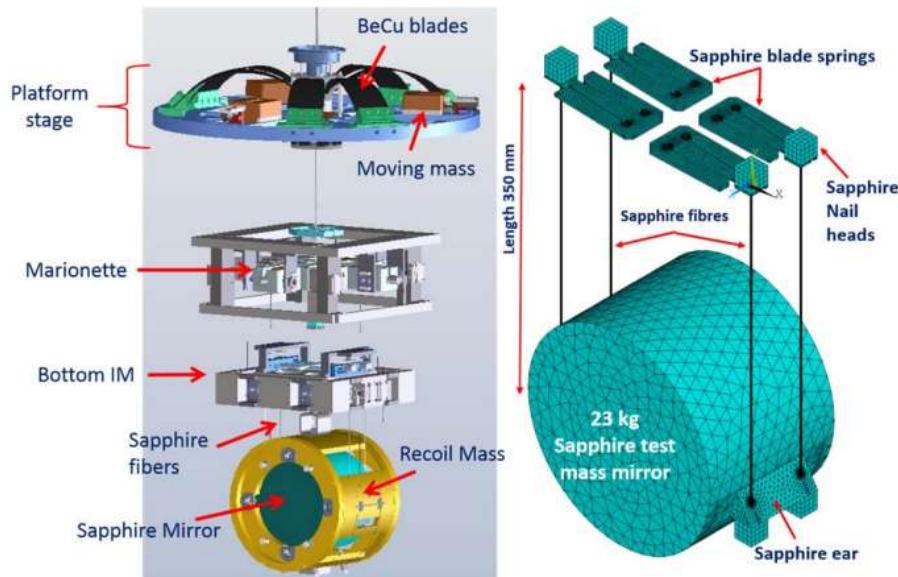
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Standard quantum limit of GW detector

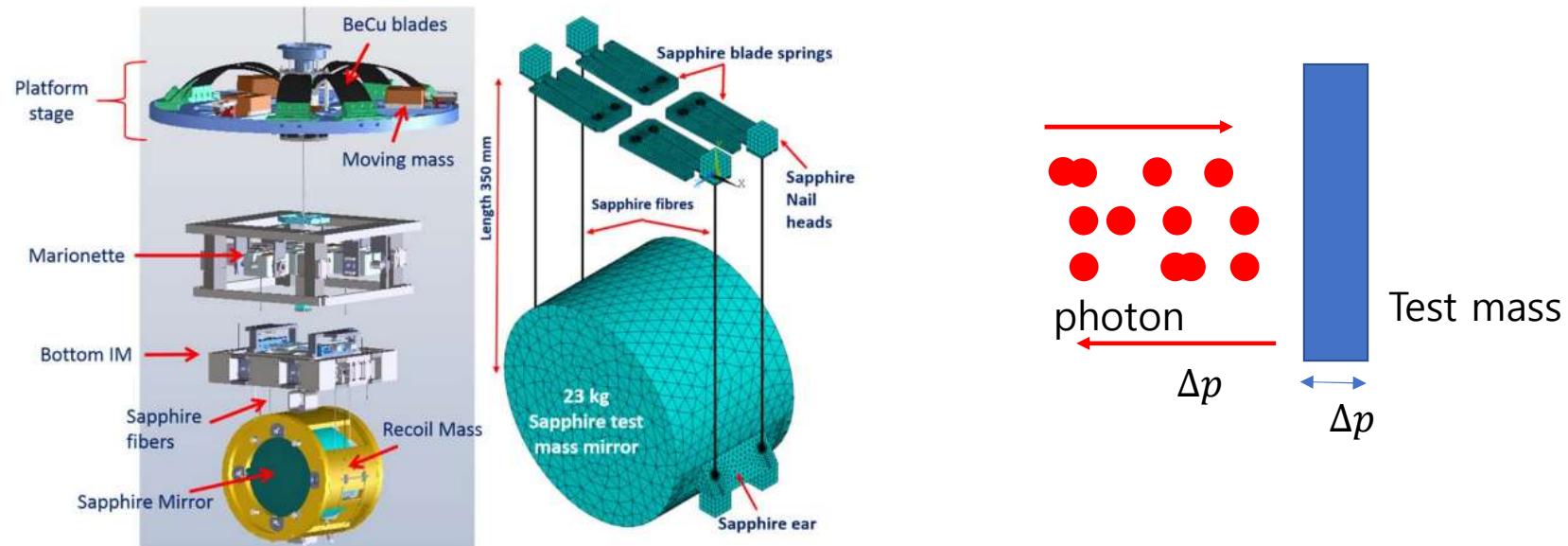


Standard quantum limit of gravitational wave detector
Shot noise + Radiation pressure noise

Radiation pressure noise

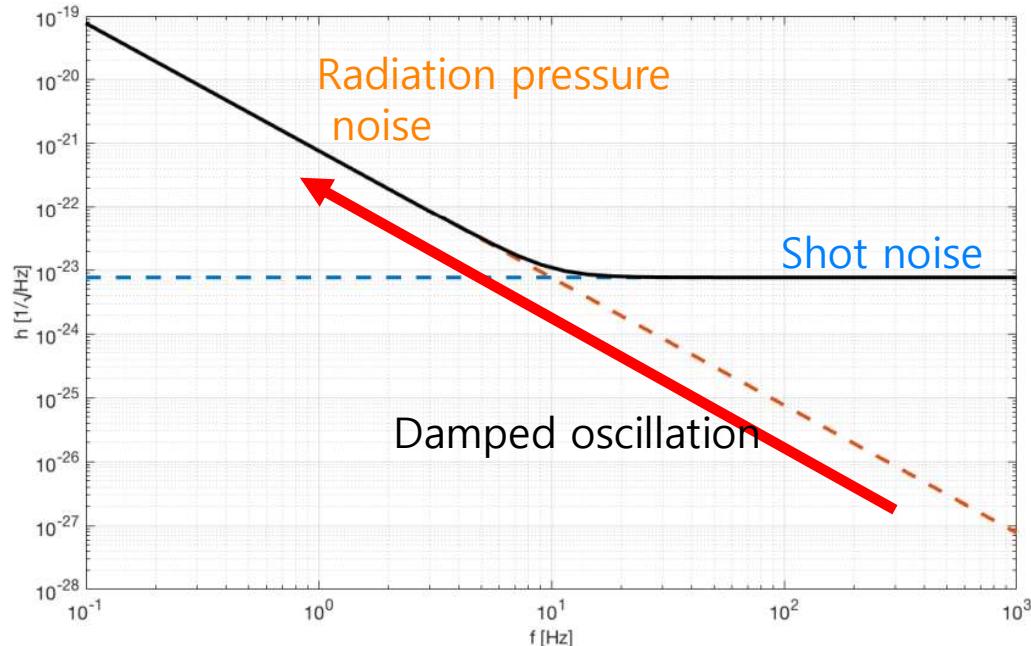


Radiation pressure noise



- Stored energy is very high (750 kW)
- Desired sensitivity is very high ($10^{-21} \sim 10^{-24}$)

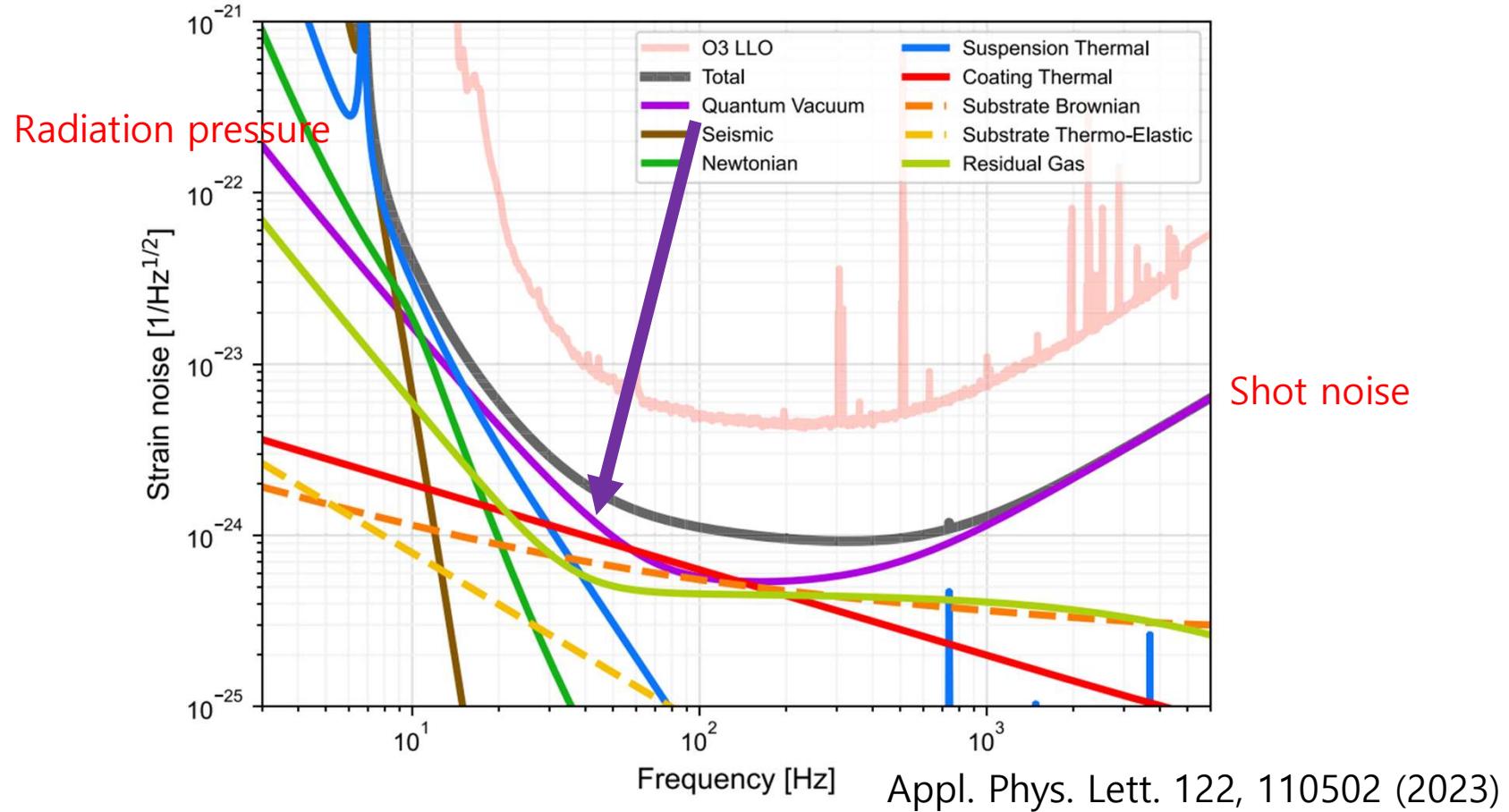
Standard quantum limit of GW detector



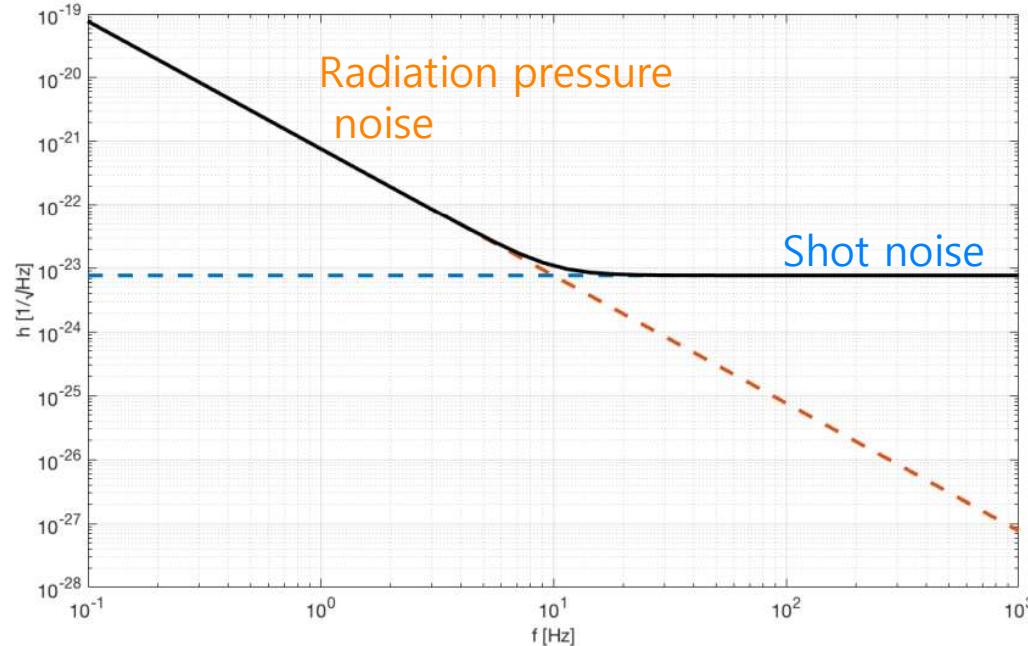
Standard quantum limit of gravitational wave detector
Shot noise + Radiation pressure noise

- Radiation pressure noise
 - Photon pressure fluctuation at **mirror**
 - > **Intensity noise**
- Shot noise
 - Photon number fluctuation at **PD**
 - > **Phase noise**

LIGO sensitivity

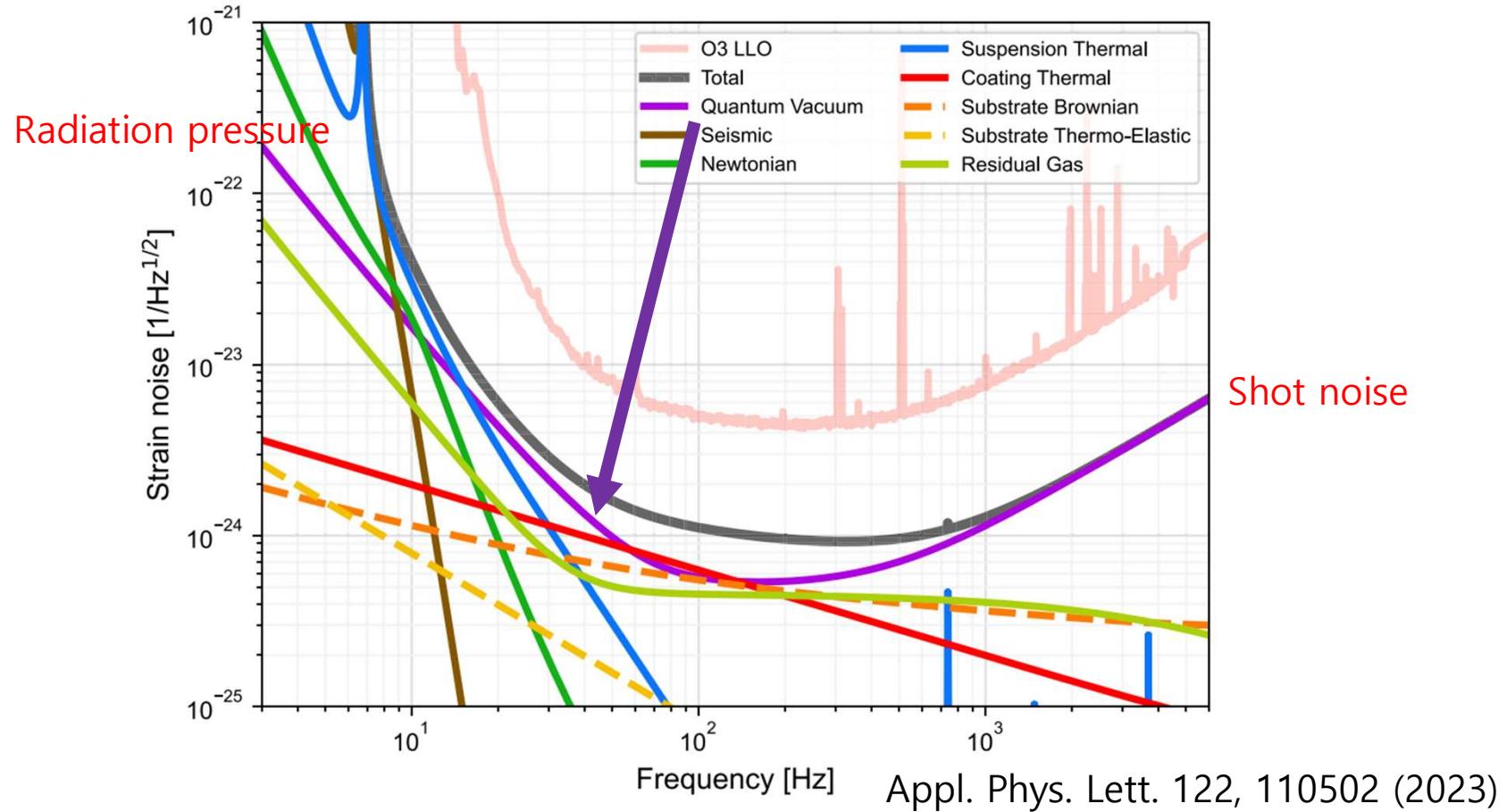


Standard quantum limit of GW detector

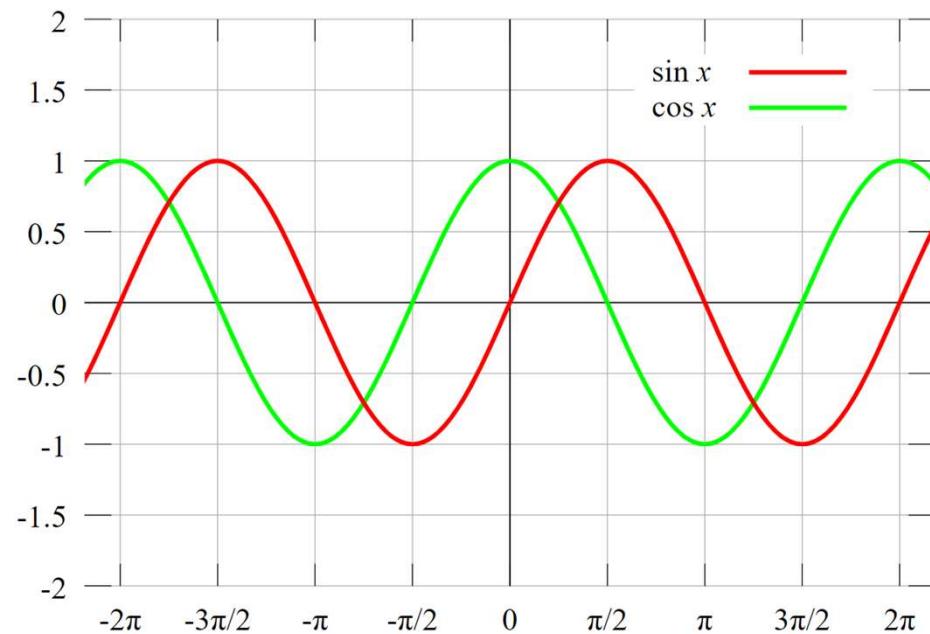


Standard quantum limit of gravitational wave detector
Shot noise + Radiation pressure noise

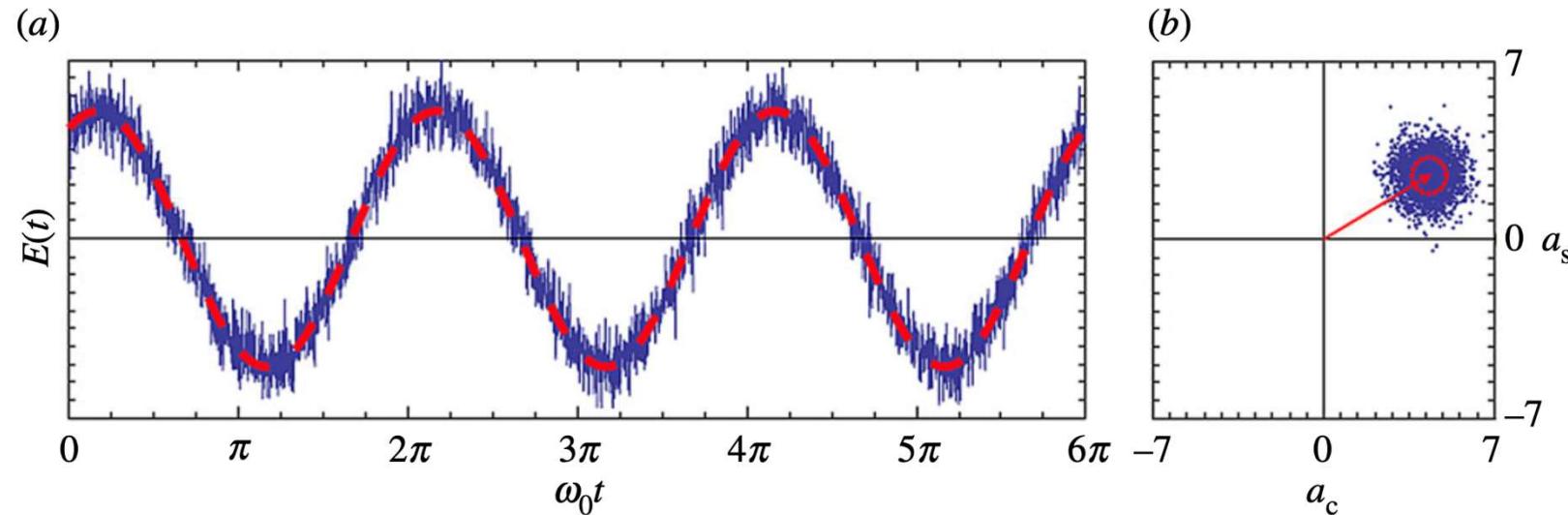
LIGO sensitivity



Classical electromagnetic wave

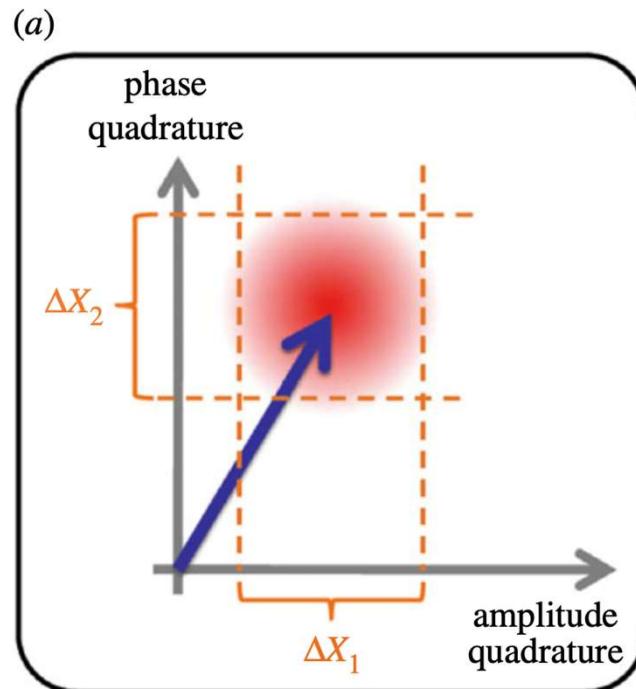


Quantum noise of coherent light



Heurs M. 2018 Gravitational wave detection using laser interferometry beyond the standard quantum limit. Phil. Trans. R. Soc. A 376: 20170 289.

Phase and amplitude quadrature



$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

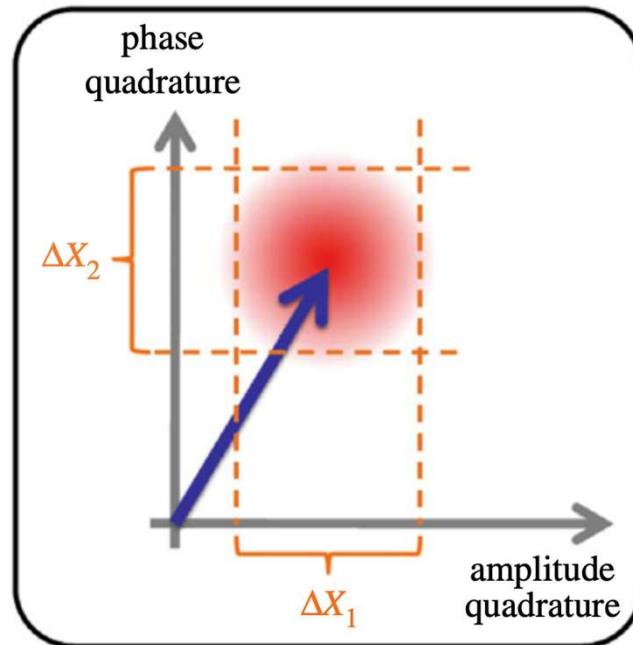
Uncertainty principle

Heurs M. 2018 Gravitational wave detection using laser interferometry beyond the standard quantum limit. Phil. Trans. R. Soc. A 376: 20170 289.

Phase and amplitude quadrature



(a) Shot noise



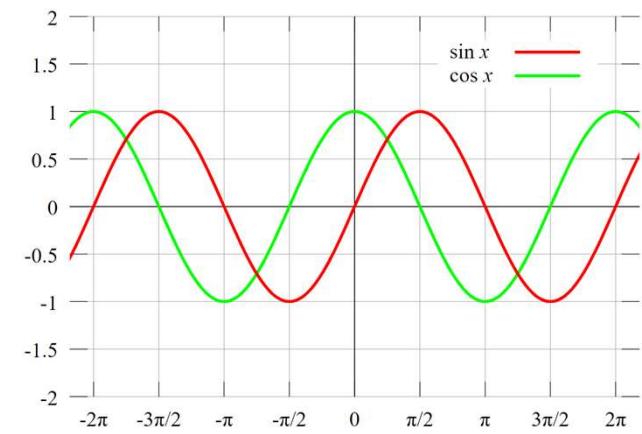
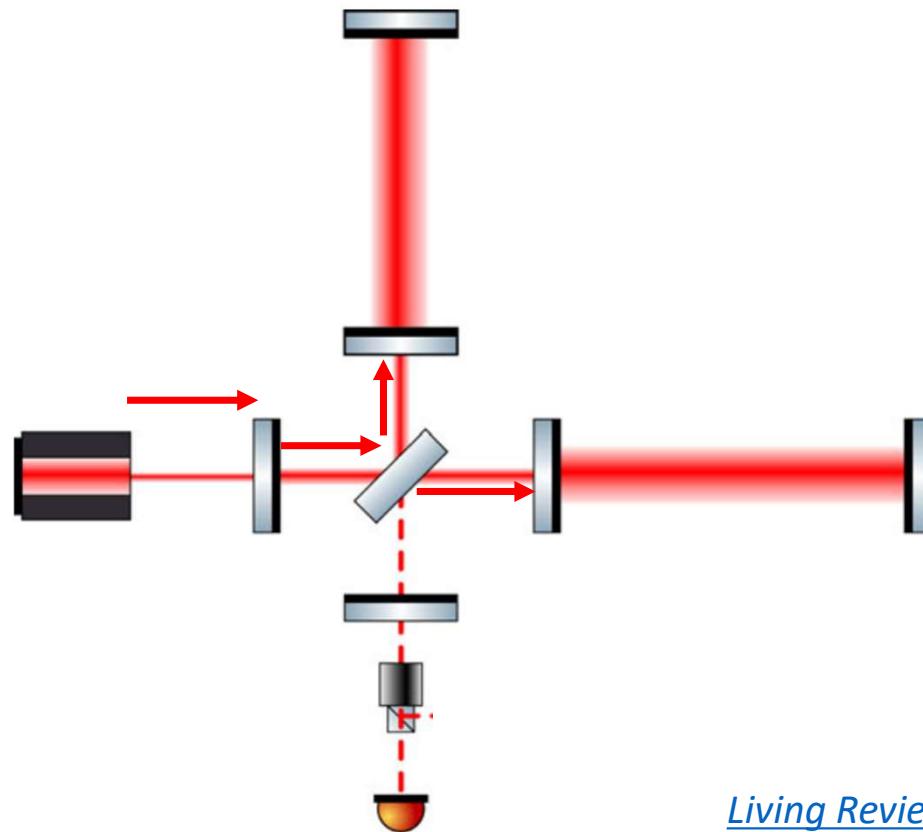
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Uncertainty principle

Radiation pressure noise

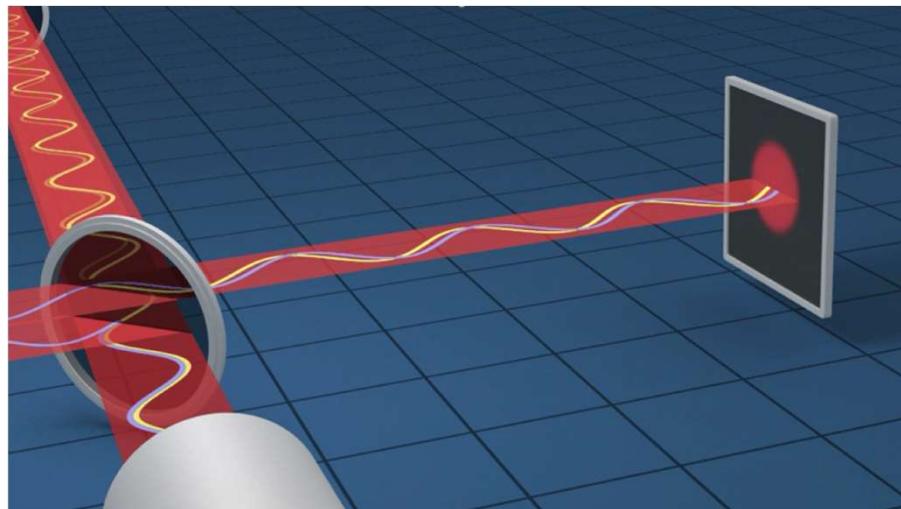
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Gravitational wave detector

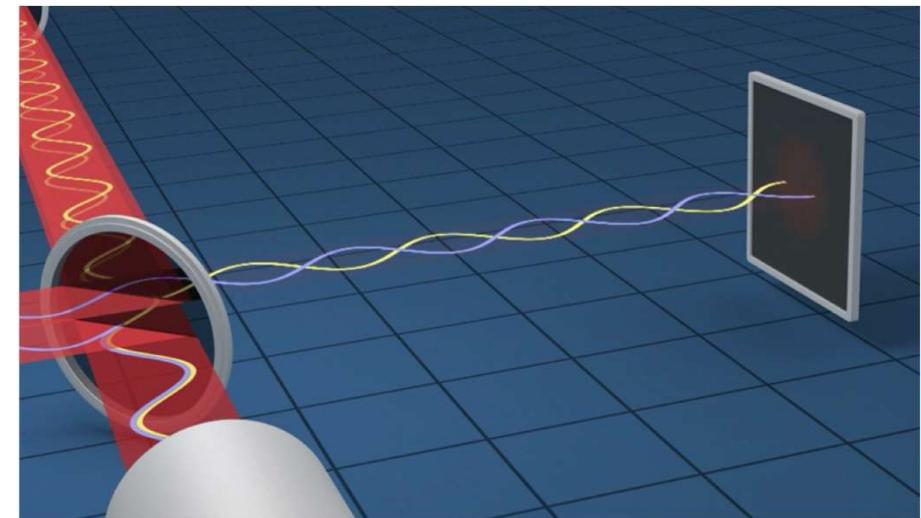


[Living Reviews in Relativity](https://lrr.springeropen.com/articles/10.12893/lrr-2019-0002) volume 22, Article number: 2 (2019)

Bright port and dark port

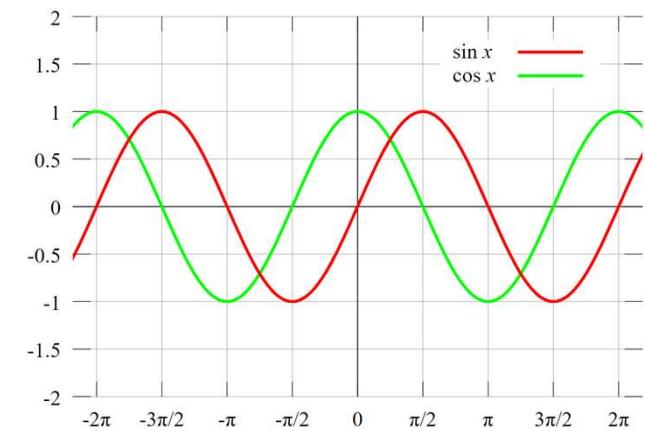
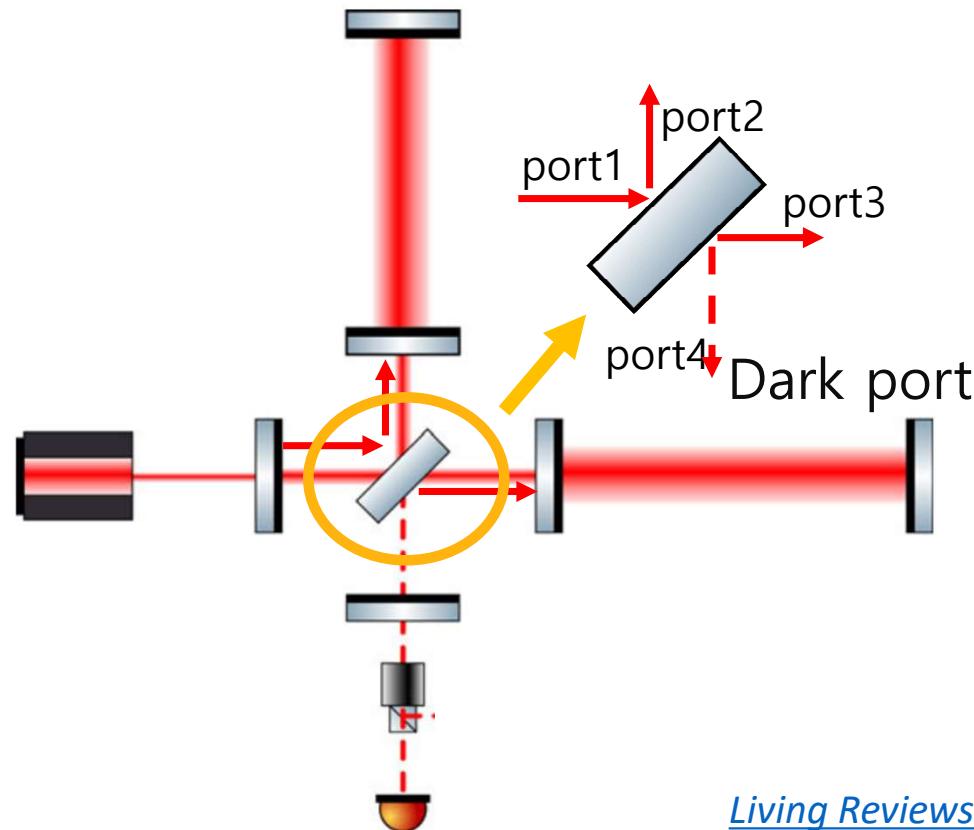


Bright



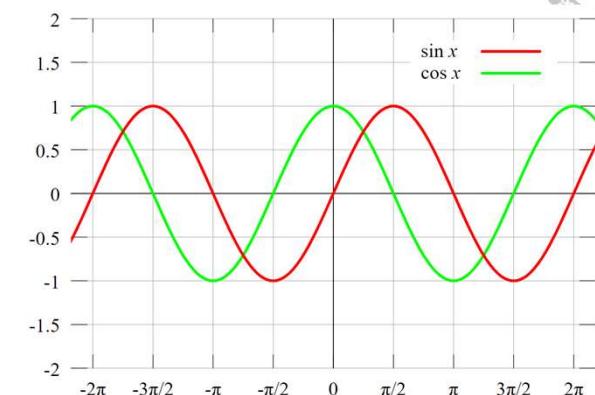
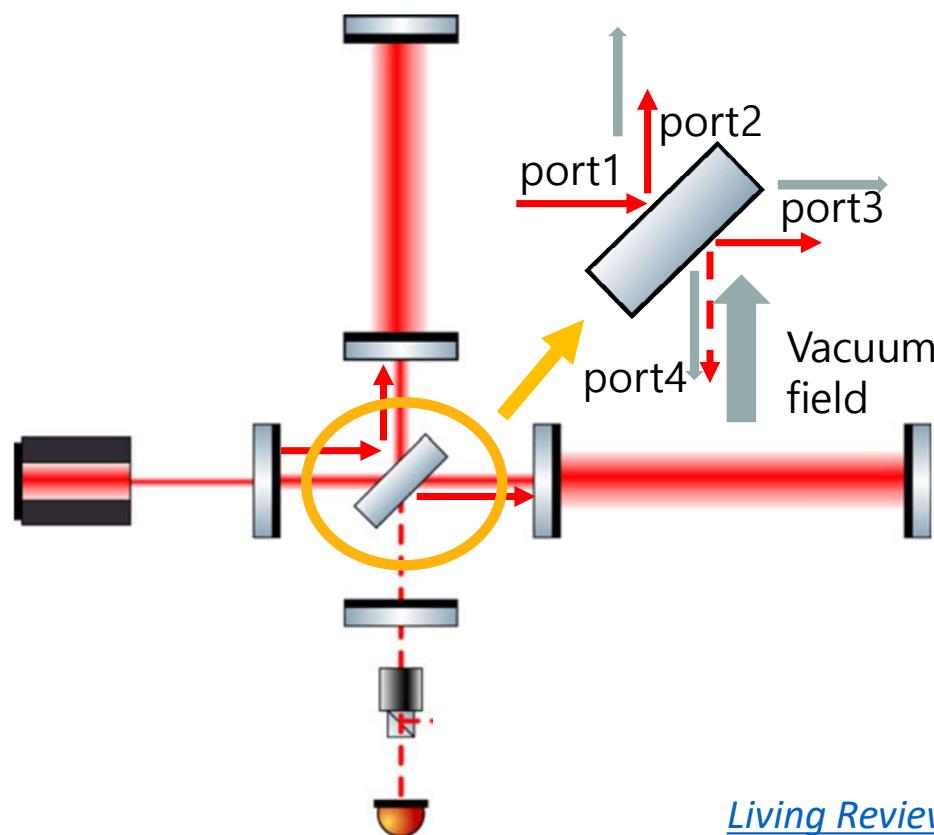
Dark

Gravitational wave detector



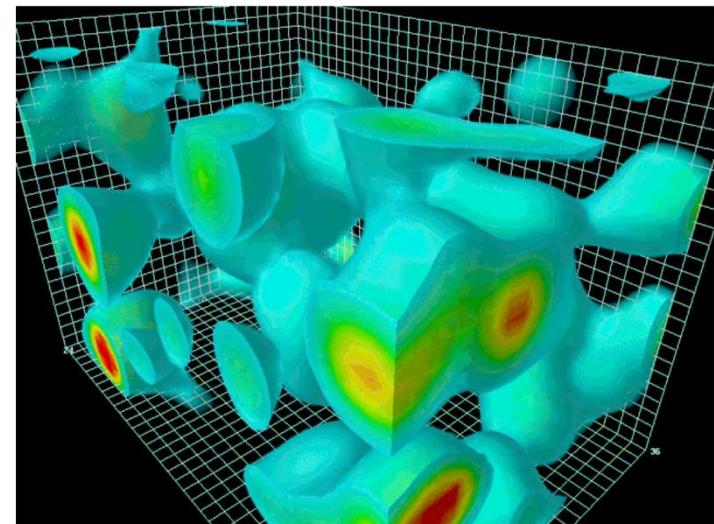
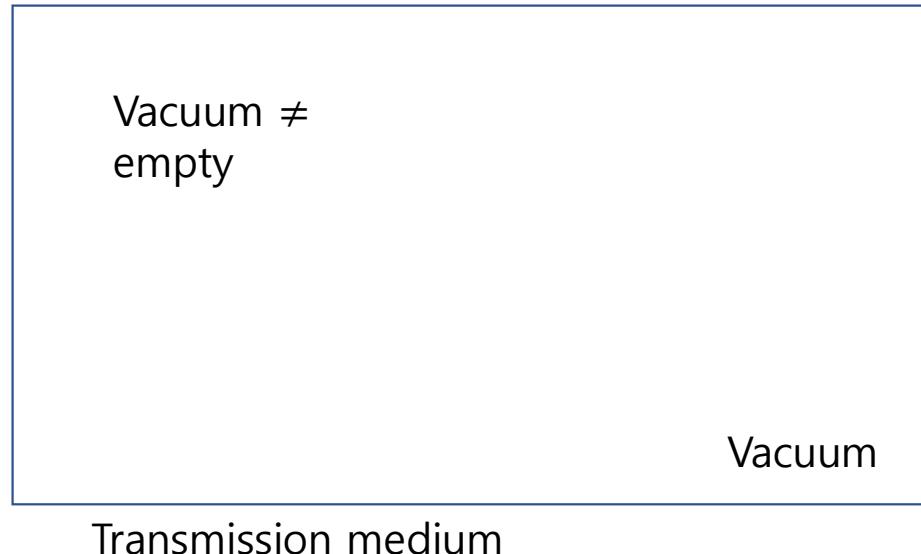
[Living Reviews in Relativity](#) volume 22, Article number: 2 (2019)

Quantum noise of gravitational wave detector



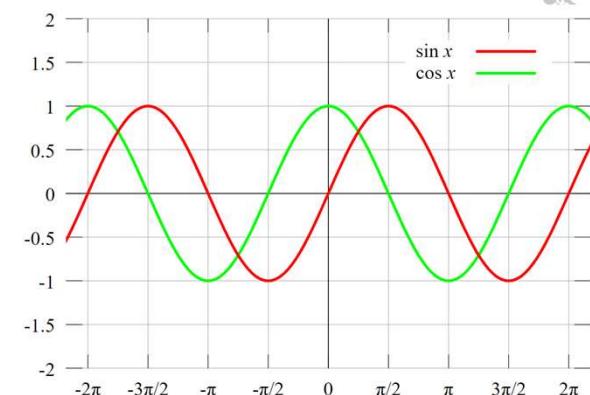
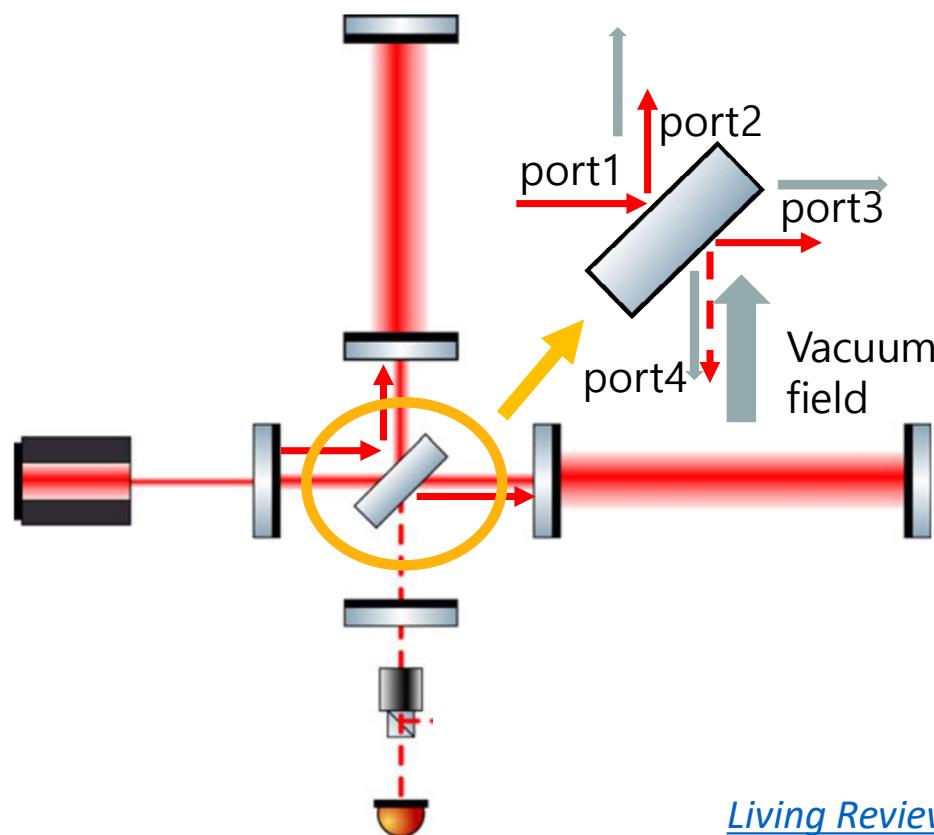
[Living Reviews in Relativity](#) volume 22, Article number: 2 (2019)

Quantum vacuum fluctuation



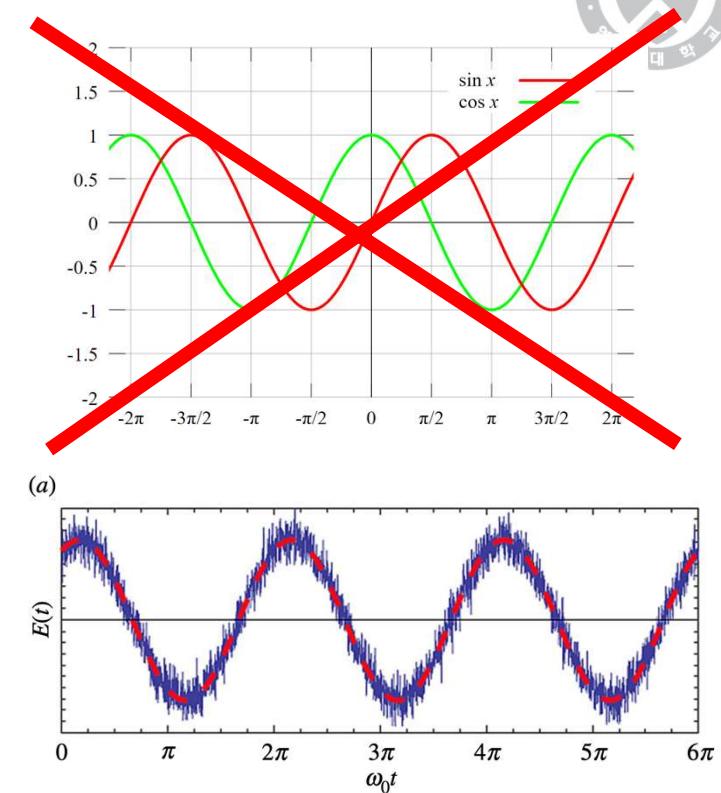
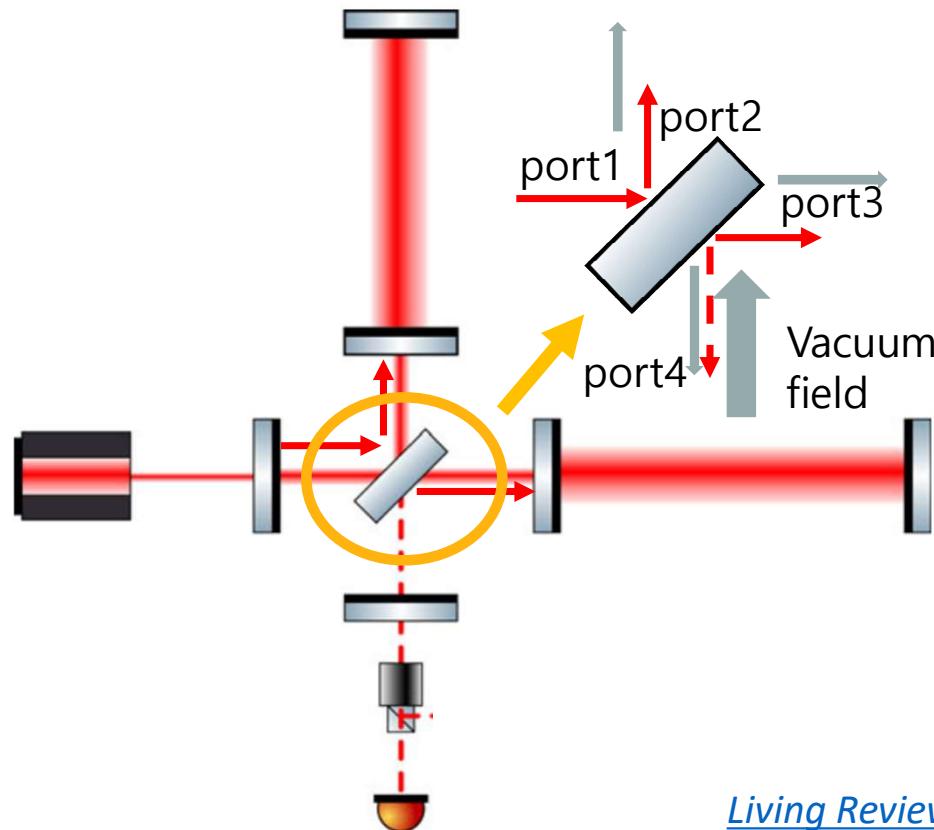
Average energy = zeropoint energy

Quantum noise of gravitational wave detector



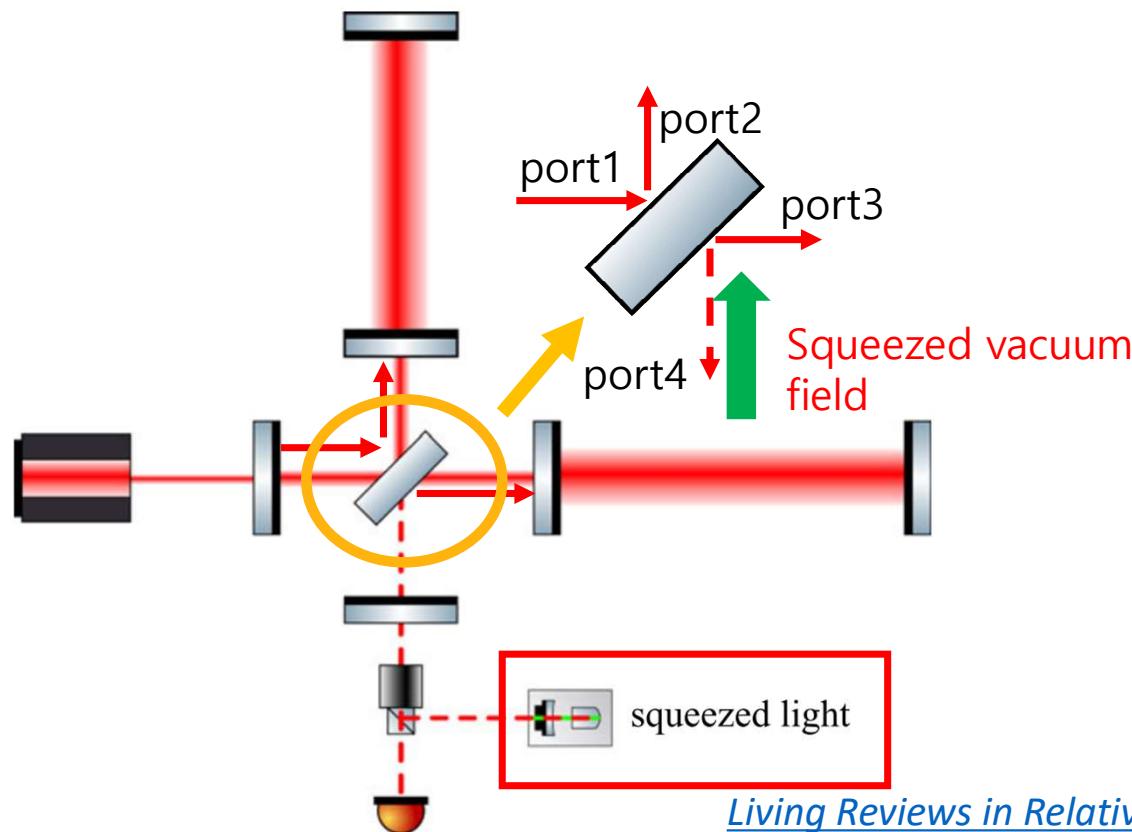
[Living Reviews in Relativity](#) volume 22, Article number: 2 (2019)

Quantum noise of gravitational wave detector



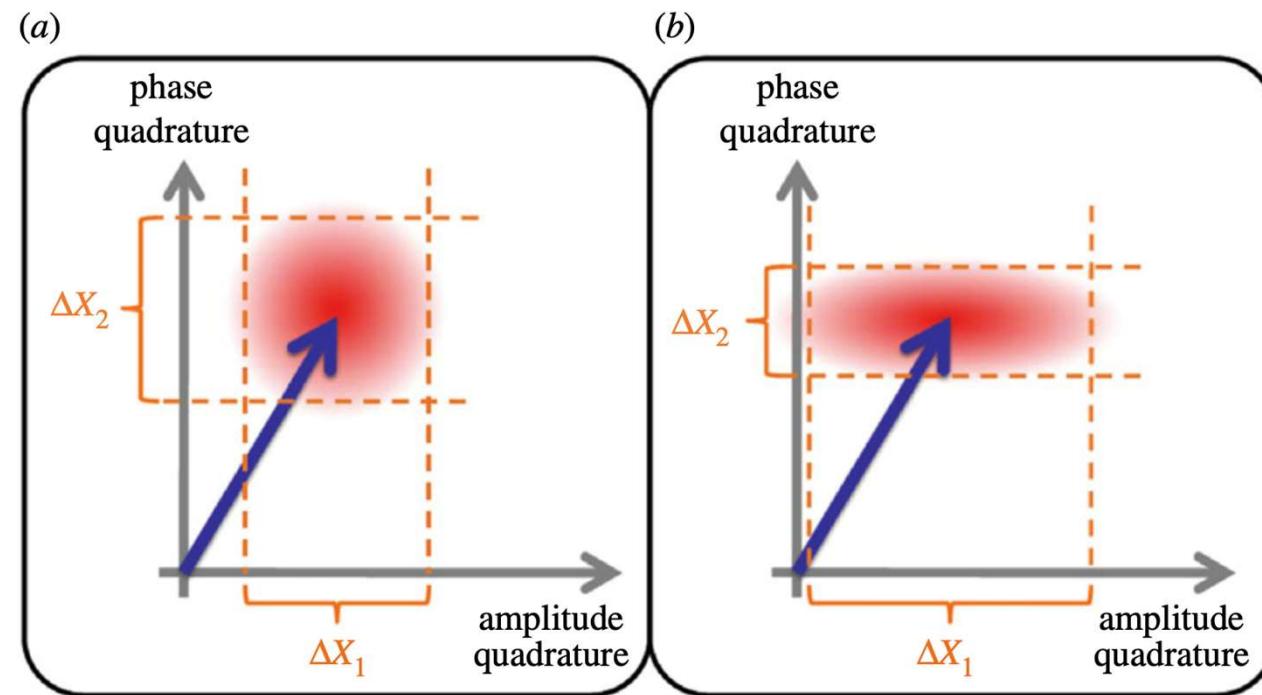
[Living Reviews in Relativity](#) volume 22, Article number: 2 (2019)

Gravitational wave detector



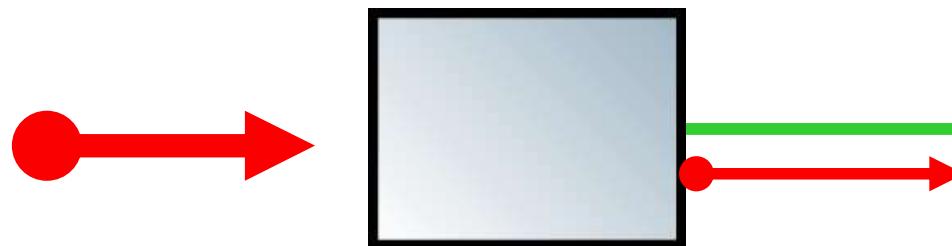
[*Living Reviews in Relativity*](#) volume 22, Article number: 2 (2019)

Squeezed light

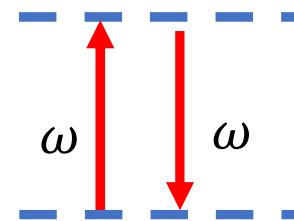


Heurs M. 2018 Gravitational wave detection using laser interferometry beyond the standard quantum limit. Phil. Trans. R. Soc. A 376: 20170 289.

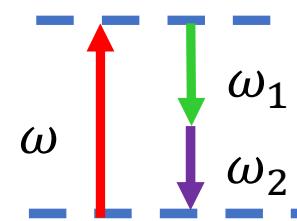
Non linear crystal



Non-linear crystal

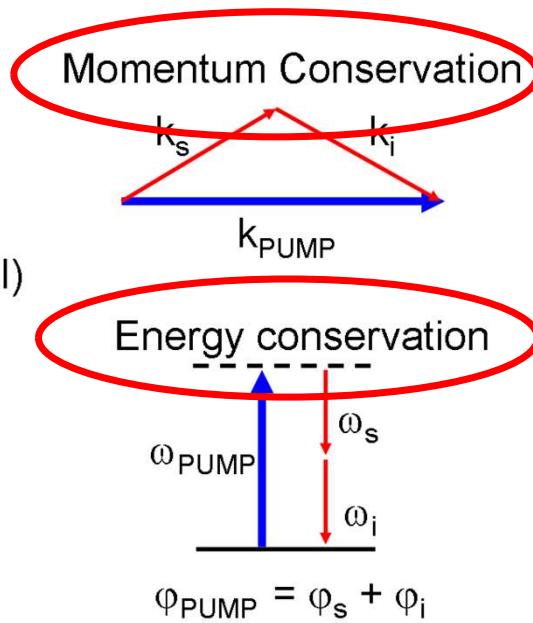
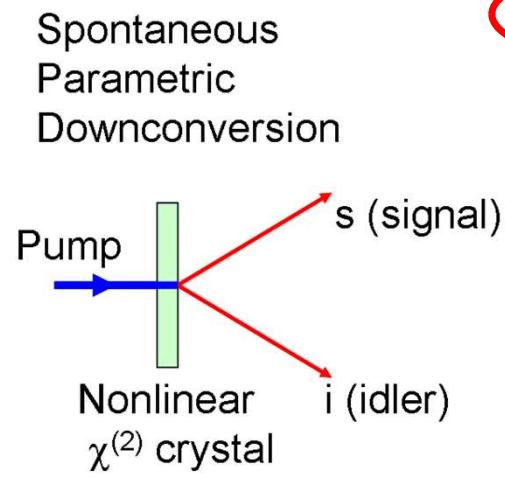


linear optics



Non linear optics

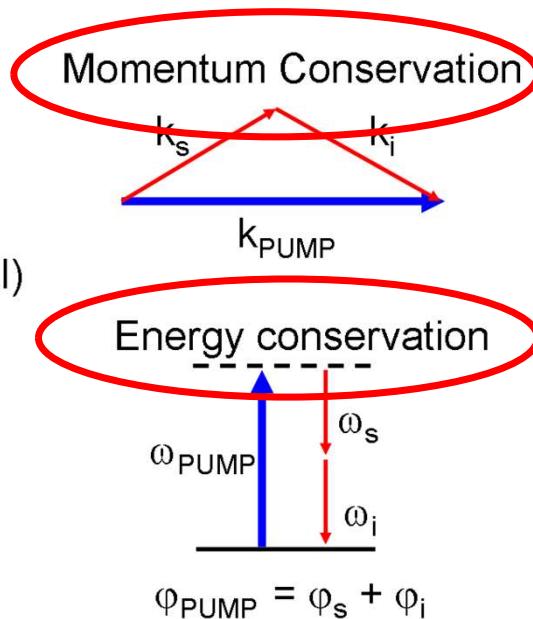
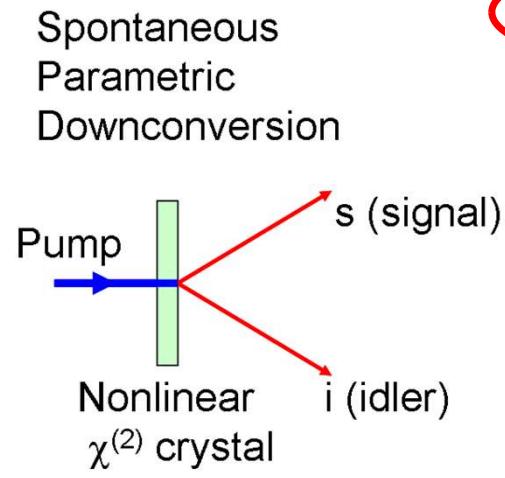
Parametric down conversion



Every field which are participated in Non-linear conversion process, Must obey energy, momentum conservation.

[J S Lundeen](#)

Parametric down conversion

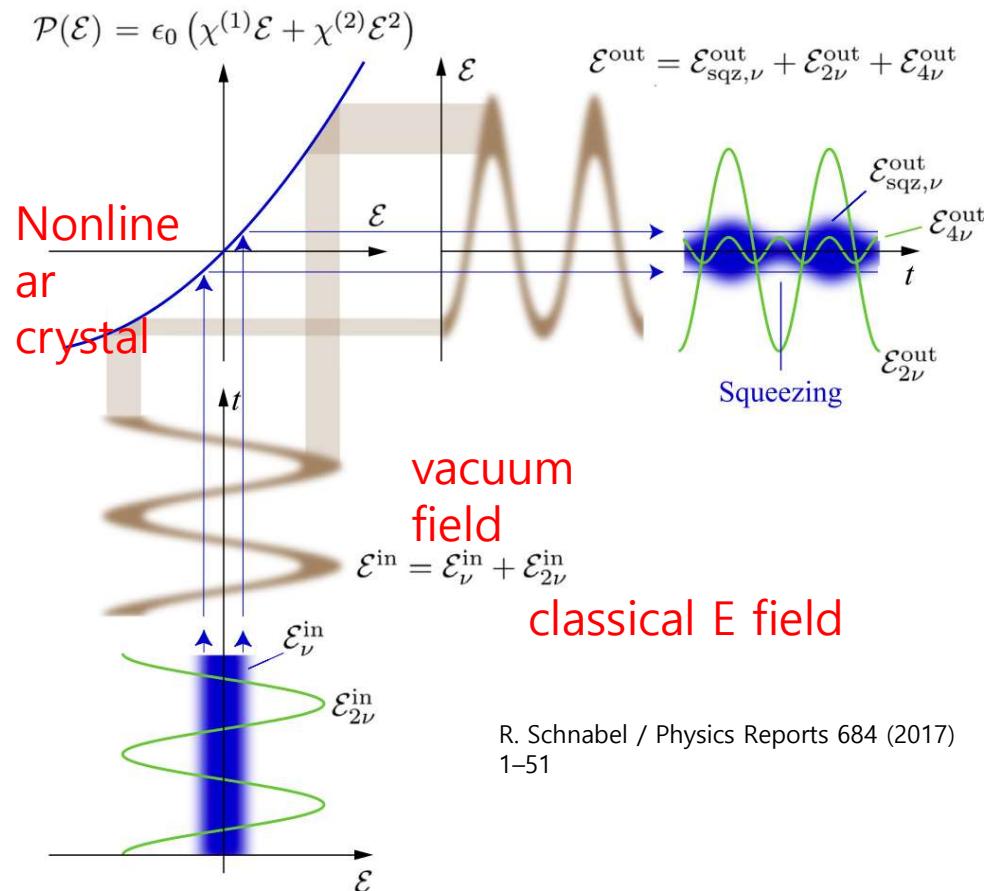


Every field which are participated in Non-linear conversion process,
Must obey energy, momentum conservation.

Even for vacuum field!!

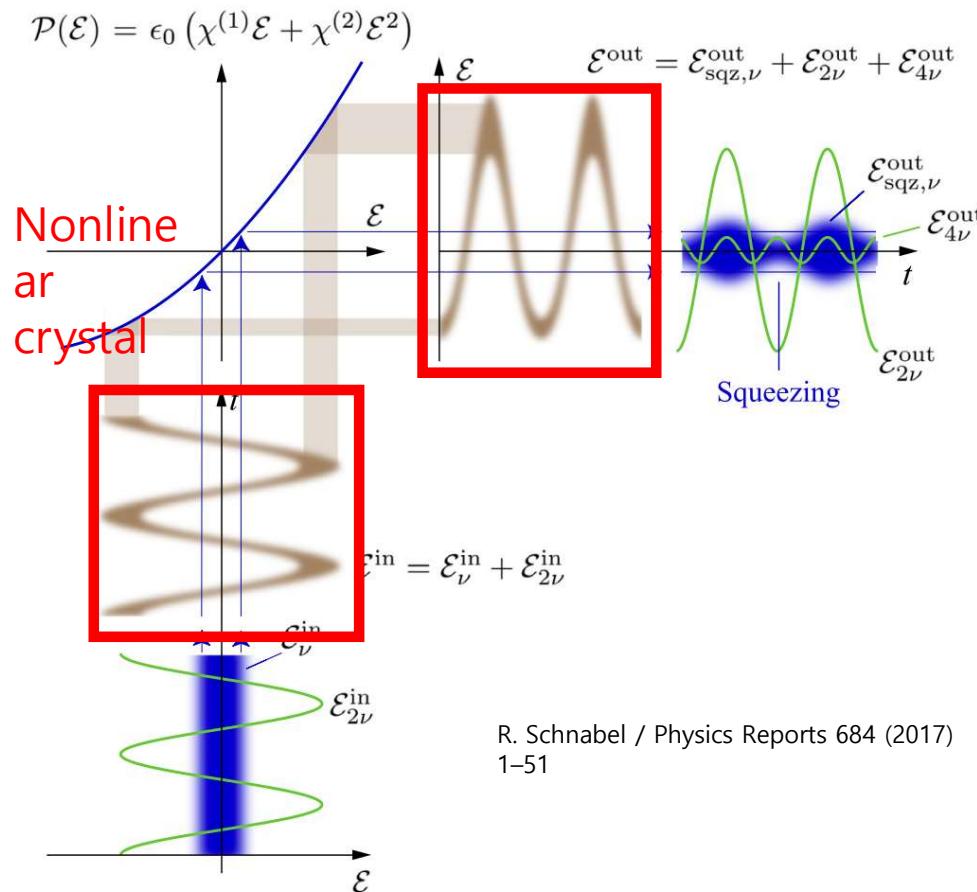
[J S Lundeen](#)

Parametric down conversion

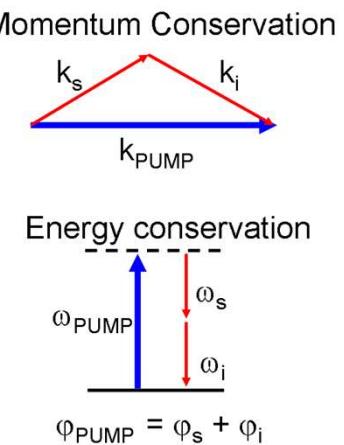
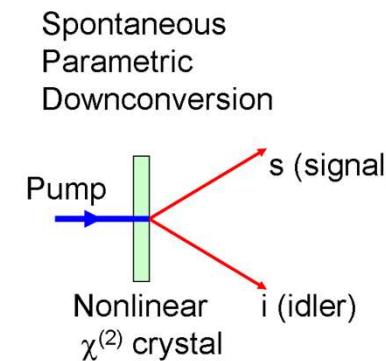


R. Schnabel / Physics Reports 684 (2017)
1–51

Parametric down conversion

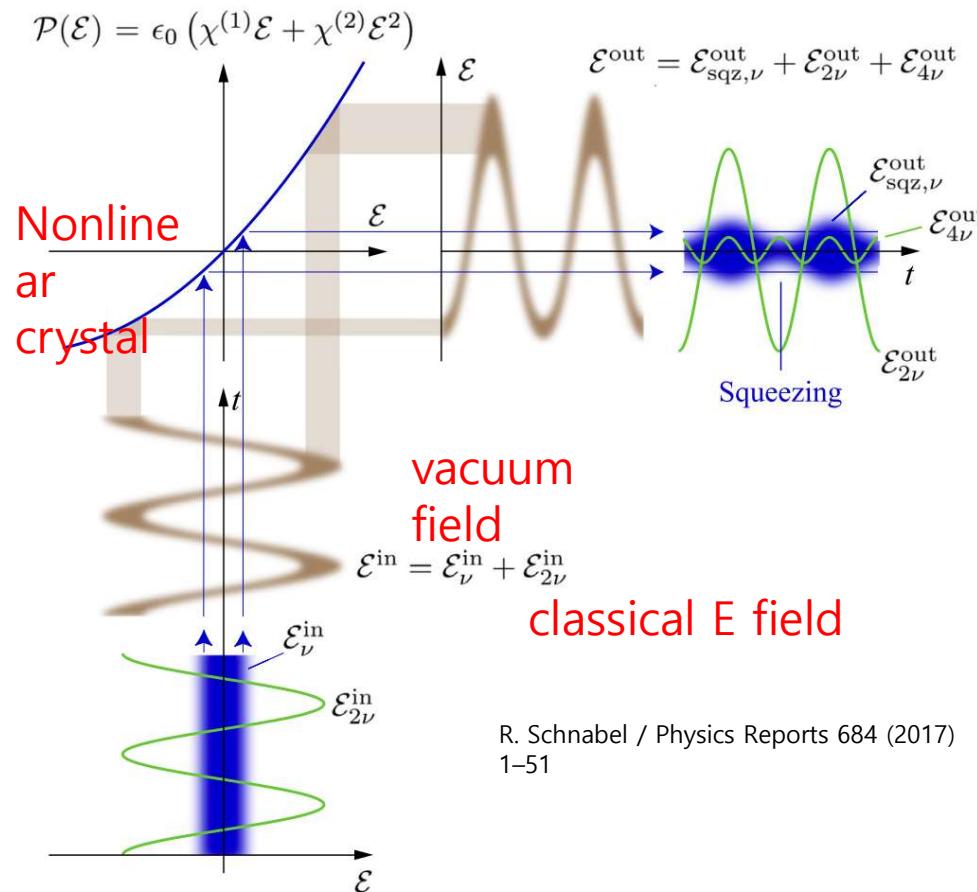


R. Schnabel / Physics Reports 684 (2017)
1–51



Boundary condition

Parametric down conversion



R. Schnabel / Physics Reports 684 (2017)
1–51

State of vacuum is changed during conversion process due to momentum and energy conservation

Boundary condition



Squeezed vacuum

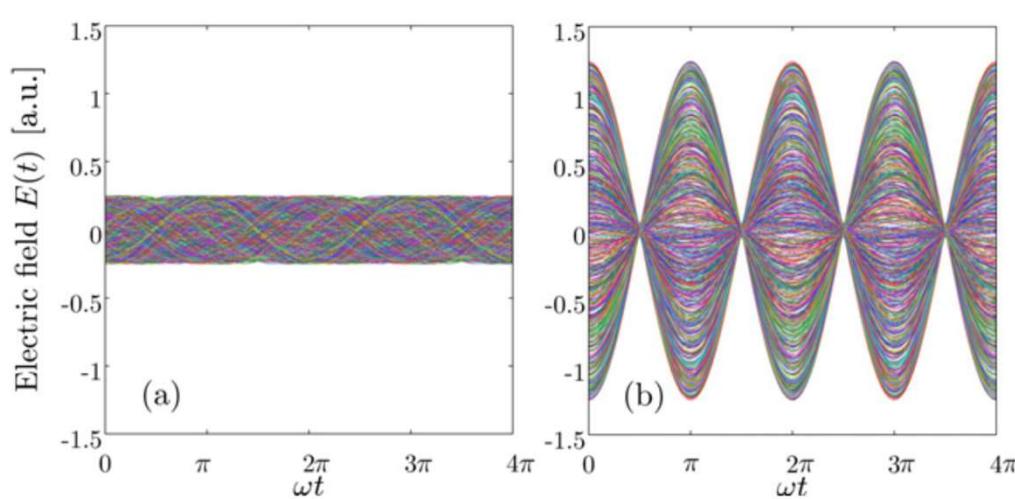
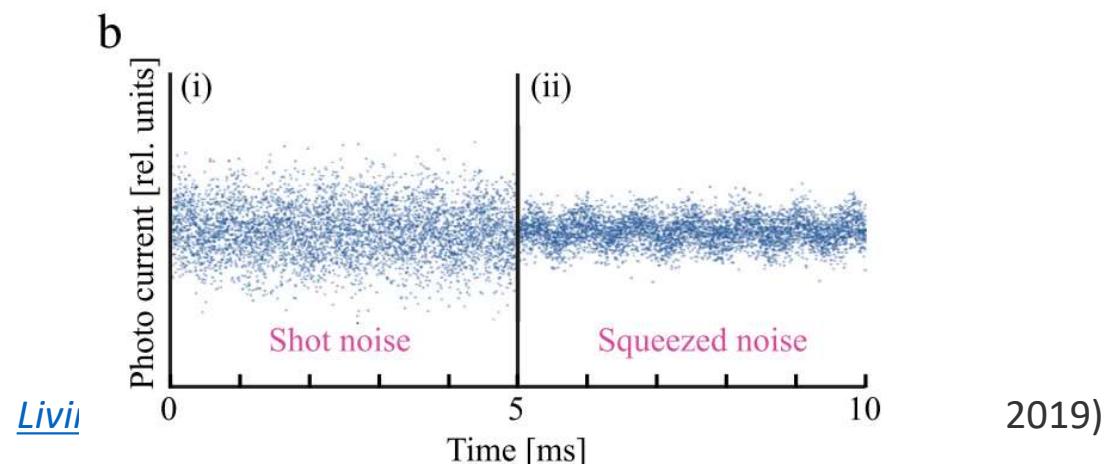
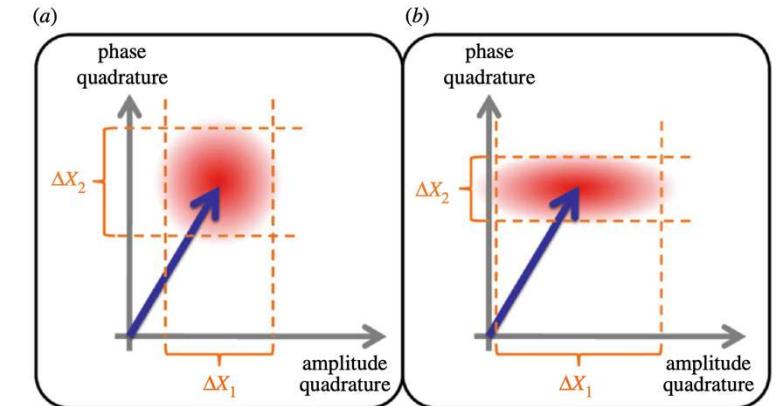
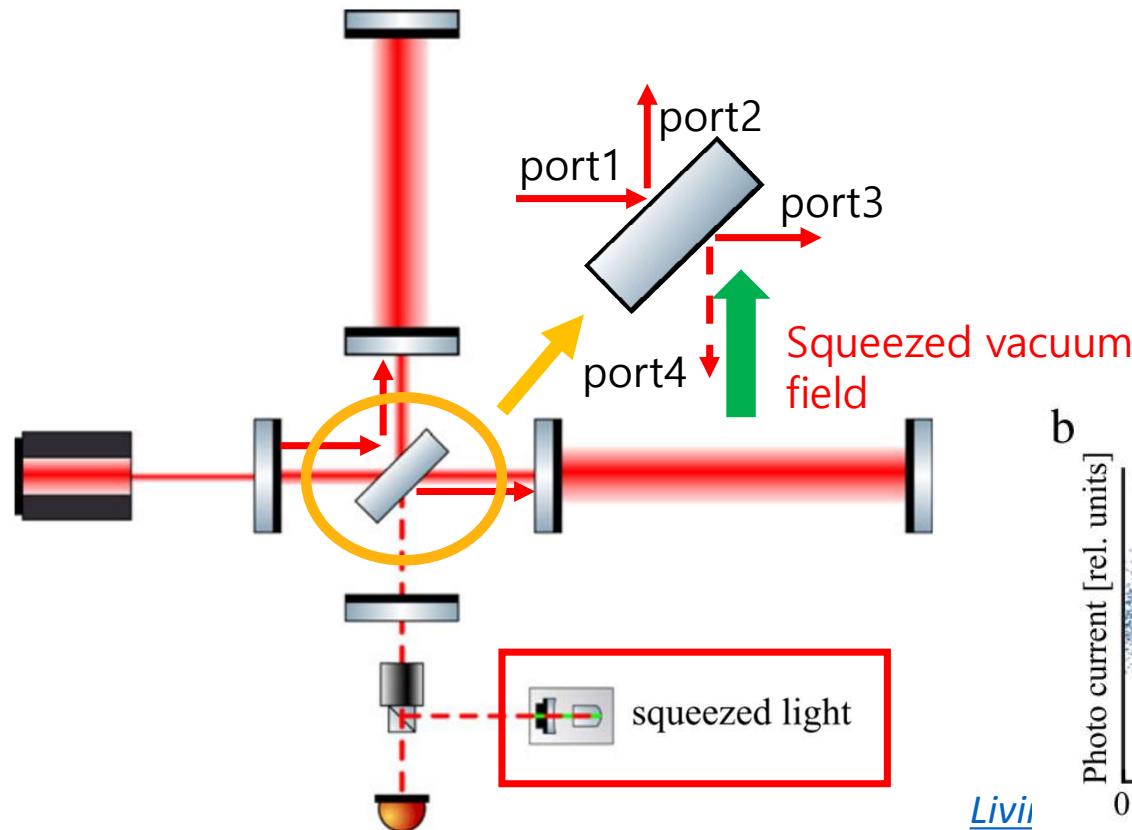
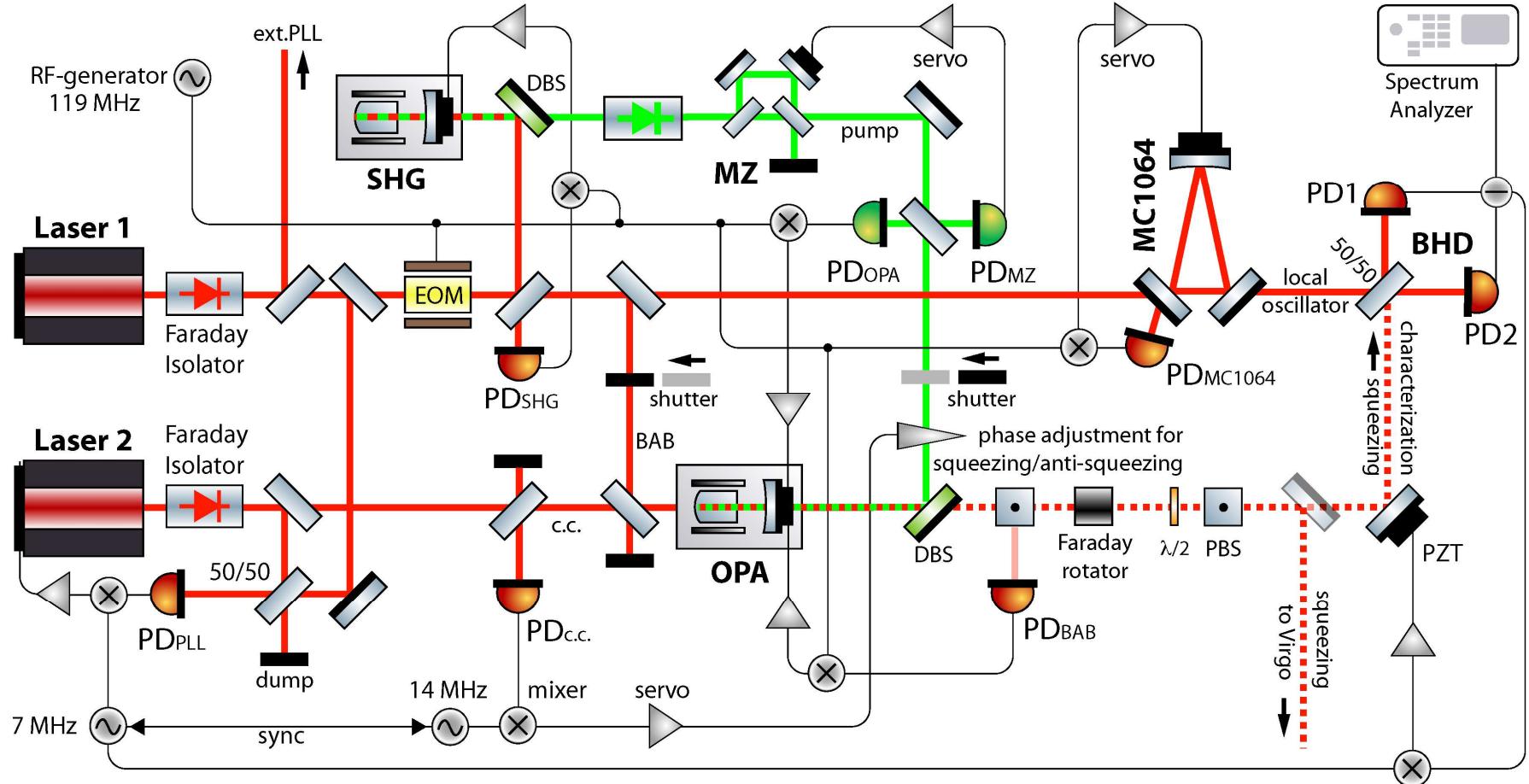


Figure 1.6: Simulation of electric field in time for (a) vacuum state and for (b) squeezed vacuum.

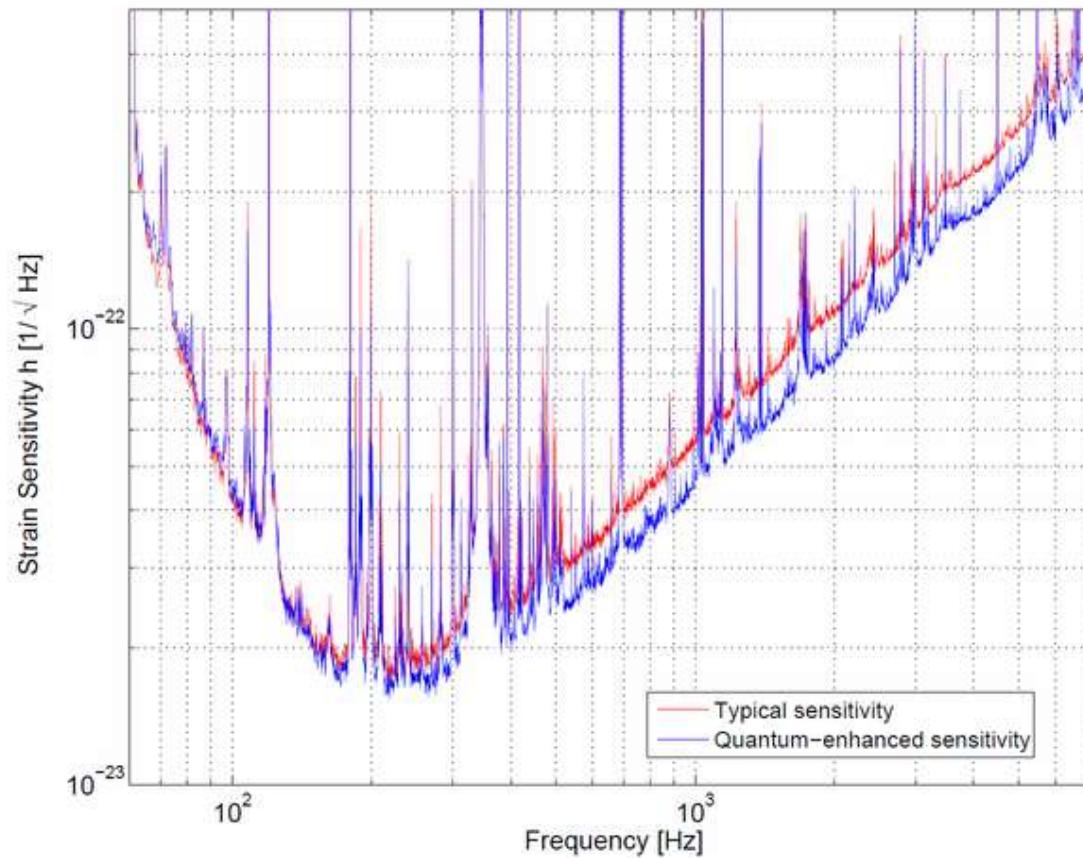
Gravitational wave detector



Livit



Quantum noise enhancement of LIGO



오분순삭

01:37



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오늘순식

MARKET IN M

01:35

아니요



지금 지쳤나요?

오늘순식

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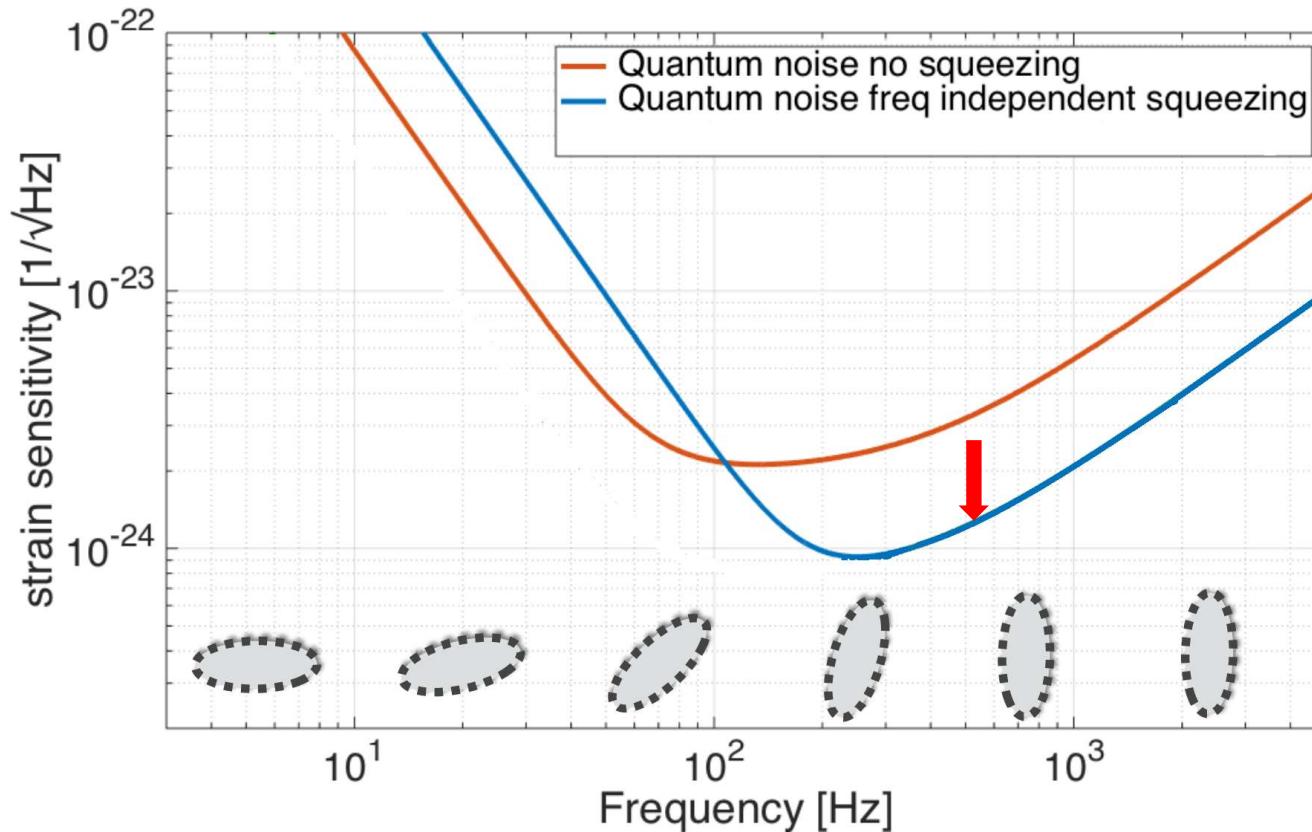
단호함과 양반된 움짤 그래프



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Frequency independent squeezing



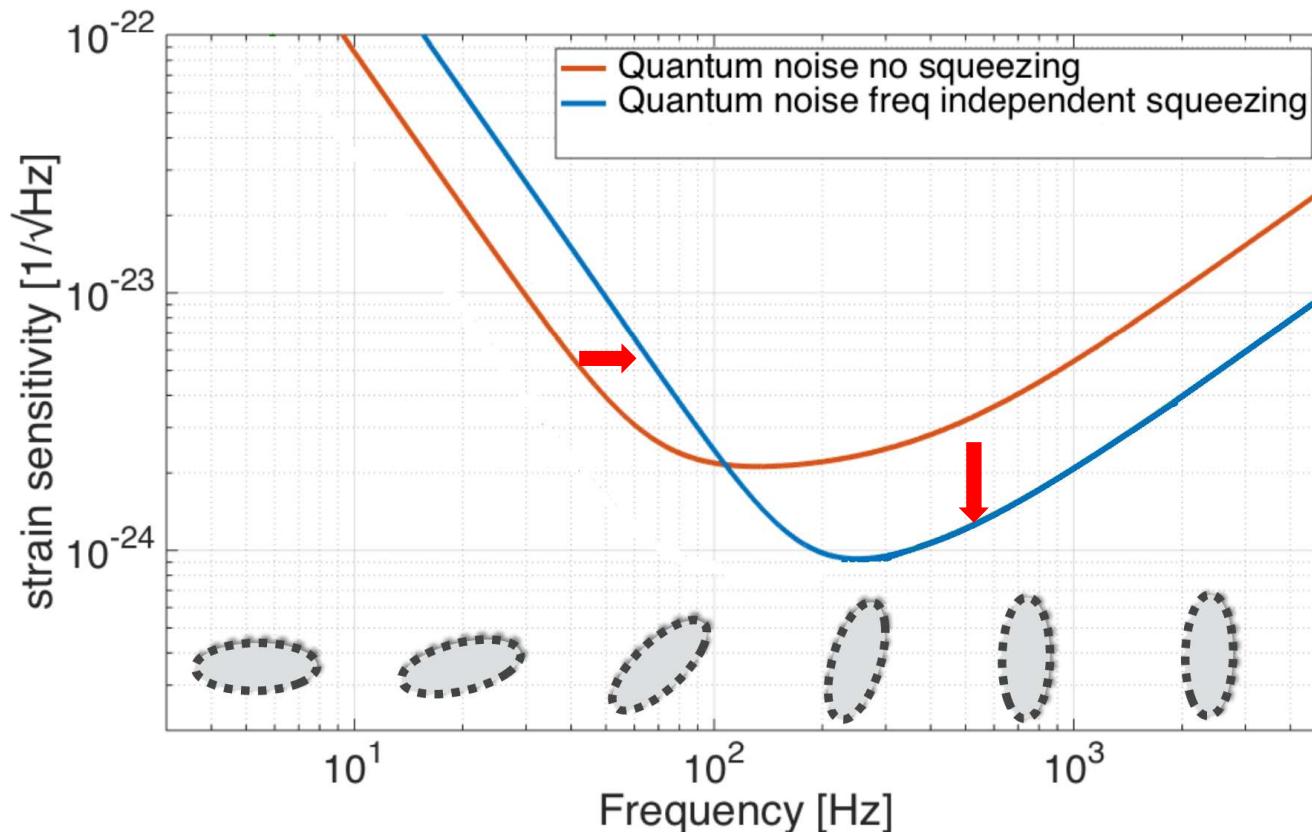
Optical and noise studies for Advanced Virgo and filter cavities for quantum noise reduction

in gravitational-wave interferometric detectors, Eleonora Capocasa, UNIVERSITÉ PARIS DIDEROT (2017)

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Frequency independent squeezing



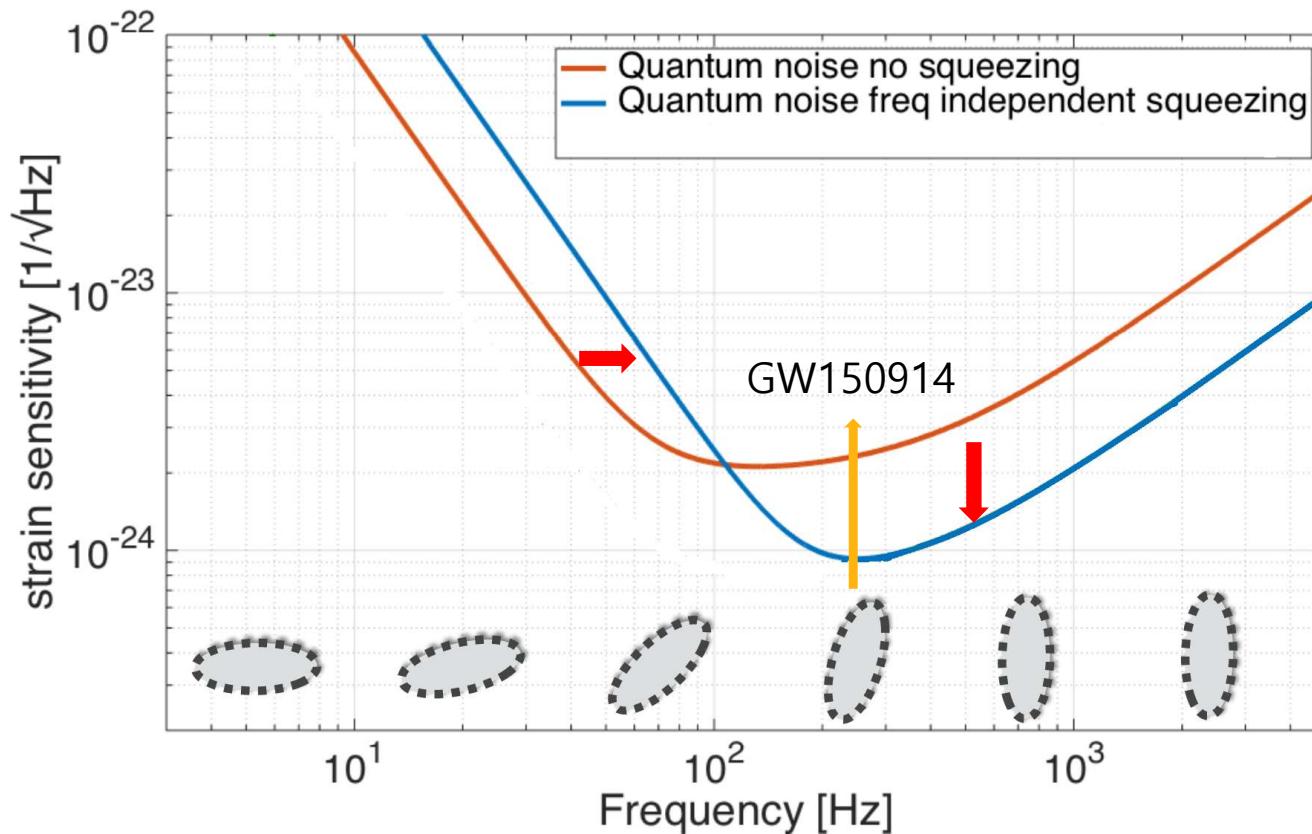
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Frequency independent squeezing



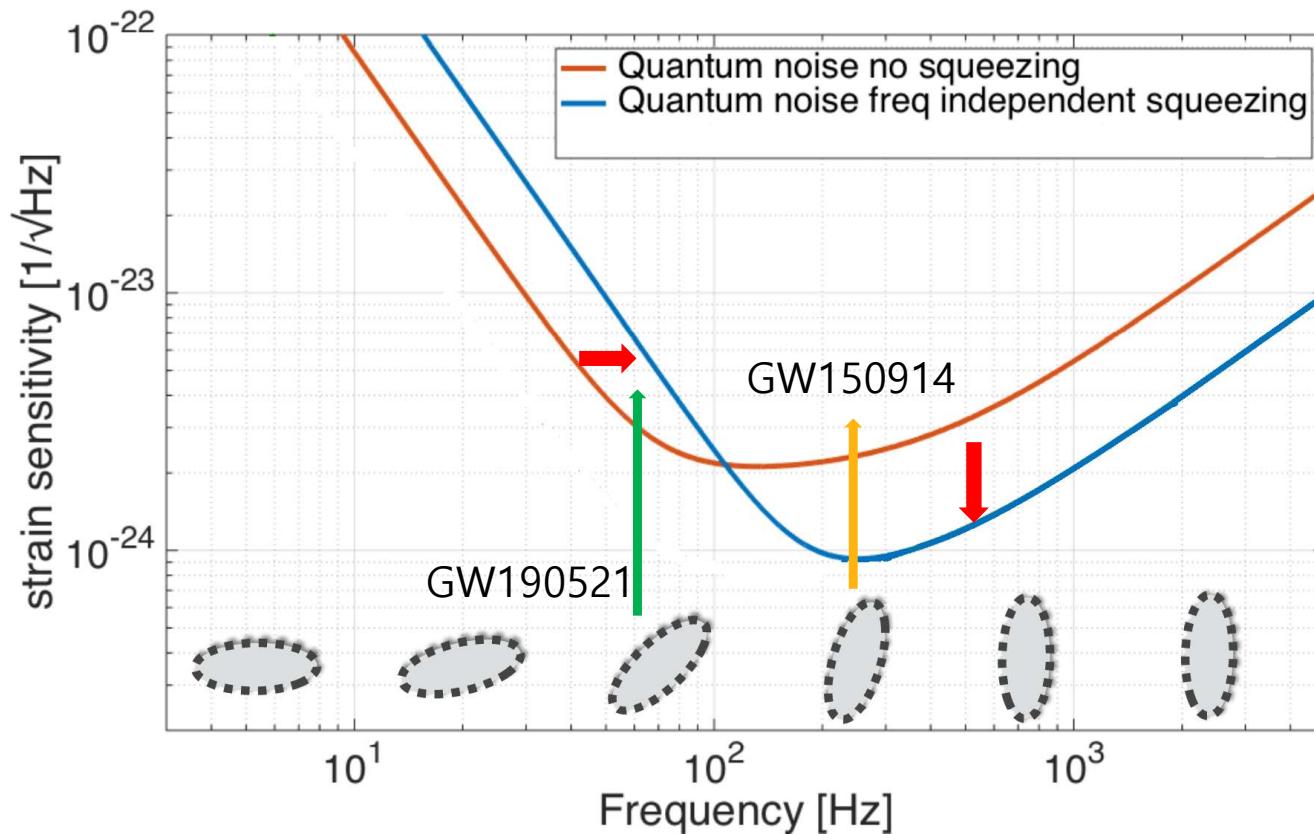
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Frequency independent squeezing



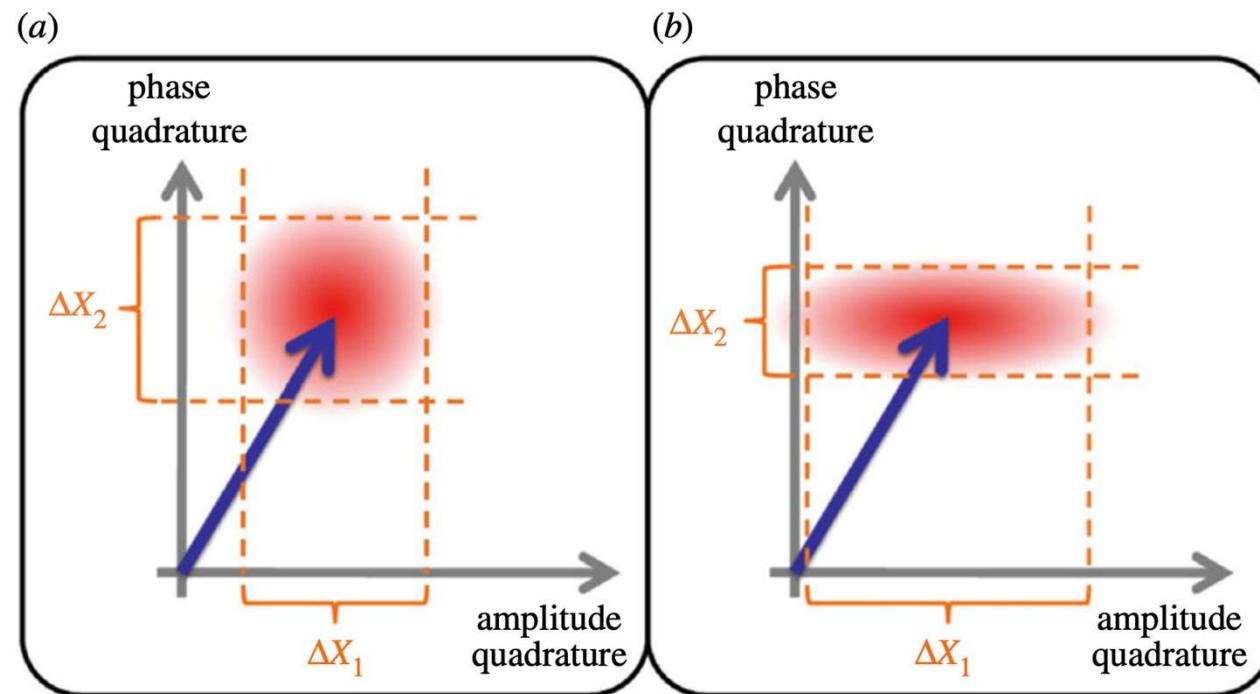
Optical and noise studies for Advanced Virgo and filter cavities for quantum noise reduction

in gravitational-wave interferometric detectors, Eleonora Capocasa, UNIVERSITÉ PARIS DIDEROT (2017)

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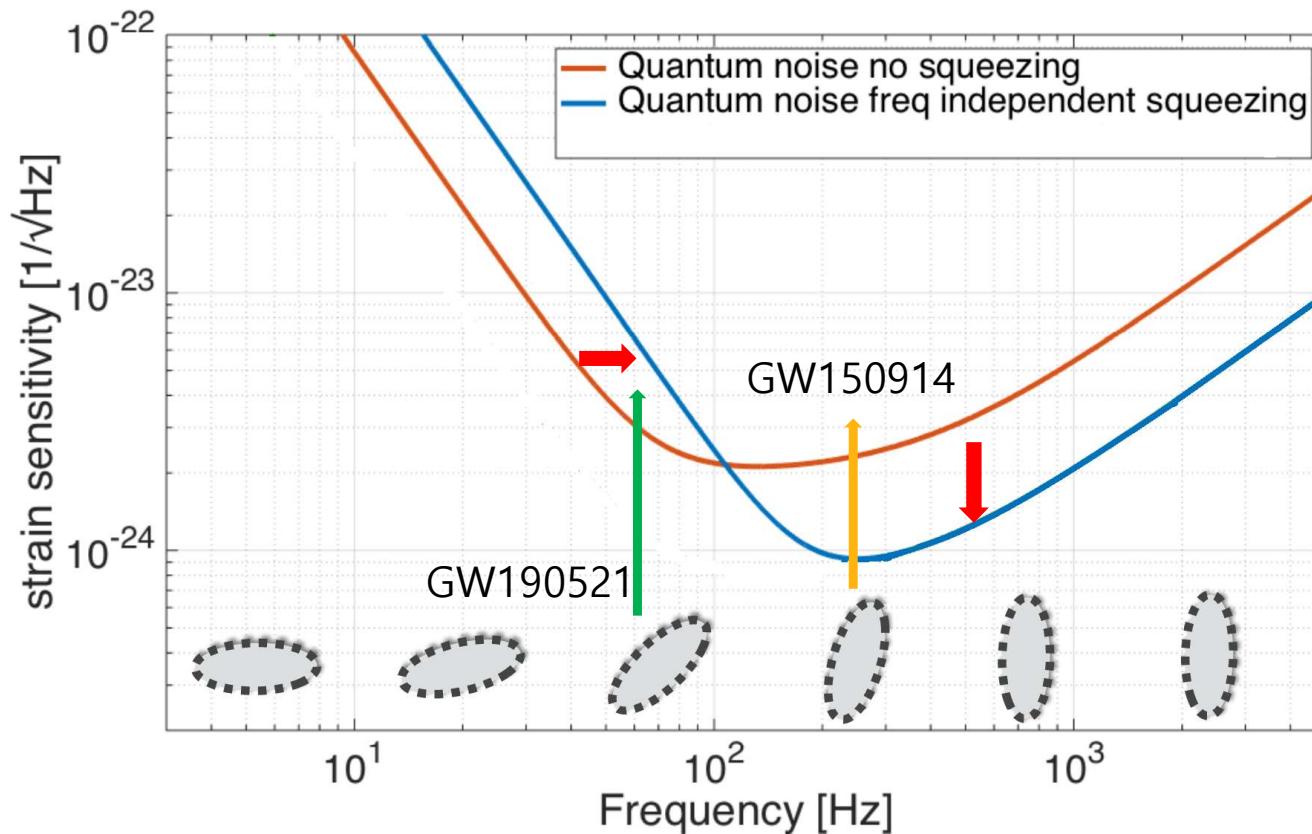
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Squeezed light



Heurs M. 2018 Gravitational wave detection using laser interferometry beyond the standard quantum limit. Phil. Trans. R. Soc. A 376: 20170 289.

Frequency independent squeezing



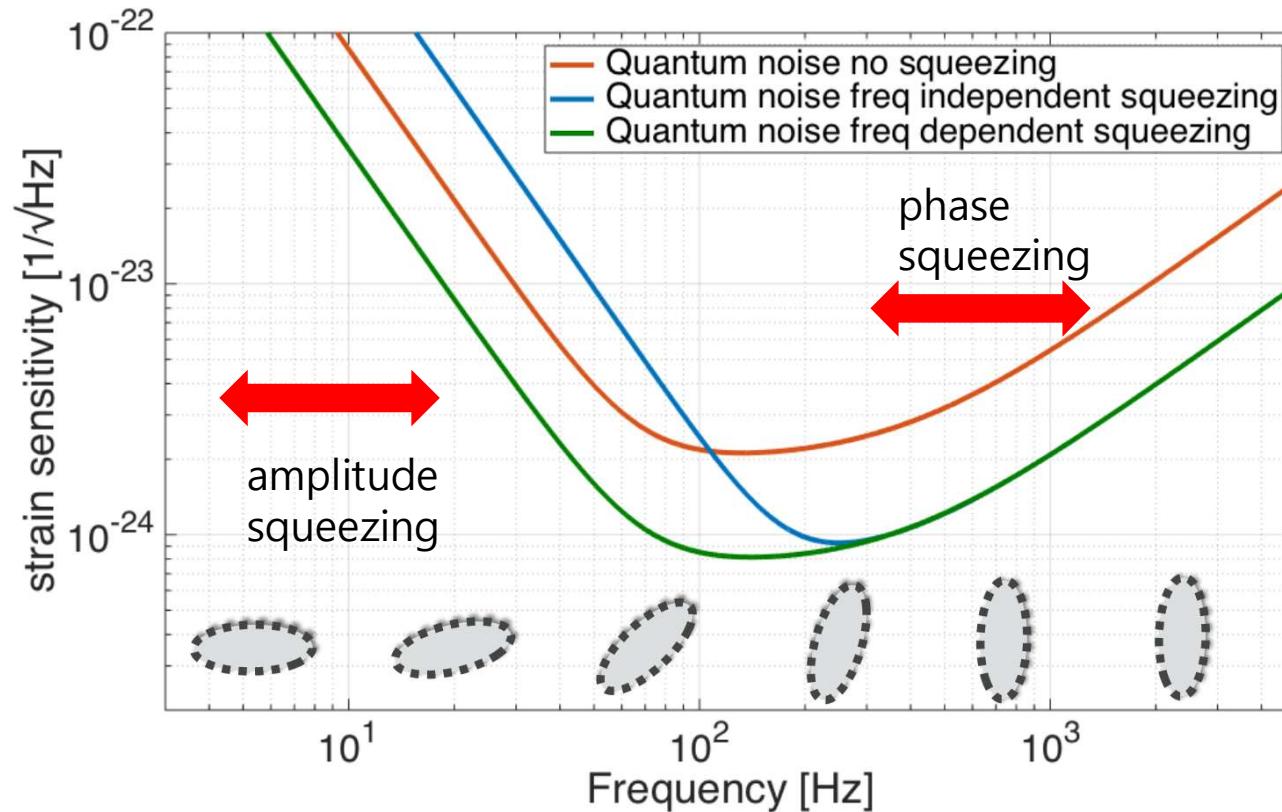
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Frequency dependent squeezing(FDS)



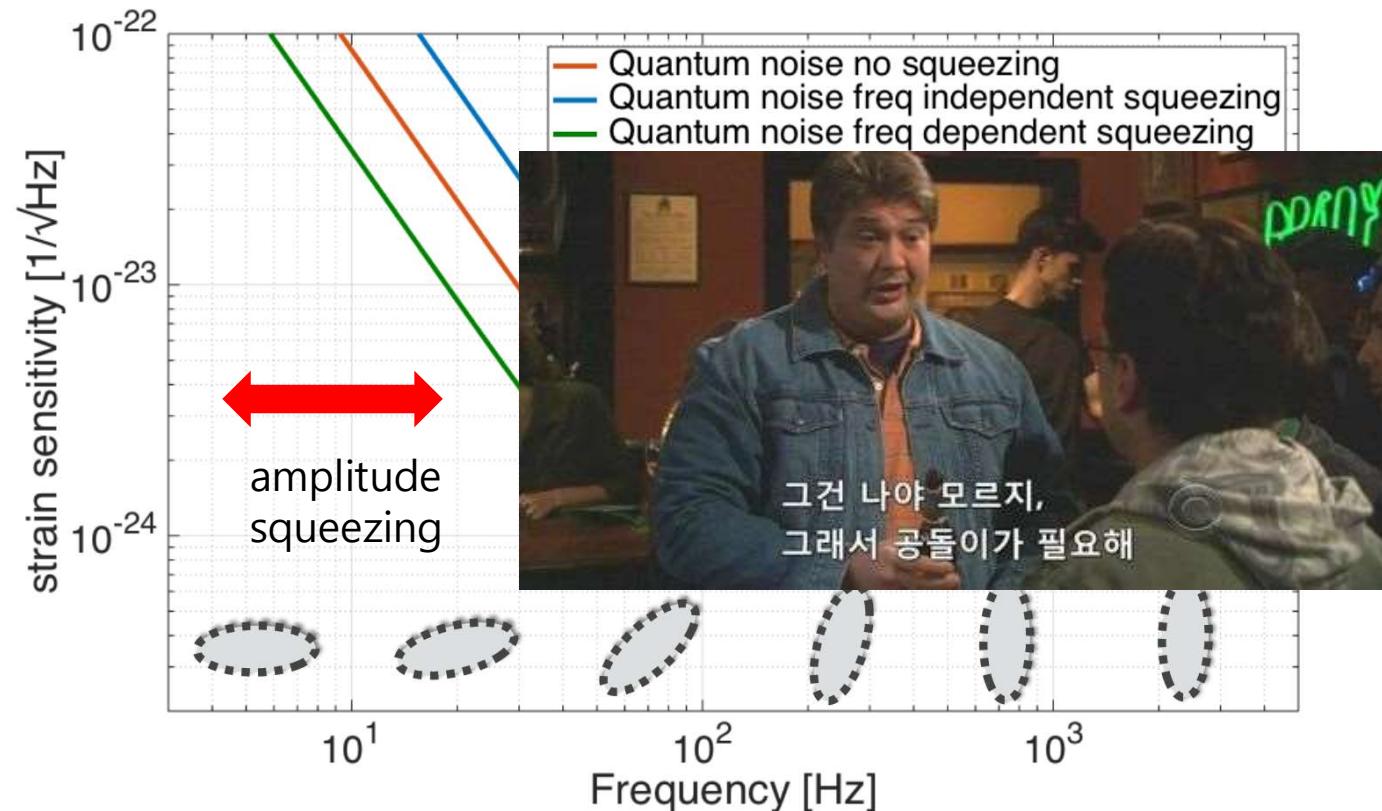
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Frequency dependent squeezing(FDS)



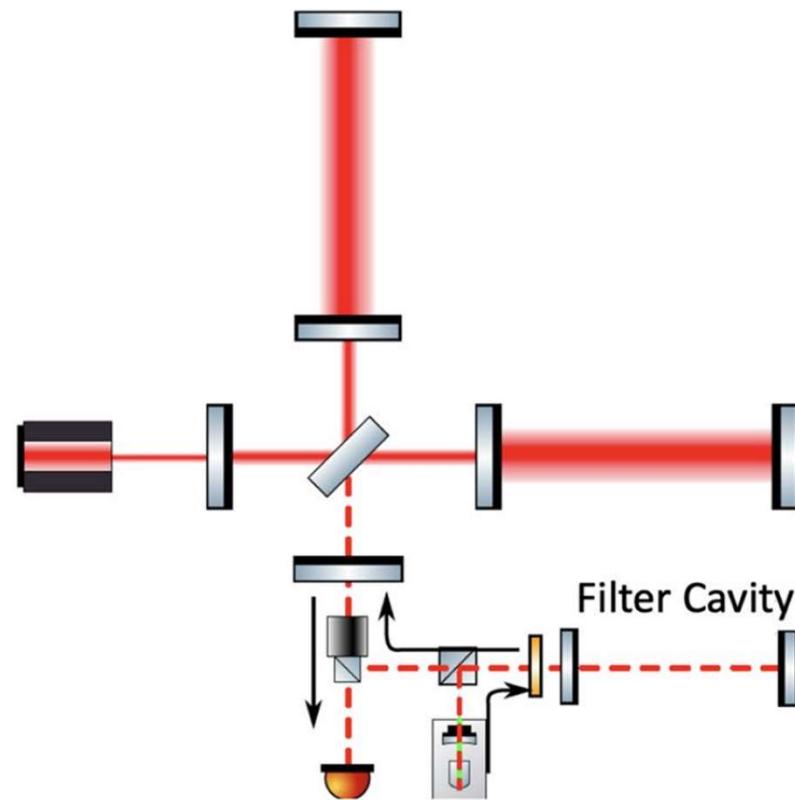
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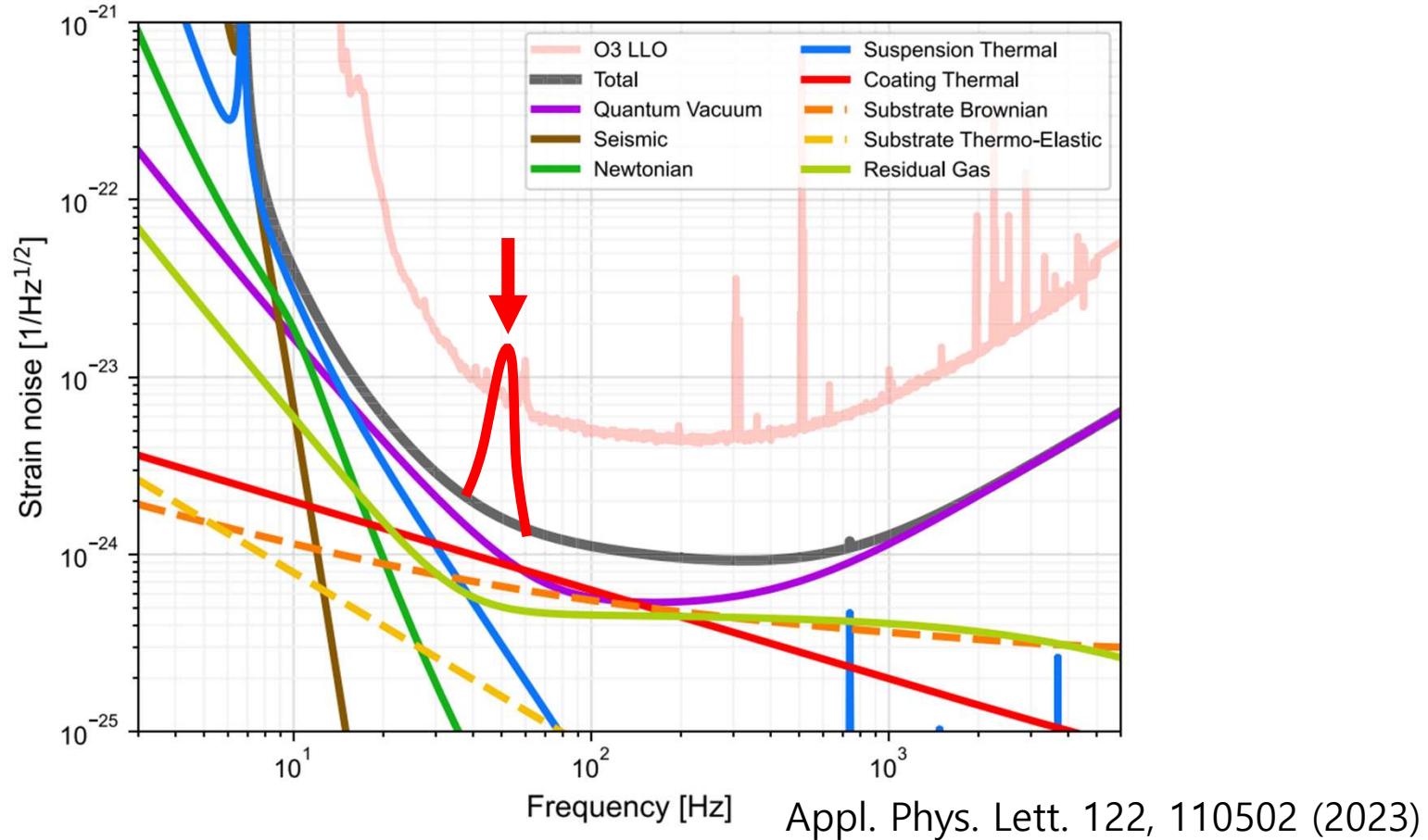
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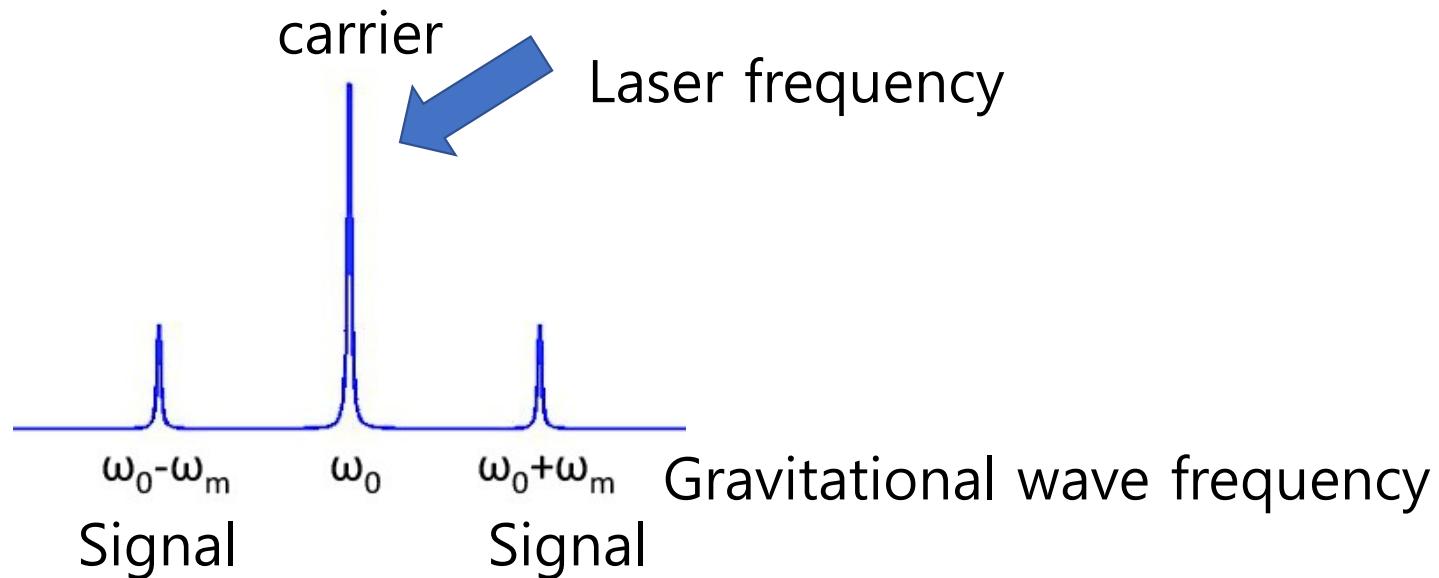
Frequency dependent squeezing using filter cavity



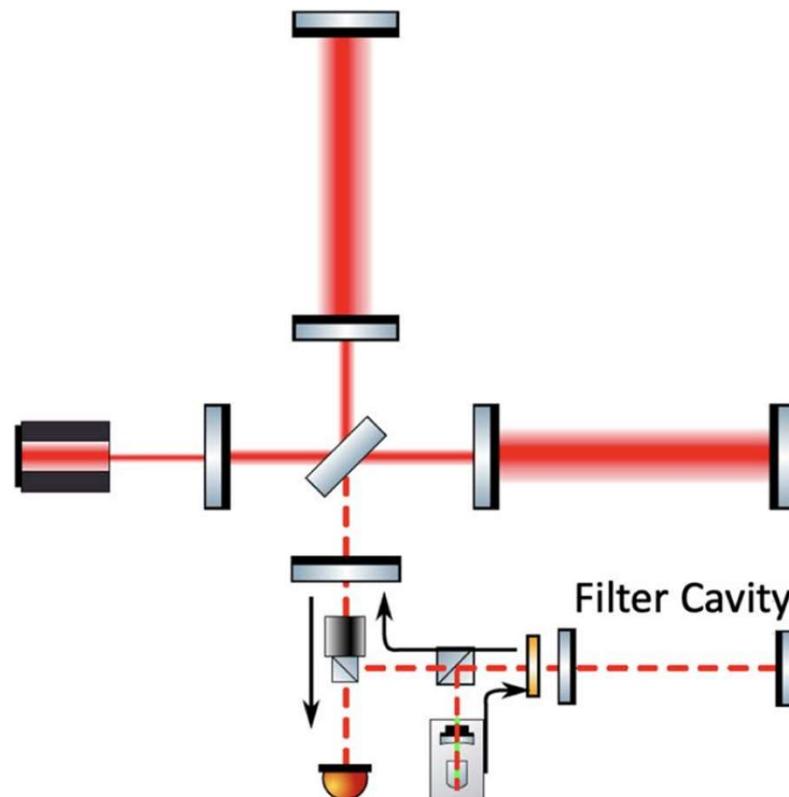
LIGO sensitivity



Side band figure

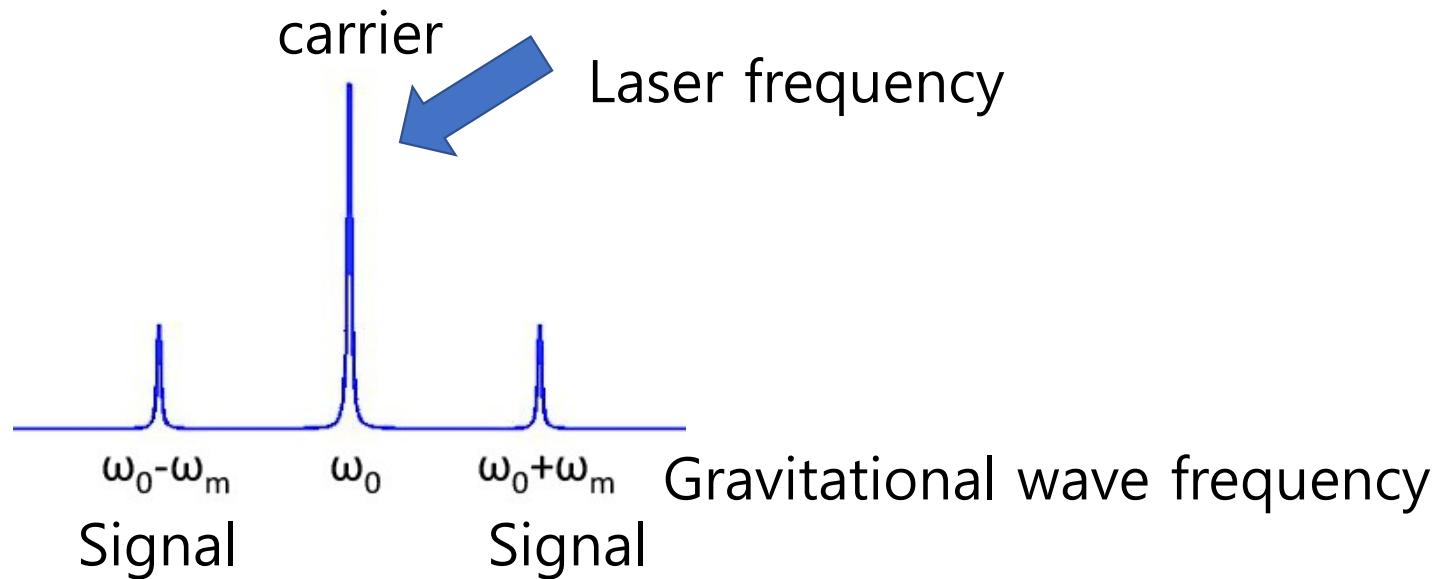


Frequency dependent squeezing using filter cavity

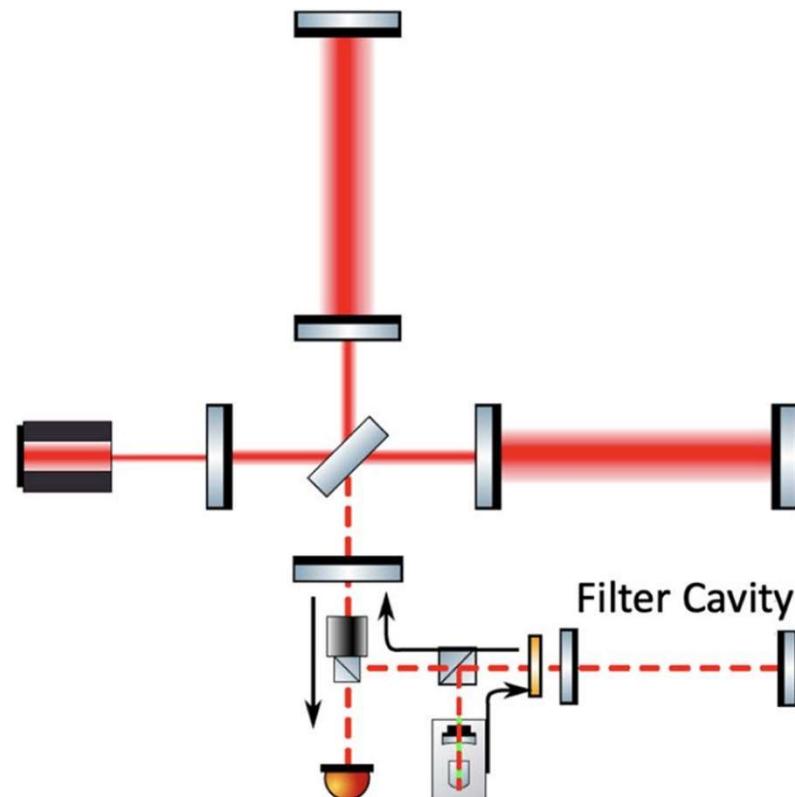


$$E^{(+)}(t) = \int_0^{\infty} E(\omega) e^{-i\omega t} \frac{d\omega}{2\pi}$$

Side band figure



Frequency dependent squeezing using filter cavity



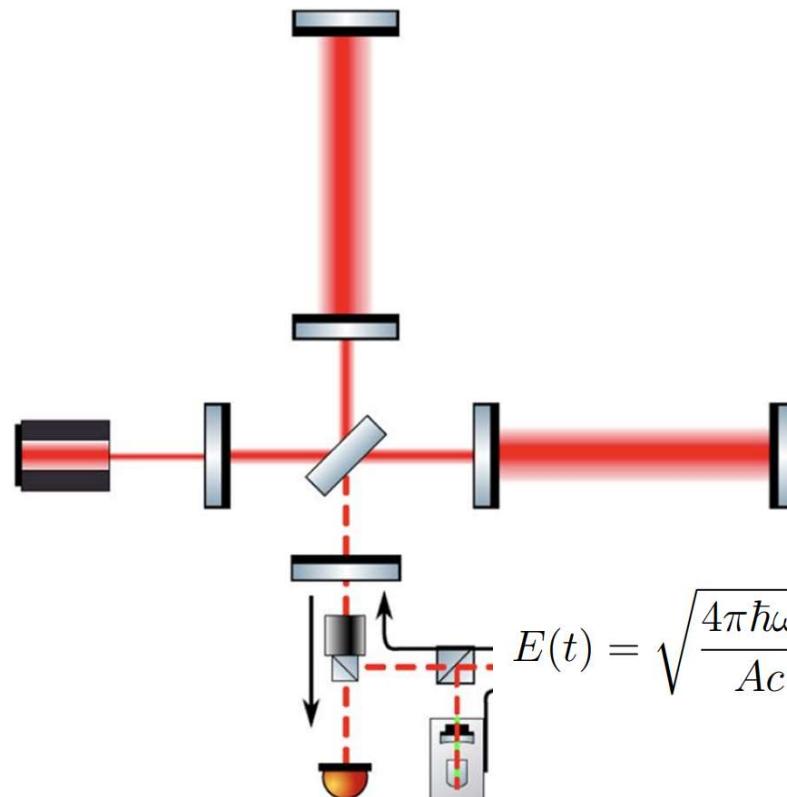
$$E^{(+)}(t) = \int_0^{\infty} E(\omega) e^{-i\omega t} \frac{d\omega}{2\pi}$$

$$E^{(+)}(t) = e^{-i\omega_0 t} \int_0^{\infty} (E(\Omega) e^{-i\Omega t} + E(-\Omega) e^{+i\Omega t}) \frac{d\Omega}{2\pi}$$

$$\omega = \omega_0 + \Omega$$

GW signal frequency
Laser frequency

Frequency dependent squeezing using filter cavity



$$E(t) = \sqrt{\frac{4\pi\hbar\omega_0}{Ac}} [a_1(t) \cos(\omega_0 t) + a_2(t) \sin(\omega_0 t)]$$

Amplitude Quadrature : $a_1(t) = \frac{a(t) + a^\dagger(t)}{\sqrt{2}}$

Phase Quadrature : $a_2(t) = \frac{a(t) - a^\dagger(t)}{i\sqrt{2}}$

$$E(t) = \sqrt{\frac{4\pi\hbar\omega_0}{Ac}} \left[\cos(\omega_0 t) \int_0^{\infty} (a_1(\Omega) e^{-i\Omega t} + a_1^\dagger(\Omega) e^{+i\Omega t}) \frac{d\Omega}{2\pi} + \sin(\omega_0 t) \int_0^{\infty} (a_2(\Omega) e^{-i\Omega t} + a_2^\dagger(\Omega) e^{+i\Omega t}) \frac{d\Omega}{2\pi} \right]$$

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아니요

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집에 가고 싶으세요?

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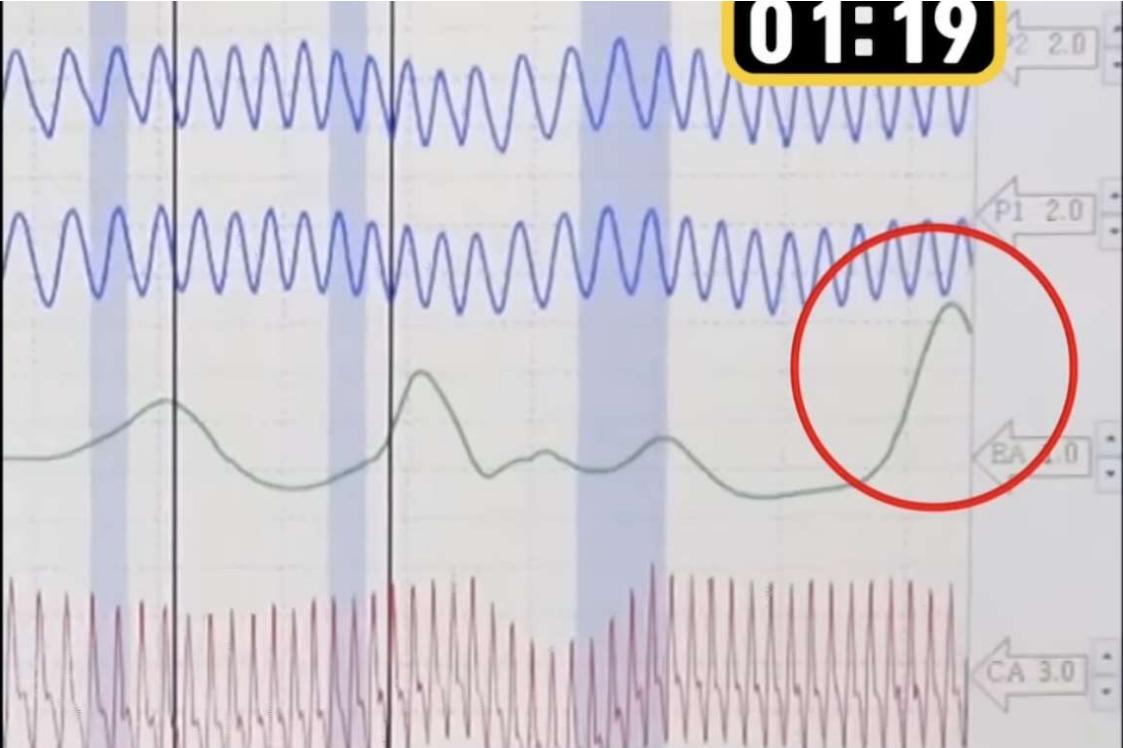
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도전酱

오봉순작

01:19



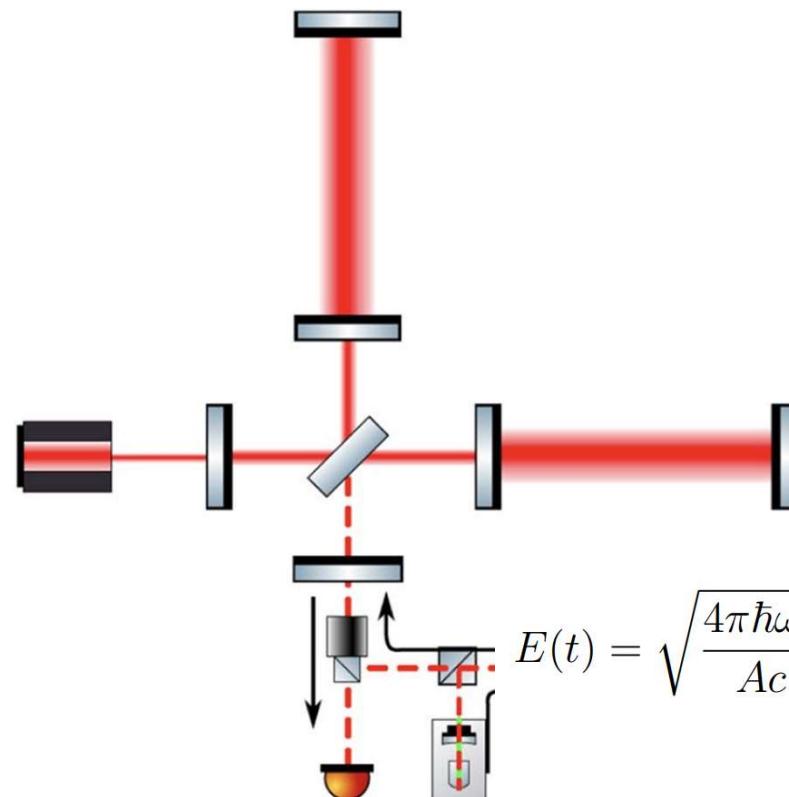
집에 올시
가고 싶음...



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Frequency dependent squeezing using filter cavity



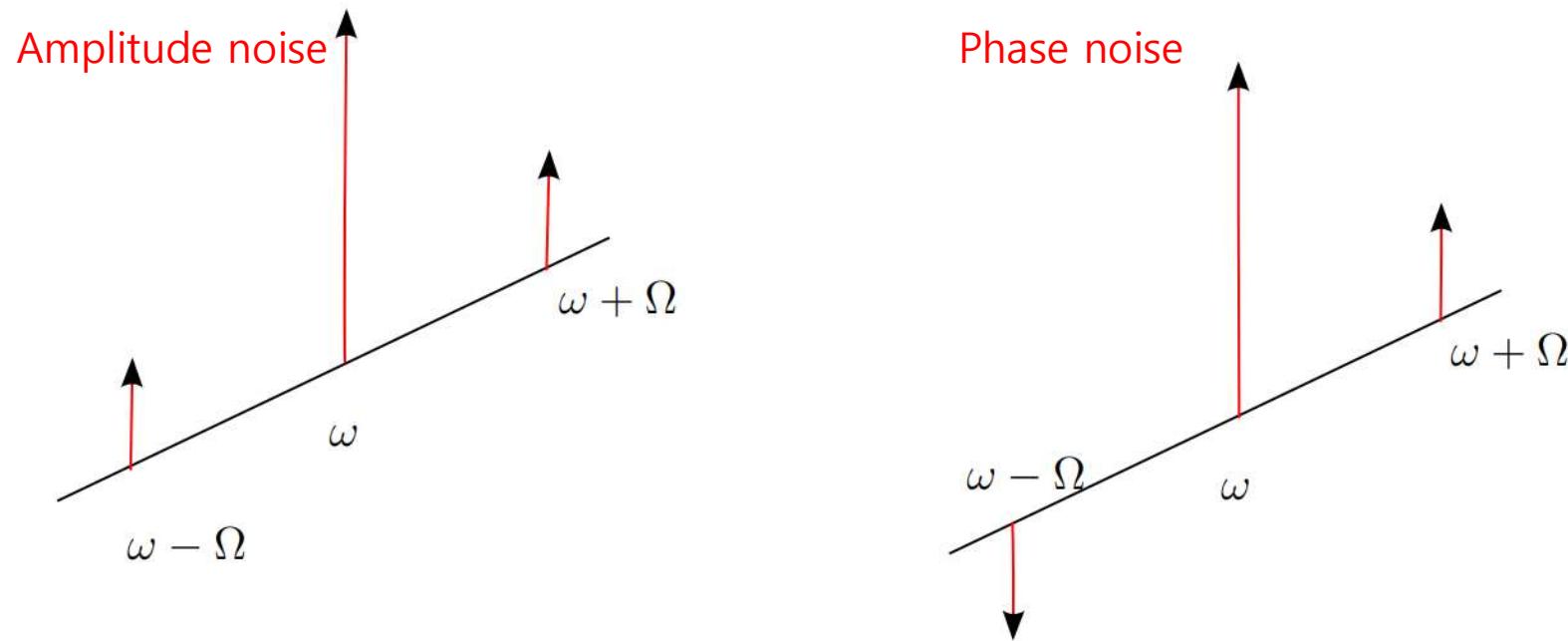
$$E(t) = \sqrt{\frac{4\pi\hbar\omega_0}{Ac}} [a_1(t) \cos(\omega_0 t) + a_2(t) \sin(\omega_0 t)]$$

Amplitude Quadrature : $a_1(t) = \frac{a(t) + a^\dagger(t)}{\sqrt{2}}$

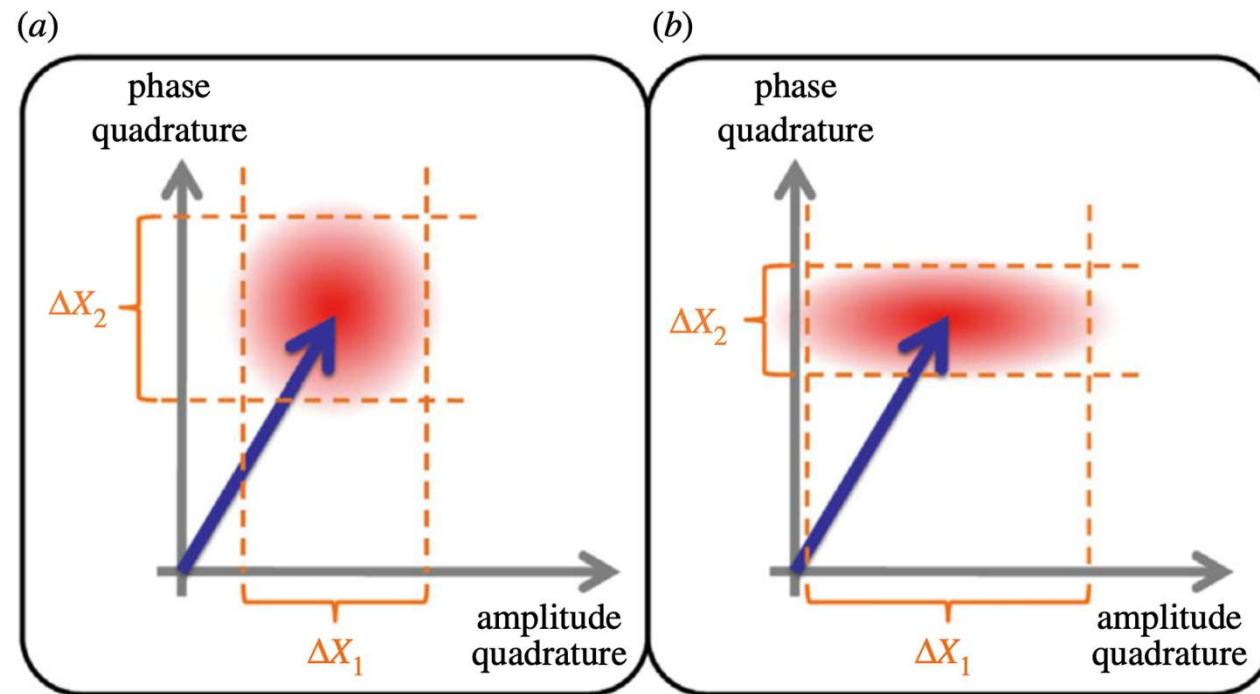
Phase Quadrature : $a_2(t) = \frac{a(t) - a^\dagger(t)}{i\sqrt{2}}$

$$E(t) = \sqrt{\frac{4\pi\hbar\omega_0}{Ac}} \left[\cos(\omega_0 t) \int_0^{\infty} (a_1(\Omega) e^{-i\Omega t} + a_1^\dagger(\Omega) e^{+i\Omega t}) \frac{d\Omega}{2\pi} + \sin(\omega_0 t) \int_0^{\infty} (a_2(\Omega) e^{-i\Omega t} + a_2^\dagger(\Omega) e^{+i\Omega t}) \frac{d\Omega}{2\pi} \right]$$

Quantum noise side band figure

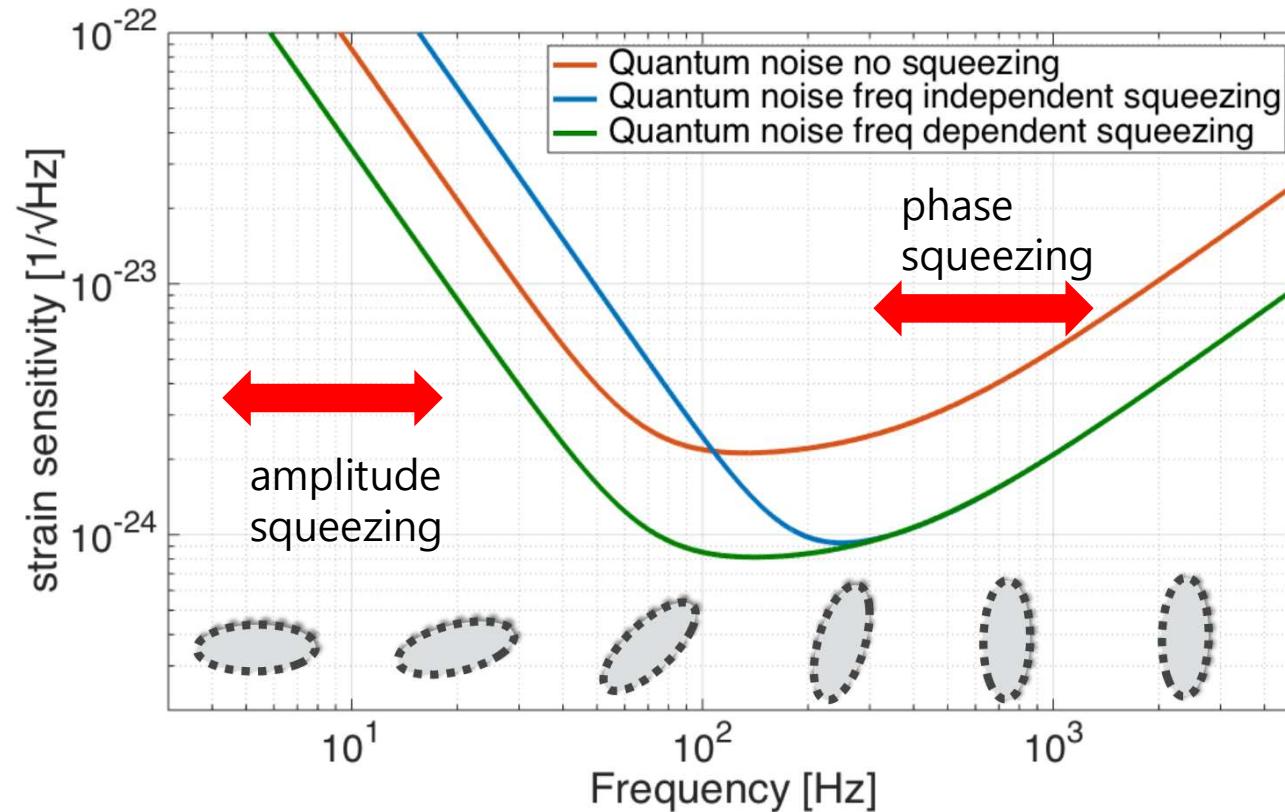


Squeezed light



Heurs M. 2018 Gravitational wave detection using laser interferometry beyond the standard quantum limit. Phil. Trans. R. Soc. A 376: 20170 289.

Frequency dependent squeezing(FDS)



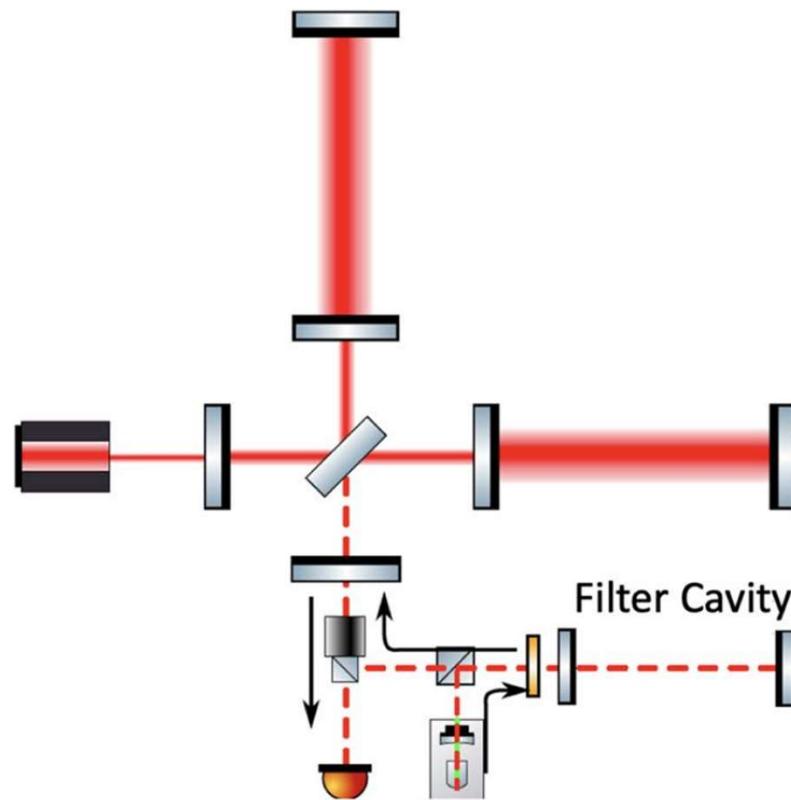
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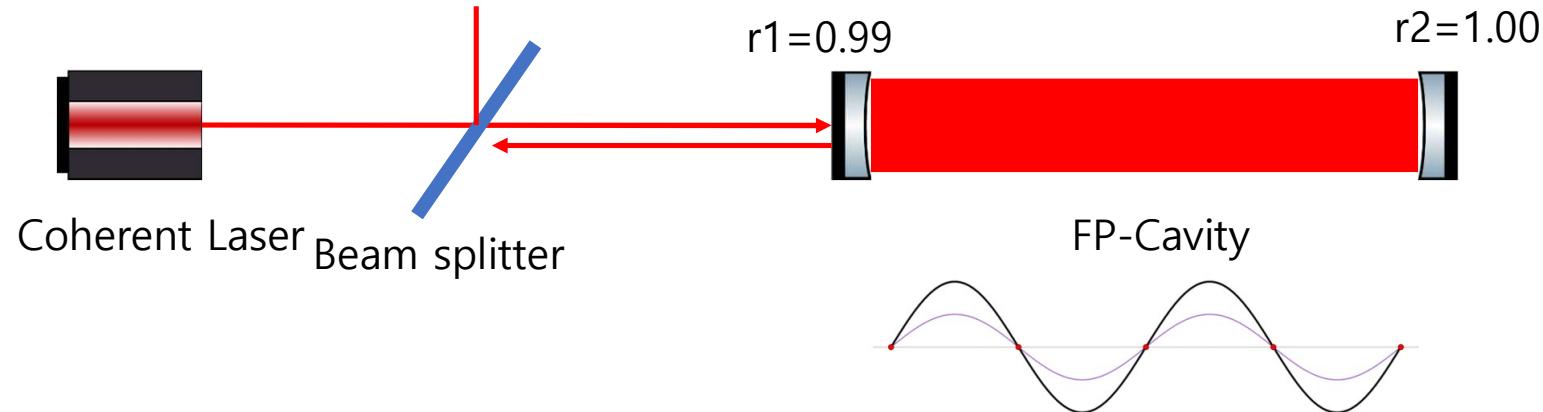
Frequency dependent squeezing using filter cavity



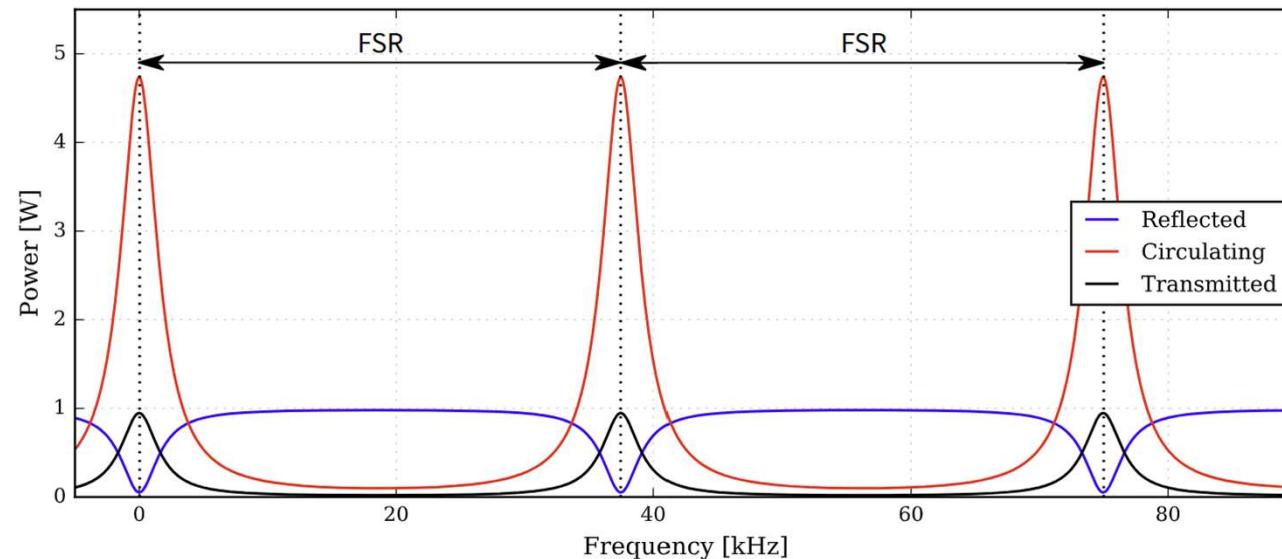
Fabry-perot cavity



Number of round trip > 250
4 km x 250 ~ 1000 km



Free spectral range of cavity



Detuned cavity



Simple picture

Stefan Hild et al, "Detuned arm cavities", 3rd GEO simulation workshop, Hannover, June 2007

B:

- less carrier light in cavity => less GW sidebands are produced.
- Since one GW sideband is resonant, it gets enhanced.

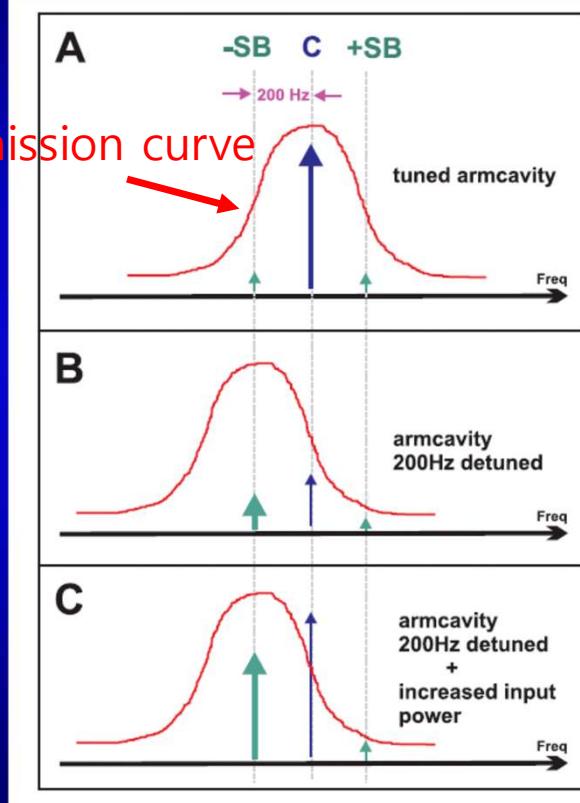
=> Smaller GW signal

C:

- optical power is restored in the cavity by larger PR-gain.
 - Same amount of GW sidebands are produced.
 - Since one GW sideband is resonant, it gets enhanced.
- Overall we win GW signal.

=> Larger GW signal

Transmission curve



Quantum noise side band figure

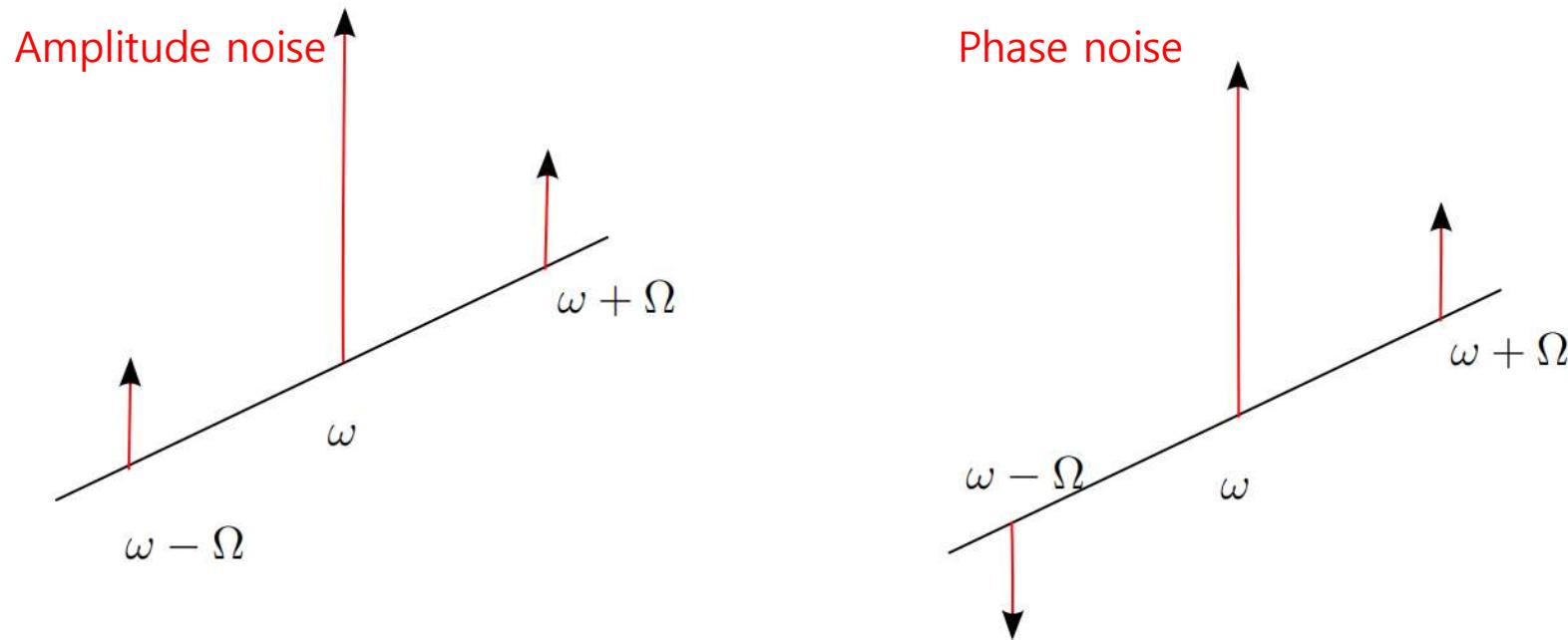
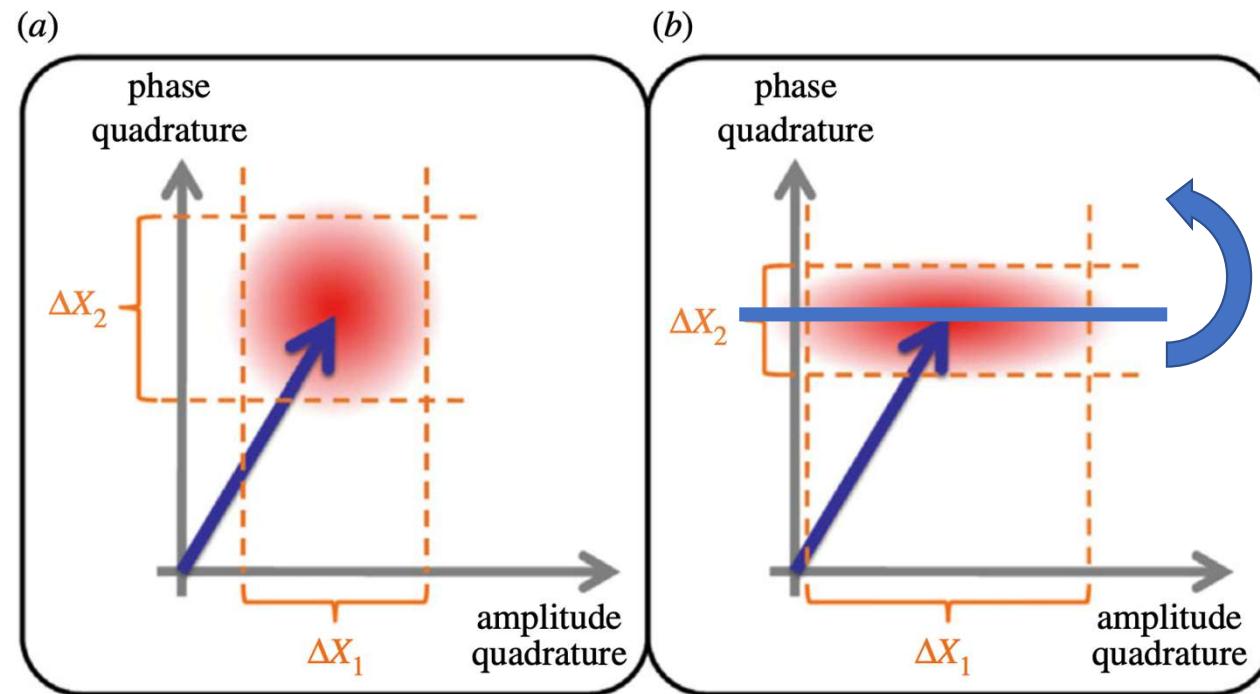


Figure 1-2: Phasors of amplitude noise (left) and phase noise (right) in the sideband picture. In the frame rotating at the carrier frequency ω the carrier is still in these diagrams while the sidebands rotate at Ω , the signal at $\omega + \Omega$ rotating clockwise while the idler at $\omega - \Omega$ rotates counter clockwise. (Sidebands have equal amplitudes)

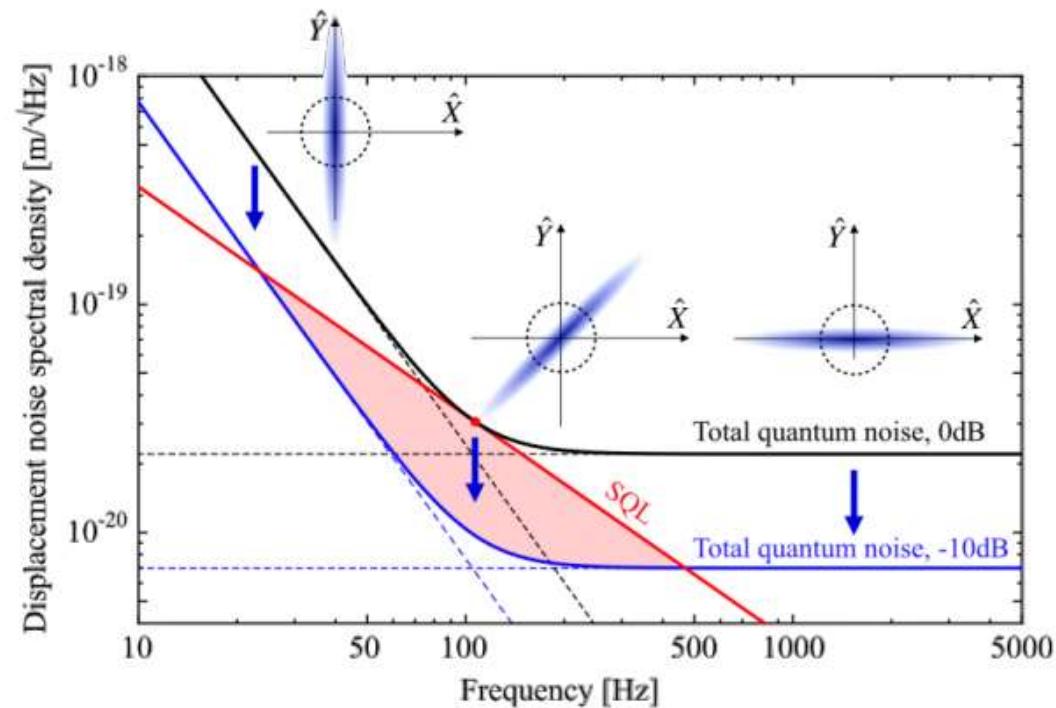
Squeezed light



Squeeze angle
rotation

Heurs M. 2018 Gravitational wave detection using laser interferometry beyond the standard quantum limit. Phil. Trans. R. Soc. A 376: 20170 289.

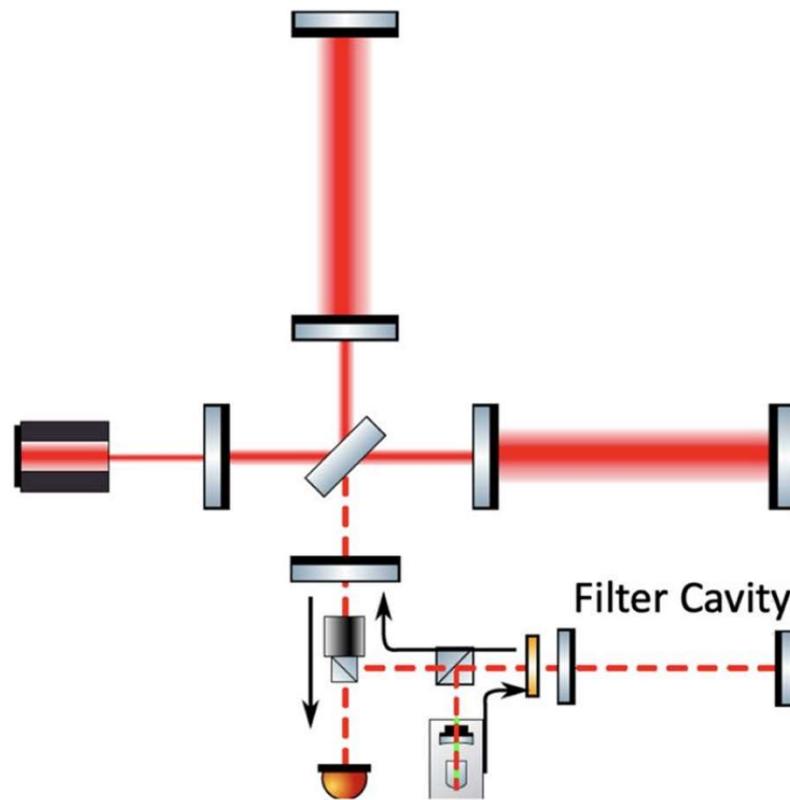
Squeeze angle rotation



$$\alpha_p = \arctan \left(\frac{2\gamma_{fc}\Delta\omega_{fc}}{\gamma_{fc}^2 - \Delta\omega_{fc}^2 + \Omega^2} \right)$$

$$t_{st} = \frac{1}{\gamma_{fc}} = \frac{\sqrt{2}}{\Omega_{SQL}} \simeq 3 \text{ ms}$$

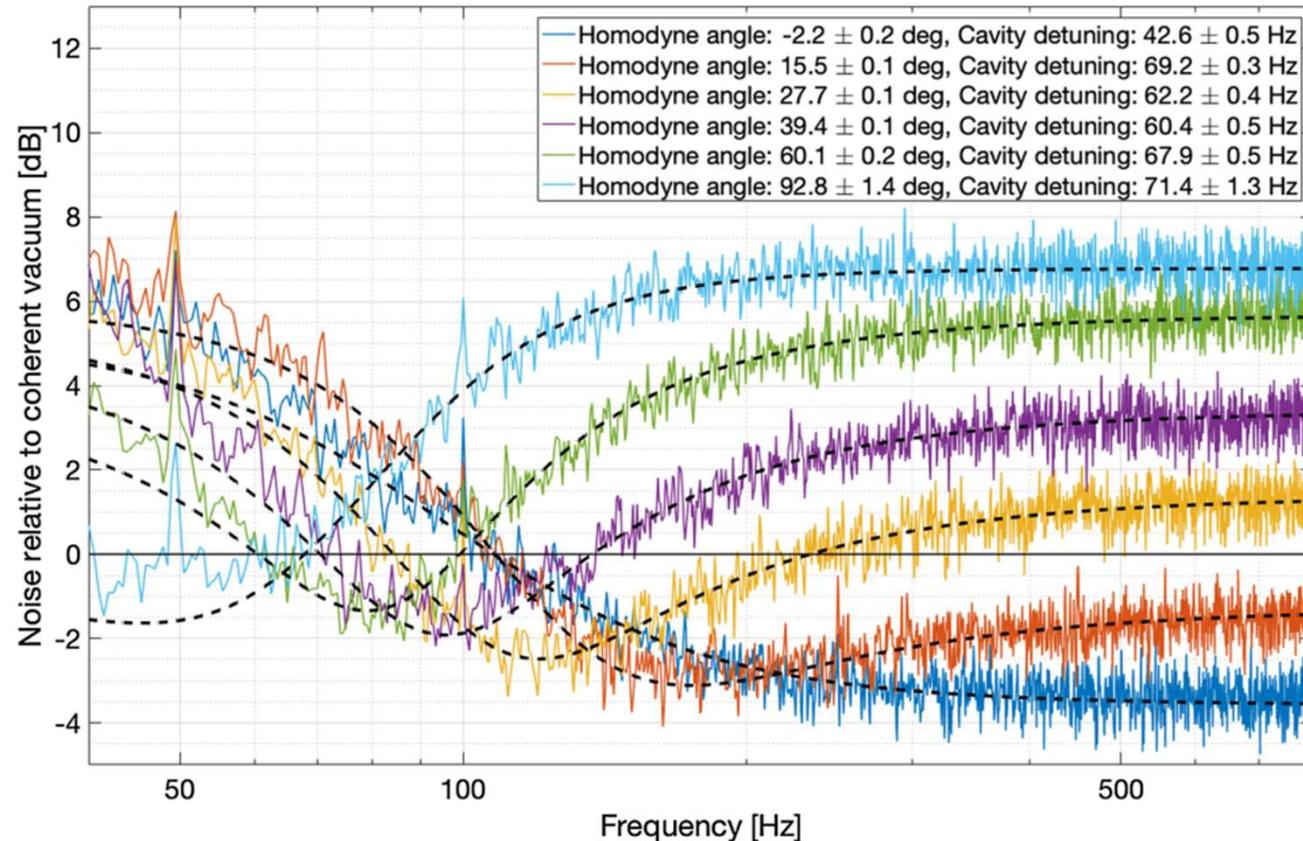
Frequency dependent squeezing using filter cavity



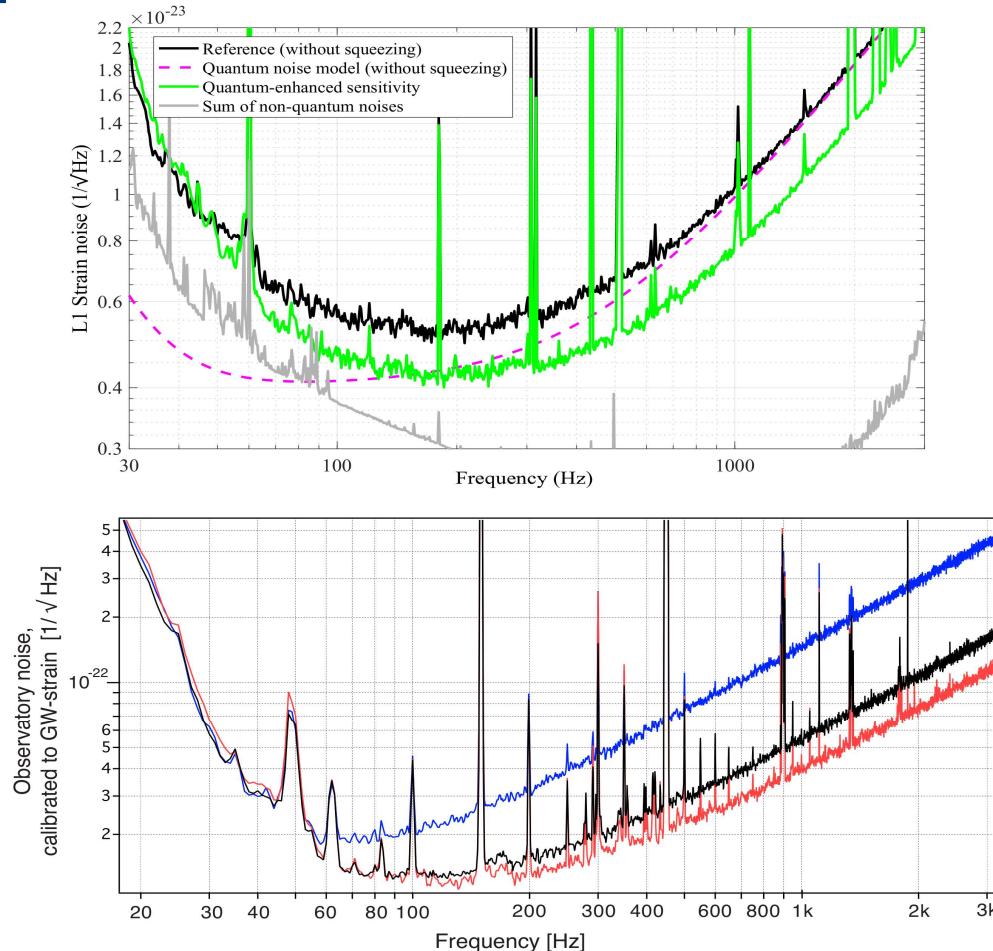
Frequency dependent squeezing - KAGRA



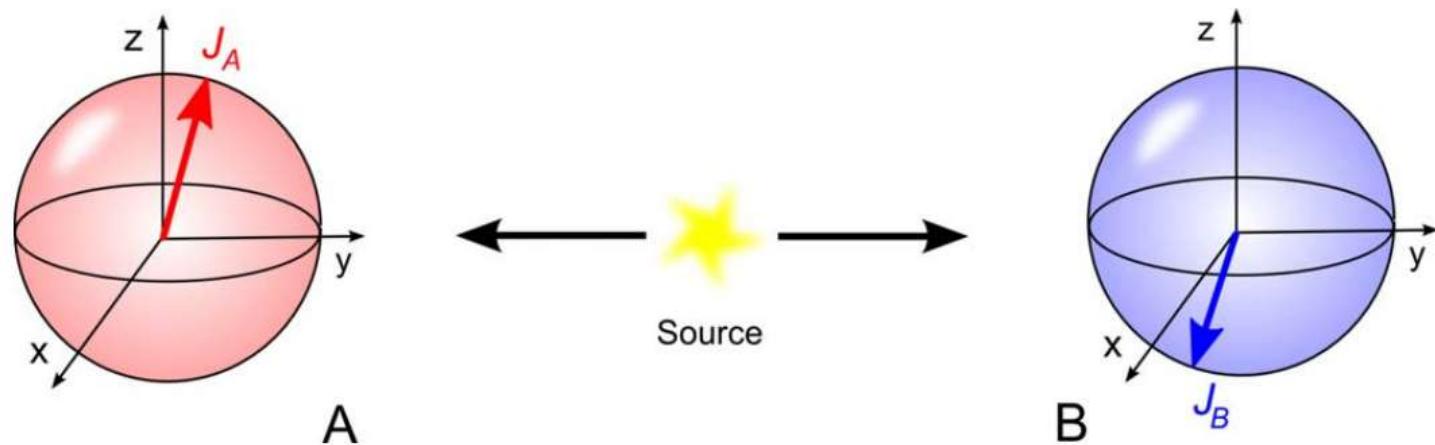
PHYSICAL REVIEW LETTERS 124, 171101 (2020)



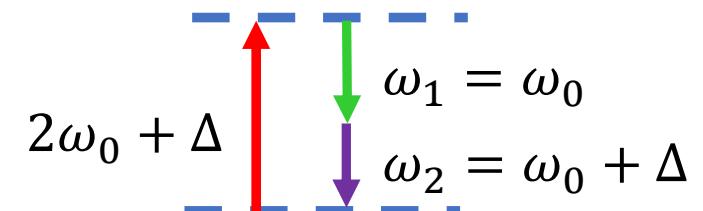
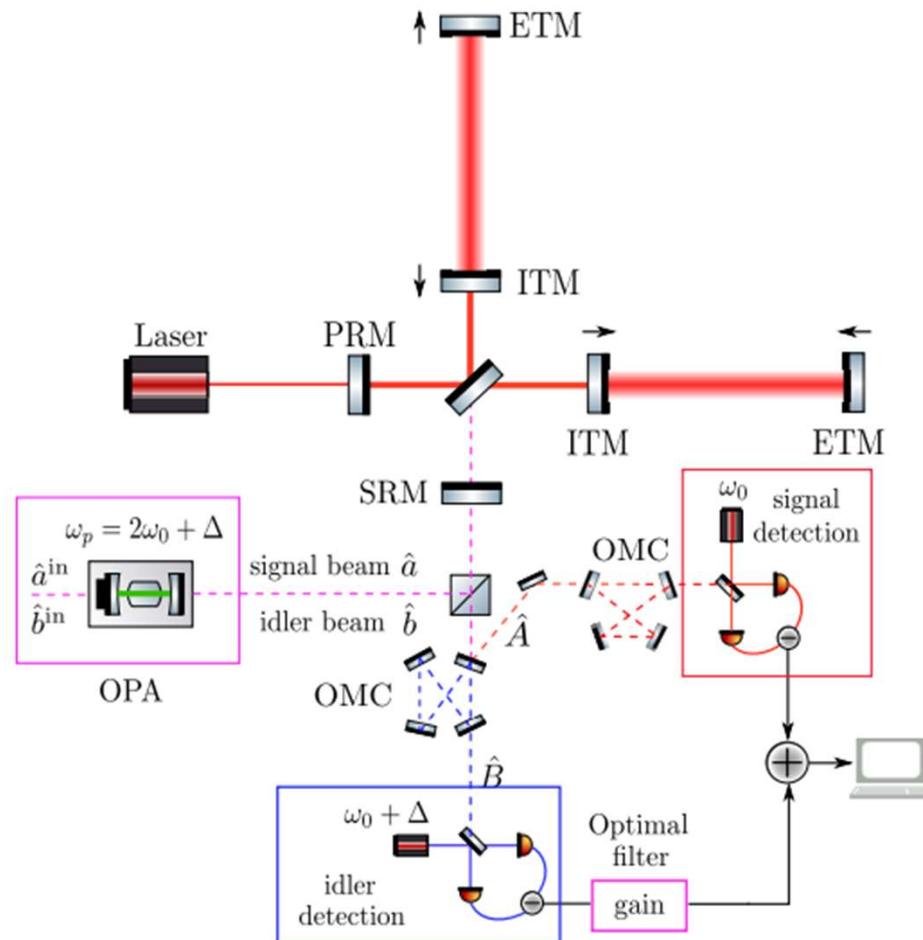
Quantum noise of gravitational wave detector



Einstein-Podolsky-Rosen paradox

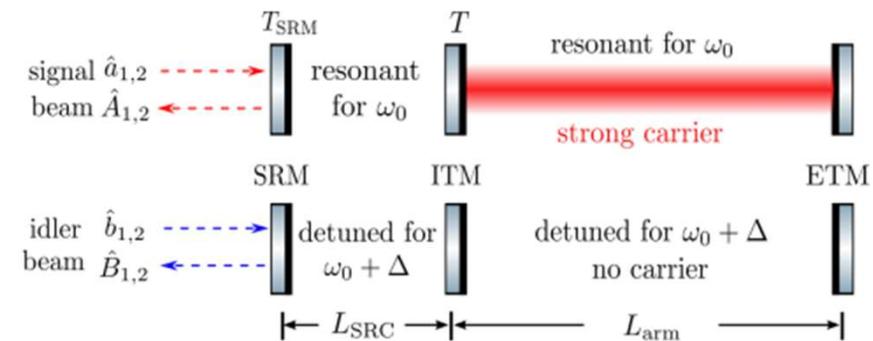
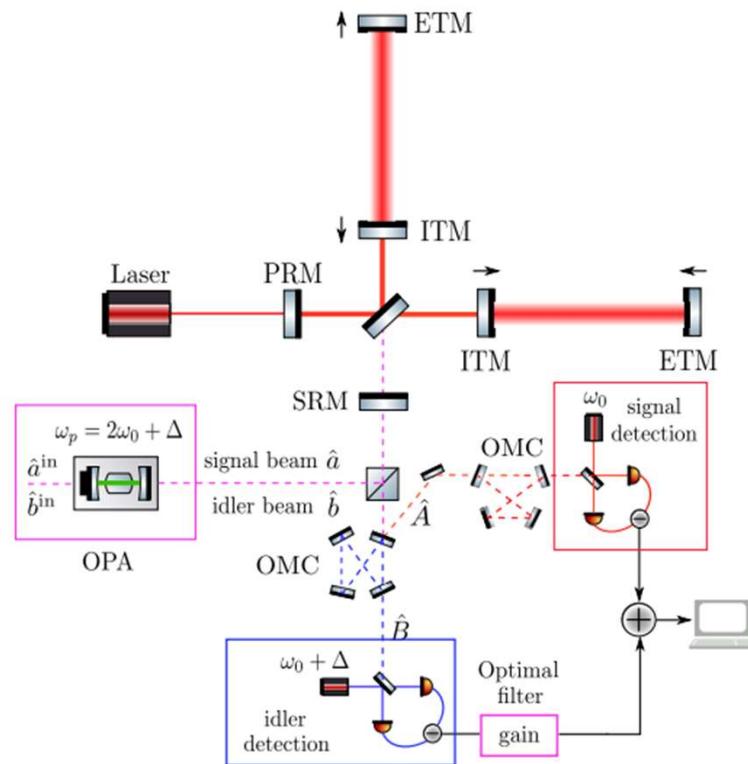


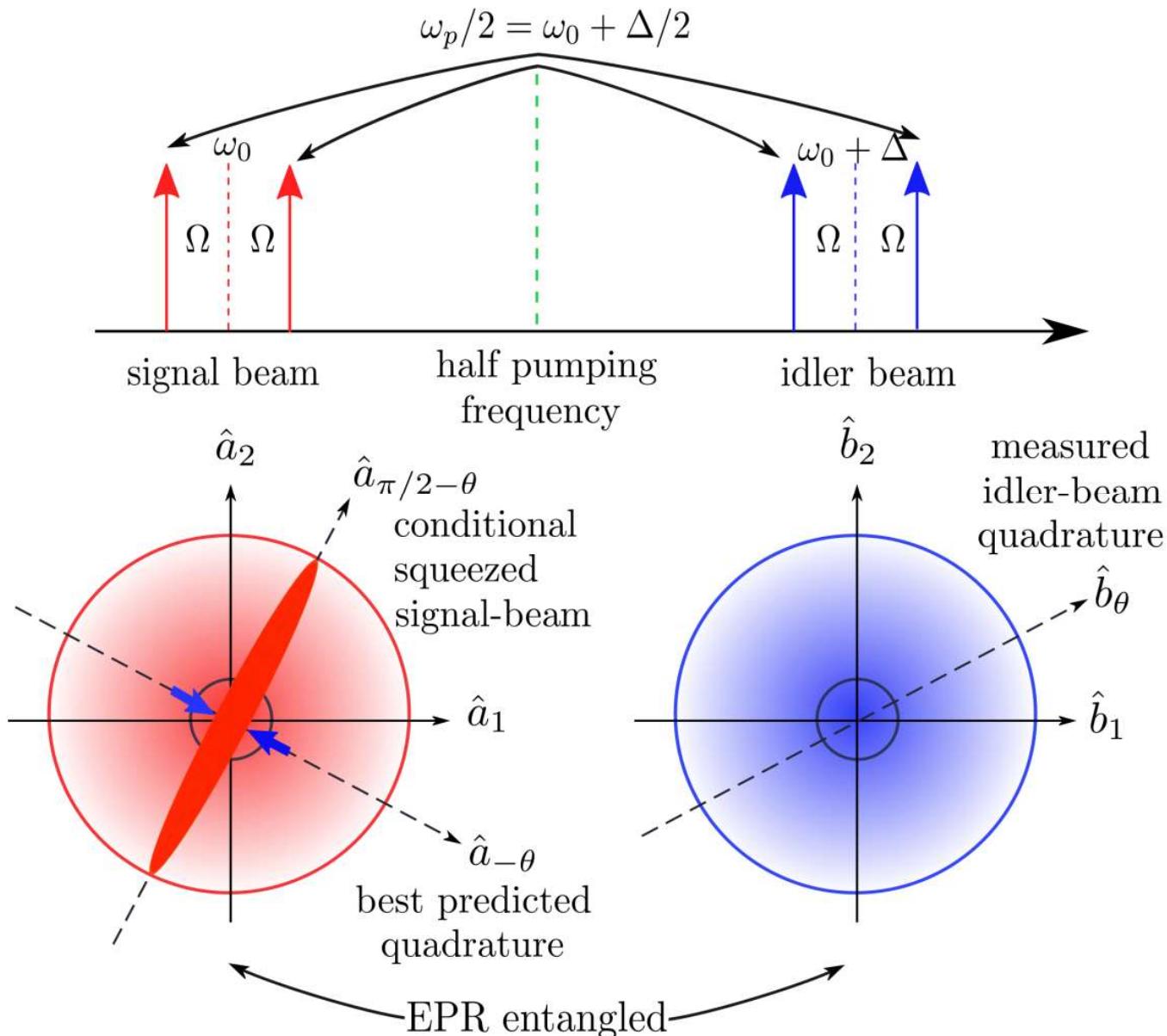
EPR squeezing for gravitational wave detector



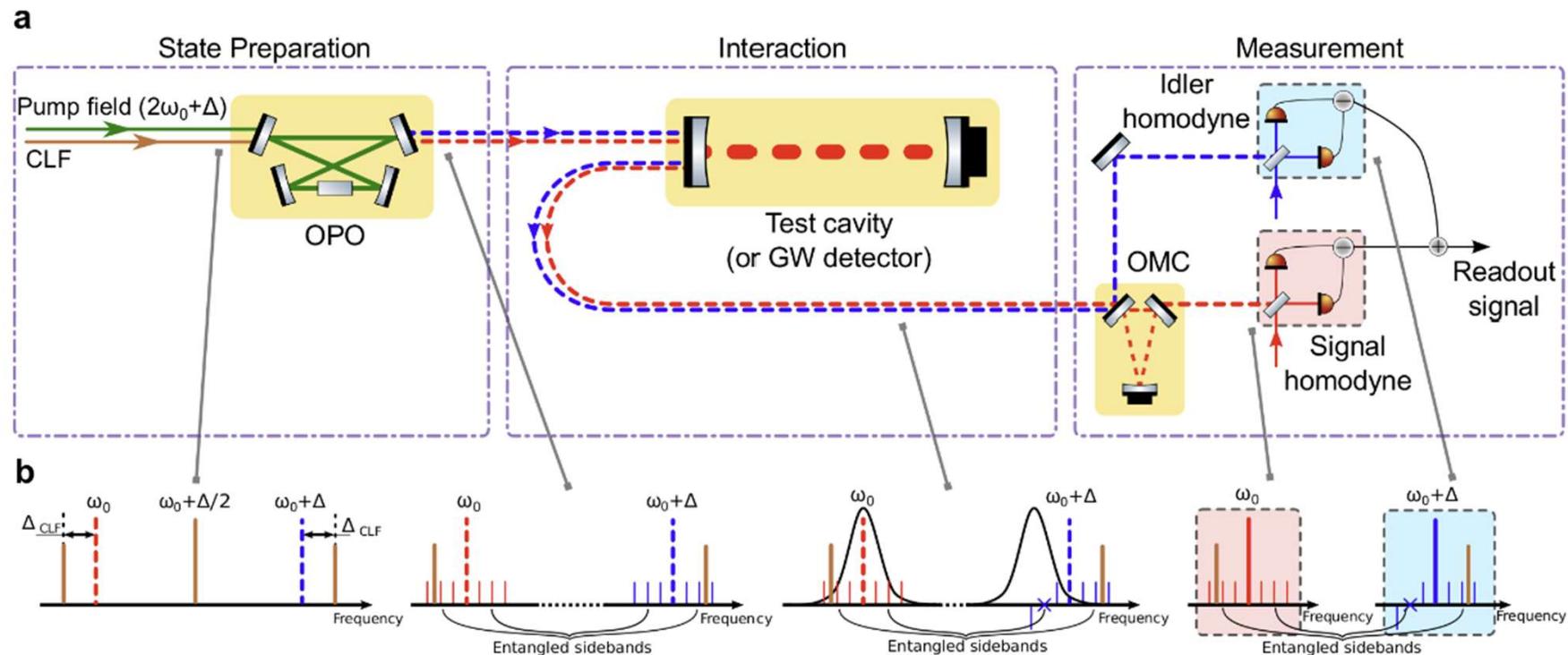
Non linear optics

EPR squeezing for gravitational wave detector





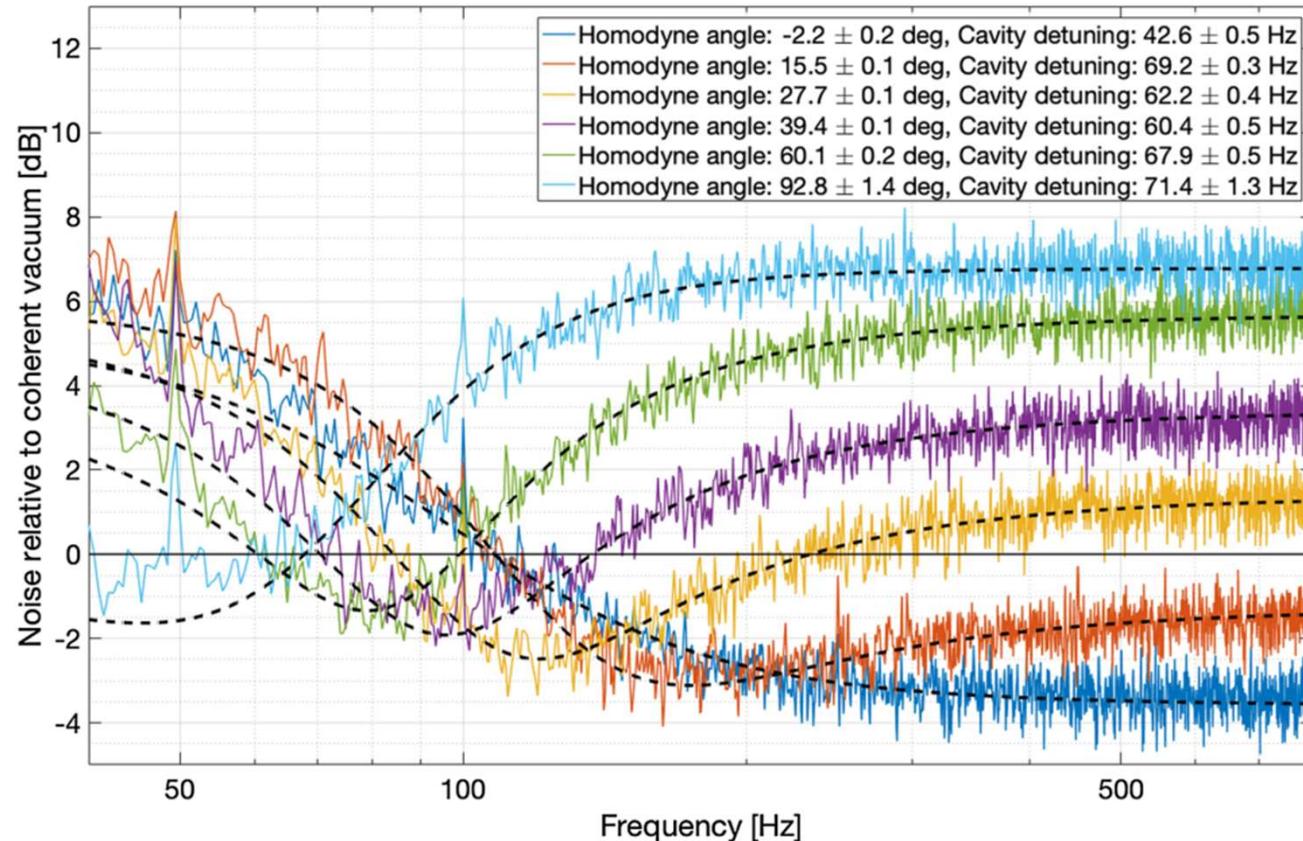
Generation and control of frequency dependent squeezing via EPR entanglement



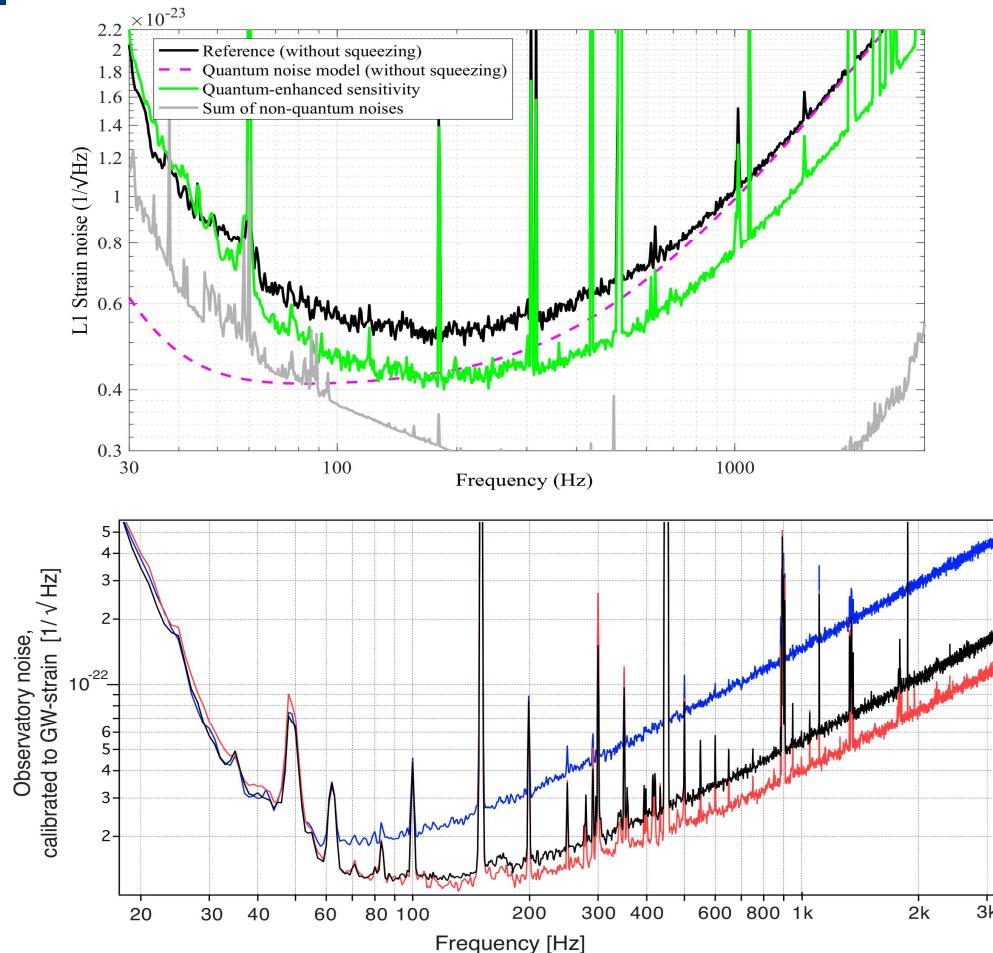
Frequency dependent squeezing - KAGRA



PHYSICAL REVIEW LETTERS 124, 171101 (2020)



Quantum noise of gravitational wave detector



도문소식

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구욱~



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